Searching for new physics in high-mass ditau events at ATLAS



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Ph.D. thesis defense, Penn

Outline

1. Introduction

Standard Model, LHC, ATLAS

- 2. TRT commissioning threshold calibration, hit efficiency
- 3. Tau performance

reconstruction, cut-based ID, pile-up robustness

4. SM $Z \rightarrow \tau \tau$

selection design, observation, cross section

5. $Z' \rightarrow \tau \tau$

fake factor method, exclusion

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Particle Physics

Fundamental questions of particle physics:

- 1. What is matter?
- 2. How does it interact?



Four fundamental forces at low energies:

- 1. Gravity
- 2. Electromagnetism
- 3. Strong force
- 4. Weak force

- very weak, no complete quantum theory
- binds atoms, chemistry
- nuclear range, binds nuclei
- nuclear range, radioactivity, solar fusion

Standard Model

- In QFT, *fields* are actually what is fundamental, and particles are localized excitations in the fields.
- Gauge symmetries determine the character of the forces between fermion fields through gauge bosons.

The SM gauge group is



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Strong

force

July 4, 2012 CERN announces the discovery of a new particle by ATLAS and CMS, consistent with the Higgs boson



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July 5 cover of the New York Times: *Physicists Find Elusive Particle Seen as Key to the Universe*

Why the Standard Model?

- Why the gauge group $SU(3)_C \times SU(2)_L \times U(1)_Y$?
- Why are there **3** generations of quarks and leptons?
- Why are lepton and hadron charges quantized in the same units? Why the existing hypercharges? $Q_{\rm EM}=T_{\rm 3L}+Y/2$

Is it because ...

- the gauge group of Nature is actually bigger? $SO(10) \rightarrow SU(5) \times U(1)$ Georgi-Glashow $SO(10) \rightarrow SU(4)_C \times SU(2)_L \times SU(2)_R$ Pati-Salam 1974
- e.g. Pati-Salam SO(10): $Q_{EM} = T_{3L} + T_{3R} + 1/2(B L)$
- Could it be that the SM is the product of a larger symmetry breaking process than just electroweak symmetry breaking?

Running of the couplings

- After precision measurements of the SM couplings at LEP, one could run the couplings according to the RGEs to higher energies.
- The SM *couplings apparently converge*, motivating the possibility of grand unification.
- Moreover, the unification seems to require Supersymmetry, but the extrapolation is over 10¹⁴ orders, and we need more experimental clues.



- Z' bosons occur in theories with additional U(1) symmetries.
- Best limits on $Z' \rightarrow ee/\mu\mu$ are

M > 2.86 TeV ATLAS [ATLAS-CONF-2013-017] M > 2.96 TeV CMS [CMS-PAS-EXO-12-061]

• Important to test the couplings to all lepton flavors. Ryan Reece (Penn)



We need high energies

- To probe physics at the TeV scale and beyond we need a high-energy collider.
- The cross section ⇒ production rate grows significantly with the collision energy, √s.
- *W*, *Z*, top, Higgs, *Z'*, ...



Overall view of the LHC experiments.

- 27 km circumference
- 1232 dipoles: 15 m , 8.3 T
- 100 tons liquid He, 1.9 K
- p-p collisions at $\sqrt{s} = 7-8$ TeV
- inst. luminosity = 10^{32} - 10^{34} cm⁻²s⁻¹

- 10¹¹ protons / bunch
- 1000 bunches/ beam
- 20 MHz , 50 ns bunch spacing
- 1-40 interactions / crossing
- 0.5×10^9 interactions / sec



ATLAS Detector

Magnets: 2T solenoid, 4T toroid barrel and end-caps **Muon Spectrometer:** $|\eta| < 2.7$, drift-tube chambers

.....

25m

T.Rex

Both ATLAS and CMS have:

- 3000 scientists, 170+ institutions
- tracking, calorimetry, muon spec.
- 100 M readout channels

44m

- 1 MB/event written at 500 Hz
- O(10⁴) TB of data/year/exp.
- world-wide grid computing

LAr electromagnetic calorimeters

Pixel detector

Transition radiation tracker



Toroid magnets

Muon chambers

Solenoid magnet

Tile calorimeters

LAr hadronic end-cap and

forward calorimeters





• Started working with Mike Hance on calibration scans and threshold normalization.





TRT threshold calibration



- Fit a parametrization to the data from a scan varying the analog-todigital thresholds in the discriminator chips on the TRT front-end.
- Wrote an algorithm for calibrating the thresholds *channel-by-channel* to reach a uniform noise occupancy by floating the functional form.

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TRT threshold calibration



- Developed a GUI making it easy for shifters to archive scans to a database for monitoring long-term detector health.
- Still used in the regularly scheduled calibration periods between beam fills.
- Supported TRT as part of DAQ on-call team.^{II} Ryan Reece (Penn)



TRT hit efficiency



• $\varepsilon(x) = \frac{n_{\text{hits}}(x)}{n_{\text{hits}}(x) + n_{\text{holes}}(x)}$

- *Hits* are easy to count directly from the data.
- *Holes* are counted by extrapolating along a track to determine which straws were crossed.
- Extrapolation tools use a detailed description of detector material to stochastically model bremsstrahlung and multiple scattering.

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'09-'10

TRT hit efficiency

'09-'10



- I wrote an algorithm that uses extrapolator tools to calculate the TRT straw-hit efficiency.
- Gives an important data/MC comparison to test the TRT digitization, the step in the MC production where the response detector and electronics are simulated.
- **Published in the first JHEP paper** documenting the ATLAS detector performance with the first 900 GeV collision from 2009.



Tau performance

 π

 π

What's a tau?

- Only lepton massive enough to decay hadronically.
- Decay in beam pipe: $c\tau \approx 87 \ \mu m$
- 65% hadronic
 50% 1-prong, 15% 3-prong.
- **Signature:** narrow jet with 1 or 3 tracks, possibly additional EM clusters.
- Challenge: large multijet background at hadron colliders.
- **Importance:** often preferred coupling to new physics:

SM $H \rightarrow \tau \tau$, $H^+ \rightarrow \tau^+ \nu$, $Z' \rightarrow \tau \tau$, high-tan β SUSY,...









Timeline of taus at ATLAS

- Nov 2010: Obs. of W→τν (546 nb⁻¹)
- Feb 2011: Obs. of Z→ττ
 (8.5 pb⁻¹)
- July 2011: $W \rightarrow \tau \nu$ and $Z \rightarrow \tau \tau$ cross section measurements (36 pb⁻¹)
- Feb 2012: $Z \rightarrow \tau \tau$ cross section (1.5 fb⁻¹)
- June 2012: SM H→ττ excluded 3-4×SM at m_H≈125 GeV [arXiv:1206.5971]
- Several other analyses: MSSM $H \rightarrow \tau \tau$, $t\bar{t}$ with τ , $H^+ \rightarrow \tau \nu$, $Z' \rightarrow \tau \tau$, SUSY $\tau + MET$, ...



• Now eagerly waiting to see if $H \rightarrow \tau \tau$ will be excluded at 1×SM this year?

Tau reconstruction

- Seeded by anti-k_t jets (R=0.4) of
 3-D topological calorimeter clusters.
- 2. **Define the four-momentum** as the jet-axis with a tau-specific calibration.
- 3. **Associate tracks** with the jet that are consistent with the chosen vertex.
- 4. **Calculate discriminating variables** from the combined calorimeter and tracking information, later used to identify hadronic tau decays with BDT and likelihood based discriminants.



Cut-based tau ID

Safe Cut Variables

 $\times 10^3$

- My timing with the development of tau reconstruction software and the arrival of first collision data allowed me to contribute to the *first data/MC comparisons* of tau ID variables and *develop the first cut-based ID* used with ATLAS data.
- Prefers narrow calorimeter deposits and closely associated tracks.
- Used in the *first observation* of $W \rightarrow \tau \nu$ and $Z \rightarrow \tau \tau$.



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2010

- 1 $R_{\rm EM} = \frac{\sum \Delta R E_{\rm T}}{\sum E_{\rm T}}$, summed over cells from the first 3 layers of the EM calorimeter within $\Delta R < 0.4$.
- 2 $R_{\text{track}} = \frac{\sum \Delta R p_{\text{T}}}{\sum p_{\text{T}}}$, summed over tracks associated to the tau within $\Delta R < 0.2$.
- 3 $f_{trk,1} = \frac{p_T(\text{lead track})}{p_T(\tau_h)}$, the transverse momentum fraction of the leading track.

[ATLAS-CONF-2010-012, ATLAS-CONF-2010-059, ATLAS-CONF-2010-086, ATLAS-CONF-2011-010] 24

E_T-parametrized ID

Lorentz boost implies R should collimate as

$$R(E_{\rm T}) \propto \frac{1}{E_{\rm T}}$$

- Multiplying by E_T flattens out E_T -dependence.
- Example plots here are for 3-prong, but the π^0 s of 1-prong taus also collimate.
- Construct a smooth family of curves between the signal and background that have efficiency that is approximately flat in E_T.

$$R_{\rm cut}(E_{\rm T}; x_{\rm cut}) E_{\rm T} = (1 - x_{\rm cut}) f_{\rm sig}(E_{\rm T}) + x_{\rm cut} f_{\rm bkg}(E_{\rm T})$$

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Efficiency / Rejection

Simple cuts



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Pile-up



- 1-40 pile-up interactions / crossing
- The additional tracks and clusters
 Mean Number of Interactions per Crossing

 from pile-up are especially challenging for tau identification, which
 discriminates hadronic tau decays from jets with isolation-related
 track and calorimeter quantities.
- Efforts in 2011-2012 involved re-defining or adding corrections to identification variables to be more robust against the increasing pile-up.

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ATLAS Online Luminosity

 $\sqrt{s} = 8 \text{ TeV}, \int \text{Ldt} = 14.0 \text{ fb}^{-1}, \langle u \rangle = 20.0$ $\sqrt{s} = 7 \text{ TeV}, \int \text{Ldt} = 5.2 \text{ fb}^{-1}, \langle u \rangle = 9.1$

Identification and pile-up

2011

Important offline variable in 2010-2011:
 EM radius - "width of jet in calorimeter"

$$R_{\rm EM} = \frac{\sum_{\{\Delta R < 0.4\}} E_{\rm T}^{\rm EM}(\text{cell}) \,\Delta R(\text{cell, jet})}{\sum_{\{\Delta R < 0.4\}} E_{\rm T}^{\rm EM}(\text{cell})}$$

• Strong pile-up dependence due to using calorimeter deposits in the wide cone: $\Delta R < 0.4$.



Pile-up robust variables

2011



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[ATLAS-CONF-2011-152]

0.9

f_{core}

Pile-up corrections

 Also beginning in 2012, the variables with the largest pile-up dependence (*f*_{core} and *f*_{track}) are corrected with terms that are linear in the number of reconstructed vertices.

$$f_{\text{core}} = \frac{\sum_{\{\Delta R < 0.1\}} E_{\text{T}}^{\text{EM}}(\text{cell})}{\sum_{\{\Delta R < 0.2\}} E_{\text{T}}^{\text{EM}}(\text{cell})} + (0.3\%/\text{vertex}) \times N(\text{vertex})$$

Tight/Medium/Loose working points of the BDT and LLH are defined
 (≈40%, 60%, 70% efficient), optimized as function of p_T and in separate N(vertex) categories.



SM $Z \rightarrow \tau \tau$



$Z \rightarrow \tau \tau$ studies

- Focus on lep-had final state. Trigger on e or μ .
- Able to select $Z \rightarrow \tau \tau$ control sample for *studying tau ID*.
- Important to establish understanding of this *irreducible background to new physics* with taus: $H \rightarrow \tau \tau$ and $Z' \rightarrow \tau \tau$ searches.
- Complicated background composition.
 Multijet, W+jet, and Z+jet backgrounds all compete at the same order.

Data-driven background estimates:

- Jet fake rate mis-modeled in ATLAS Monte Carlo. W+jet background normalized with high-m_T control region.
- Multijet background modeled from same-sign data.

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'09-'11

$\sum \cos \Delta \phi$ for suppressing W+jet



$Z \rightarrow \tau \tau$ cross section

'10-'11



- Claimed observation of $Z \rightarrow \tau \tau$ with 8.5 pb⁻¹.
- Measured cross section to 10% with 36 pb⁻¹, consistent with SM.
- Published in PRD.

 $\sigma_{\text{combined}} = 0.97 \pm 0.07 (\text{stat.}) \pm 0.07 (\text{sys.}) \pm 0.03 (\text{lumi.}) \text{ nb}$

 $\sigma_{\rm theory} = 0.96 \pm 0.05 \text{ nb}$ at NNLO

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[Phys. Rev. D 84, 112006 (2011) arxiv:1108.2016v2] 34

New physics: $Z' \rightarrow \tau \tau$



$Z' \rightarrow \tau \tau \rightarrow \mu \tau_h$

- Select OS back-to-back tau decays.
- Count high-mass events.

Event selection

- $p_{\rm T}(\mu) > 25$, $p_{\rm T}(\tau_h) > 35$ GeV
- 1-prong τ_h
- $|\Delta \phi(e, \tau_{\rm h})| > 2.7$
- opposite sign μ and τ_h
- $m_{\rm T}(e, \tau_{\rm h}, E_{\rm T}^{\rm miss}) > 600 {\rm ~GeV}$



total SM = 1.4 ± 0.4 events

 $Z'(1000) = 5.5 \pm 0.7$

- Fake factor methods used to model multijet and *W*+jet backgrounds
- Need to be modeled in data-driven ways for two reasons:

1. populate the model in the high-mass tail

2. jet $\rightarrow \tau_h$ fake rate is mis-modeled in Monte Carlo.

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'11-'12
Double fake factor procedure



- The multijet contamination is estimated from the rate of non-isolated leptons, in both the signal region that passes tau ID, and the sample that fails.
- Then, the corrected number of tau candidates failing ID are weighted to predicted the W+jet background.
- This way, the corrections are small at each step.

Multijet background estimation



• Predict the number of multijet events:

$$N_{\text{multijet}}(p_{\text{T}}, \eta, x) = f_{\mu-\text{iso}}(p_{\text{T}}, \eta) \cdot N_{\text{multijet}}^{\text{fail } \mu-\text{iso}}(p_{\text{T}}, \eta, x)$$

$$= f_{\mu-\text{iso}}(p_{\text{T}},\eta) \cdot \left(N_{\text{data}}^{\text{fail }\mu-\text{iso}}(p_{\text{T}},\eta,x) - N_{\text{MC}}^{\text{fail }\mu-\text{iso}}(p_{\text{T}},\eta,x) \right)$$

W+jet background estimation

W+jet control region

- $m_{\rm T}(\mu, E_{\rm T}^{\rm miss}) = 70-200 \,\,{\rm GeV}$
- isolated lepton
- In a W+jet control region, divide tau candidates into pass and fail identification.
- Define fake factor:

$$f_{\tau}(p_{\mathrm{T}},\eta) \equiv \left. \frac{N^{\mathrm{pass}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta)}{N^{\mathrm{fail}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta)} \right|_{\mathrm{W-CR}}$$

Predict the number of W/Z+jet events:

$$N_{W/Z+jet}(p_{\mathrm{T}},\eta,x) = f_{\tau}(p_{\mathrm{T}},\eta) \cdot N_{W/Z+jet}^{\mathrm{fail}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta,x)$$

$$= f_{\tau}(p_{\mathrm{T}},\eta) \cdot \left(N_{\mathrm{data}}^{\mathrm{fail}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta,x) - N_{\mathrm{multijet}}^{\mathrm{fail}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta,x) - N_{\mathrm{MC}}^{\mathrm{fail}\ \tau-\mathrm{ID}}(p_{\mathrm{T}},\eta,x) \right)$$



Combined limit



ATLAS Z' SSM Exclusions: observed (expected) @ 95% CL

- τ_hτ_h: 1.26 (1.35) TeV
- *μτ*_h: 1.07 (1.06) TeV
- *eτ*_h: 1.10 (1.03) TeV
- *eμ*: 0.72 (0.82) TeV

combined: 1.40 (1.42) TeV

Published in PLB

Model dependence

- Z'_{SSM} has the same chiral couplings as the Z of the SM, but with a higher mass.
- The visible momentum fraction in hadronic tau decays can depend on the handedness of the couplings because it decays left-handed through a W.
- To probe the dependence on the limit, we tested two extreme cases:
 - 1. V-A pure left
 - 2. V+A pure right

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Model dependence





Conclusions

- It's been an exciting time to be a student.
- Had a lot of fun learning experimental skills during the turn-on of the LHC.
- Helped commission thresholds and hit efficiency for the TRT.
- Developed cut-based tau identification.
- Studied pile-up robustness for taus.
- Measured SM $Z \rightarrow \tau \tau$ cross section.
- Searched for new physics in high-mass ditau events.
- Excited that discoveries at the LHC could just be beginning.

I will always be indebted to Penn for enabling me and teaching me so much.



Back up



Review of Higgs search results



July 4, 2012 CERN announces the discovery of a new particle by ATLAS and CMS, consistent with the Higgs boson



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July 5 cover of the New York Times: *Physicists Find Elusive Particle Seen as Key to the Universe*

Current Higgs results

Branching ratios

 bb

qq

WW

- Two channels with precise mass measurements: $H \rightarrow \gamma \gamma$ and $H \rightarrow ZZ \rightarrow 4I$.
- $H \rightarrow WW$ observes a broad but clear excess.



Current Higgs results





- The measurements in the $\gamma \gamma/ZZ/WW$ channels are consistent with a new neutral boson with m \approx 126 GeV.
- Interestingly, both ATLAS and CMS observe the signal strength in the $\gamma\gamma$ channel to be higher than the SM over 1σ , but still consistent with the SM.
- $H \rightarrow \tau \tau$ and $H \rightarrow bb$ are approaching sensitivity.

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[1207.7214, ATLAS-CONF-2013-034] 48

Current $H \rightarrow \tau \tau$ result



- A lot of shared experience between $Z/Z'/H \rightarrow \tau \tau$ analyses.
- Uses similar $\sum \Delta \phi$ cut for suppressing W+jet.
- Uses fake factor method for predicting fake backgrounds.
- Eagerly approaching sensitivity to $1 \times SM H \rightarrow \tau \tau$.
- 21.7 fb⁻¹ collected this year.



The LHC, ATLAS, and CMS







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CMS



ATLAS

ER

Run: 154822, Event: 14321500 Date: 2010-05-10 02:07:22 CEST

 $M_{\mu\mu} = 87 \text{ GeV}$

Z→µµ candidate in 7 TeV collisions

Muons in ATLAS

- Combination of muon spectrometer segments with inner detector tracks.
- Track combination matching to reject decays in flight.
- Impact parameter constraints to reject cosmic muons.

Electrons in ATLAS

- Seeded by matching calorimeter clusters from a slidingwindow algorithm to inner detector tracks.
- Candidates are selected by: track quality, track-cluster matching, narrow calorimeter cluster, high electromagnetic fraction
- Tight candidates have cuts on E/p and high thresholds hits from the transition radiation in the TRT.

Transition Radiation Tracker



TRT Read-out





TRT hit efficiency



Parsing ATLAS raw data

2007

- Official libraries for parsing the ATLAS raw data were not finalized.
- Groups doing commissioning were still writing their own tools.
- My first software project was to write a library for parsing the ATLAS raw data for TRT commissioning purposes.
- Later, I used it to parse TRT threshold scans.



TRT Straw Hits







Event summary

	$ au_{ m had} au_{ m had}$	$ au_{\mu} au_{ m had}$	$ au_e au_{ m had}$	$ au_e au_\mu$
$m_{Z'} [{ m GeV}]$	1250	1000	1000	750
$m_{\rm T}^{\rm tot}$ threshold [GeV]	700	600	500	350
$Z/\gamma^* \to \tau\tau$	$0.73 {\pm} 0.23$	$0.36 {\pm} 0.06$	$0.57 {\pm} 0.11$	$0.55{\pm}0.07$
W+jets	< 0.03	$0.28 {\pm} 0.22$	0.8 ± 0.4	$0.33 {\pm} 0.10$
$Z(\rightarrow \ell \ell) + jets$	< 0.01	< 0.1	< 0.01	$0.06 {\pm} 0.02$
$t\overline{t}$	< 0.02	$0.33 {\pm} 0.15$	$0.13 {\pm} 0.09$	$0.97 {\pm} 0.22$
Diboson	< 0.01	$0.23 {\pm} 0.07$	$0.06 {\pm} 0.03$	$1.69 {\pm} 0.24$
Single top	< 0.01	$0.19 {\pm} 0.18$	< 0.1	< 0.1
Multijet	$0.24 {\pm} 0.15$	< 0.01	< 0.1	< 0.01
Total expected background	$0.97 {\pm} 0.27$	1.4 ± 0.4	1.6 ± 0.5	3.6 ± 0.4
Events observed	2	1	0	5
Expected signal events	6.3 ± 1.1	5.5 ± 0.7	5.0 ± 0.5	$6.72 {\pm} 0.26$
Signal efficiency (%)	4.3	1.1	1.0	0.4

Systematics

Uncertainty [%]	Signal			Ba	Background			
	hh	$\mu \mathrm{h}$	eh	$e\mu$	hh	$\mu \mathrm{h}$	eh	$e\mu$
Stat. uncertainty	1	2	2	3	5	20	23	7
Eff. and fake rate	16	10	8	1	12	16	4	3
Energy scale and res.	5	7	6	2	$+22 \\ -11$	3	8	5
Theory cross section	8	6	6	5	9	4	4	5
Luminosity	4	4	4	4	2	2	2	4
Data-driven methods	_	_	_	_	$+21 \\ -11$	6	16	

Table 2: Uncertainties on the estimated signal and total background contributions in percent for each channel. The following signal masses, chosen to be close to the region where the limits are set, are used: 1250 GeV for $\tau_{had}\tau_{had}$ (hh); 1000 GeV for $\tau_{\mu}\tau_{had}$ (μ h) and $\tau_e\tau_{had}$ (eh); and 750 GeV for $\tau_e\tau_{\mu}$ (e μ). A dash denotes that the uncertainty is not applicable. The statistical uncertainty corresponds to the uncertainty due to limited sample size in the MC and control regions.

Observed variance in fake-rates



BDTMedium)

- Divide the issue into two questions:
 - 1. Why do quarks and gluons have different tau fake-rates?

2. How does the quark/gluon fraction vary among samples?

Jet width for quark/gluons

Why do quarks and gluons have different tau fake-rates?

- $\Psi(r) =$ fraction of jet energy within $\Delta R < r$.
- Quark jets are more narrow than gluon jets of the same energy.
- Tau identification prefers narrow candidates.

 $\Psi(r=0.1)$ for 200 GeV Jets



• This is consistent with samples of quark-enriched jets, like *W*+jet, having higher fake-rates.

J. Gallicchio, M. Schwartz. "Quark and Gluon Tagging at the LHC". arXiv:1106.3076.

OS vs SS W+jet

How does the quark/gluon fraction vary among samples?

Leading order W+jet production:



- The charge of the quark should correlate with the reconstructed charge of the tau candidate, therefore (a) and (b) preferably produce opposite sign W+jet events.
- OS and SS will have different quark/gluon fractions.

Fake factors more quark-like at high- p_T



J. Gallicchio, M. Schwartz. "Pure Samples of Quark and Gluon Jets at the LHC". arXiv:1104.1175



Cross-check: single fake factor method



- Avoid issues of subtracting multijet from W/Z+jets background
- $H \rightarrow \tau \tau$ uses a fake factor method covering all tau fakes from W/Z+jets *and* multijet
- instead of estimating multijet independently with isolation fake factors Ryan Reece (Penn)

Cross-check: single fake factor method



 as expected, the single fake factor method overestimates in regions where the multijet contamination is large

Cross-check: single fake factor method



Lepton isolation fitting

- Used to cross-check the fake-factor QCD estimate.
- Similar in concept: predict normalization from the rate of non-isolated leptons.
- Fit calorimeter isolation with data-driven templates from QCD and W+jet control regions. Do fits separately in 6 p_T bins:



Agrees with the lepton isolation fake-factor method:

Fits work in	n failing tau	ID region	as well:
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fake-factor	fitting
6.3(6)	4.3(4)
2.7(3)	1.8(2)
0.29(8)	0.31(8)
0.05(3)	0.05(3)
	fake-factor 6.3(6) 2.7(3) 0.29(8) 0.05(3)


Madgraph predicted Quark/Gluon



J. Gallicchio, M. Schwartz. "Pure Samples of Quark and Gluon Jets at the LHC". arXiv:1104.1175

Control plots



Figure 28: (left) The distribution of the transverse mass of the combination of the selected electron and the $E_{\rm T}^{\rm miss}$, $m_{\rm T}(e, E_{\rm T}^{\rm miss})$. in events with exactly one selected electron, no additional preselected electrons or muons, and exactly one selected 1-prong tau. (right) The distribution of the electron impact parameter, d_0 , in events with exactly one selected electron, no additional preselected electron, no additional preselected electron impact parameter, d_0 , in events with exactly one selected electron, no additional preselected electrons or muons, and exactly one 1-prong tau candidate (without ID).

Electron veto

from tau WG Z \rightarrow ee tag-and-probe with 2.6/fb from 2011

	$ \eta < 1.37$	$1.37 < \eta < 1.52$	$1.52 < \eta < 2.0$	$2.0 < \eta $
BDT medium e-veto	1.64(0.81)	1.0(1.0)	0.71(0.63)	2.90(1.42)
BDT loose e-veto	1.21(0.30)	0.96(0.46)	0.59(0.21)	1.76(0.55)

- above are the scale factors and errors for the EleBDTMedium e-veto
- These SF have huge ≈100% uncertainties, and we've also previously shown that this veto is inefficient at high p_T.



Multijets background

- Loosened tau ID requirement in the anti-isolated lepton + tau data sample used to model the multijet background.
- Sample with tau ID is already multijet dominated.
 Loosening tau ID improves stats.
- Shapes are statistically consistent. Inclusive shape scaled to the prediction with medium tau ID to give the estimate in the tail.





- Categorize Monte Carlo events by electron or jet faking tau.
- Loosen electron veto in Monte Carlo sample matched to electron fakes.
- Shapes are consistent, and only driven by the $Z \rightarrow ee$ kinematics.
- $Z \rightarrow ee$ with e-faking tau is negligible
- $Z \rightarrow ee + jet-fake$ covered with the data-driven W/Z+jet tau fake factor method.

High-p_T tau efficiency systematic

- The dominant systematic uncertainty for the Z' signal and the $Z \rightarrow \tau \tau$ background.
- Low p_T uncertainty of 4% taken from the Tau WG blessed $Z \rightarrow \tau \tau$ tag-and-probe.
- No high- p_T control sample of true taus.
- Assume mis-modeling comes from either:
 - 1. tau kinematics (TAUOLA)
 - 2. detector response to high- p_T pions \leftarrow dominant
- Instead of using true taus, measure the trend in the scale factor for fakes from dijet events.
 - $p_{\rm T} \le 100 \text{ GeV}$: $\Delta \varepsilon = 4\%$ (taken from the $Z \to \tau \tau$ measurement)
 - $p_{\rm T} > 100 \text{ GeV}$: $\Delta \varepsilon = 4 + 0.016(p_{\rm T} 100)\%$, with $p_{\rm T}$ in GeV (taken from the dijets measurement).



$Z' \rightarrow \tau \tau \rightarrow \tau_h \tau_h$



- Tau ID efficiency uncert. $\approx 11\%$ on the signal. (4% from $Z \rightarrow \tau \tau$ tag-and-probe)
- Jet/tau energy scale uncert. $\approx +22/-11\%$
- Multijet modeled by fitting the shape of the SS data. uncert. \approx +21/-11%

Data-driven multijet

Fit same-sign (SS) data with dijet function:



 $f(M_{\rm T}|p_0, p_1, p_2) = p_0 \cdot M_{\rm T}^{p_1 + p_2 \log M_{\rm T}}.$

• normalize in OS sideband with $200 < M_T < 250$ GeV

ATLAS Tau ID





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Phenomenology of tau decays

$ au^- ightarrow$	$e^- \bar{ u}_e u_{ au}$	17.8%	$\Big]$
	$\mu^- ar{ u}_\mu u_ au$	17.4%	ieptonic 55.270
	$\pi^- \pi^0 u_ au$	25.5%	
	$\pi^- \nu_{ au}$	10.9%	
	$\pi^- 2\pi^0 \nu_{ au}$	9.3%	> 1 prong 49.5%
	$K^{-} \left(N \pi^{0} ight) \left(N K^{0} ight) u_{ au}$	1.5%	
	$\pi^- 3\pi^0 \nu_{ au}$	1.0%	J
	$\pi^- \pi^- \pi^+ u_ au$	9.0%	$\frac{1}{2}$ propg 15.9%
	$\pi^ \pi^ \pi^+$ π^0 $ u_ au$	4.6%	$\int 3 \text{ prolig } 13.270$



Current tau identification variables



8. Maximum ΔR between jet-axis and core tracks

*has pile-up correction term linear in N(vertex) Ryan Reece (Penn)

Seeds of reconstruction

Once upon a time, there were two tau reconstruction algorithms.

- 1. tauRec seeded by $p_{\rm T} > 10$ GeV anti- $k_{\rm T}$ 0.4 topo-jets. "calo-seeded"
- 2. tau1p3p seeded by $p_{\rm T} > 6$ GeV inner detector tracks. "track-seeded"



Since virtually all candidates have a calo-seed, we effectively merged the variable calculation of both algorithms, using only calo-seeds.

"Performance of the tau reconstruction and identification algorithm with 14.2.20 and mc08" [ATL-COM-PHYS-2009-229] Ryan Reece (Penn)

2009

Early MV identification

• Jet-tau discrimination



Prefers narrow calorimeter jets, likelihood-based discriminant.



• Electron-tau discrimination

	Isl	Ele(%)		IsEl	.e_eg(%	()
Candidate	Overall	1P	3P	Overall	1P	3P
$ au$ from W $\rightarrow au v$	93.2	92.7	95.3	99.8	99.8	99.8
$ au$ form $A \rightarrow au au$	93.3	92.5	96.3	99.9	98.8	99.5
Electron form $W \rightarrow ev$	2.8	2.4	0.1	14.8	13.4	0.3
Electron form $A \rightarrow \tau \tau$	5.9	4.5	0.5	18.0	15.8	0.8

[ATL-COM-PHYS-2009-229]

Early sub-structure studies



- Monte Carlo based substructure studies
- Cell-based shower-shape subtraction π^0 reconstruction.
- Still unvalidated with data.

First data



- First comparisons of background distributions and the QCD fake-rate between data and Monte Carlo.
- Already see that MC over-estimates the jet fake-rate. $\Rightarrow k_W \approx 0.5$ "Reconstruction of hadronic tau candidates in QCD events at ATLAS with 7 TeV pp collisions" [ATLAS-CONF-2010-059] "Tau Reconstruction and Identification Performance in ATLAS" [ATLAS-CONF-2010-086] 88

Tau discriminants

b

Х

• Cuts

 p_{T} -parametrized cuts on R_{EM} and R_{track} , and a cut on f_{track} .

Projective likelihood

$$d = \ln\left(\frac{L_S}{L_B}\right) = \sum_{i=1}^N \ln\left(\frac{p_i^S(x_i)}{p_i^B(x_i)}\right)$$



В

S



Maturing of discriminants



- Cuts are pt-parametrized to account for the Lorentz collimation of boosted taus.
- Experience grows with LLH and BDT discriminants, which become the preferred discriminants in 2011.

"Reconstruction, Energy Calibration, and Identification of Hadronically Decaying Tau Leptons in the ATLAS Experiment" [ATLAS-CONF-2011-077, ATL-PHYS-INT-2011-068]

Seeing first hadronic taus

2010



- Nov 2010: Observation of $W \rightarrow \tau_h \nu$ [ATLAS-CONF-2010-097]
- Feb 2011: Observation of $Z \rightarrow \tau_h \tau_l$ [ATLAS-CONF-2011-010]

$W \rightarrow \tau \nu$ cross section

$$\sigma(W \to \tau \nu) = 11.1 \pm 0.3 (\text{stat.}) \pm 1.7 (\text{sys.}) \pm 0.4 (\text{lumi.}) \text{ nb}$$

2010

 $\sigma_{\rm theory} = 10.46 \pm 0.52 \text{ nb}$ at NNLO



Dominant systematics

- $\tau_{\rm h}$ efficiency 10.3%
- $\tau_{\rm h}$ energy scale 8.0%
- $\tau_{\rm h}$ + MET trigger efficiency 7.0%
- Iuminosity 3.4%
- acceptance 2.3%

"Measurement of the $W \rightarrow \tau \nu$ cross section in pp collisions at sqrt(s)= 7 TeV with the ATLAS experiment" [arXiv:1108.4101]

$Z \rightarrow \tau \tau$ cross section

2011

 $\sigma_{\rm combined} = 0.97 \pm 0.07 ({\rm stat.}) \pm 0.07 ({\rm sys.}) \pm 0.03 ({\rm lumi.})$ nb

 $\sigma_{\rm theory} = 0.96 \pm 0.05 \ {\rm nb}$ at NNLO



Dominant systematics

- $\tau_{\rm h}$ energy scale 11%
- $\tau_{\rm h}$ efficiency 8.6%
- μ efficiency 8.6%
- e efficiency 3-10%
- acceptance 3%
- luminosity 3.4%

"Measurement of the $Z \rightarrow \tau \tau$ cross section in pp collisions at sqrt(s)= 7 TeV with the ATLAS detector" [arXiv:1108.2016]

Tau vertex association

Tau track selection

- $p_{\rm T} > 1 \text{ GeV}$,
- Number of pixel hits ≥ 2 ,
- Number of pixel hits + number of SCT hits ≥ 7 ,
- $|d_0| < 1.0 \text{ mm},$
- $|z_0 \sin \theta| < 1.5 \text{ mm},$

- The d_0 and z_0 requirements depend on the choice of vertex.
- Beginning in 2012, choose the vertex with the highest JVF for that tau candidate.



Track selection efficiency



- In 2011, the track selection for tau candidates cut on the d_0 and z_0 with respect to the vertex with the highest $\sum p_T^2$.
- Selecting the vertex with the highest JVF recovers efficiency in high pile-up (Tau Jet Vertex Association).

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[ATLAS-CONF-2012-142] 95

Electron veto

- Electrons provide a track and calorimeter deposit that can fake hadronic tau decay identification.
- ATLAS provides a BDT to discriminate electrons from tau candidates, even after removing overlaps with selected electrons.
- Tight/Medium/Loose working points are defined (≈75%, 85%, 95% efficient).
- In 2012, the BDT is being reoptimized to have better efficiency at high-p_T.



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[ATLAS-CONF-2011-152, ATLAS-CONF-2012-142] 9

Tau triggering

1. Level 1: (latency 2.5 μs)

Coarse EM+Had calorimeter trigger towers $\Delta \eta \times \Delta = 0.1 \times 0.1$. Candidate passing thresholds on the sum of energies:

- 1. highest 2×1 towers
- 2. surrounding 4×4 isolation ring
- 2. Level 2: (latency 40 ms)

Fast tracking. Region-of-interest (RoI) calculation of track- and calorimeter-based ID variables. Similar selection to offline cut-based ID.

3. Event Filter: (latency 4 s)

Beginning in 2012, started using the offline BDT algorithm at the EF trigger.



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https://twiki.cern.ch/twiki/bin/view/AtlasPublic/TauTriggerPublicResults

L2 pile-up robustness

Example improvements to variable definitions to lessen sensitivity to pile-up:

- Smaller ΔR cone for calculating EM radius 0.4 \rightarrow 0.2
- Select tracks within $\Delta z < 2 \text{ mm}$ of the highest- p_T track within the RoI (cannot vertex at L2).



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https://twiki.cern.ch/twiki/bin/view/AtlasPublic/TauTriggerPublicResults

VBF triggers

- New VBF triggers *relax tau identification* required at L2 and the EF by adding requirements for forward jets.
- This increases the control sample of tau candidates that will fail identification, used to estimate the fake contribution.
- Being evaluated for the $H \rightarrow \tau \tau \rightarrow lep + \tau_{had}$ search.



₩/Z (iet 2 L2 jets $p_T > 15$ GeV, $|\Delta \eta| > 2.5$ 2 EF jets $p_{\rm T}$ > 25 GeV, $|\Delta \eta|$ > 2.8,

Trigger menu

 $M_{\rm ii} > 400 \,\,{\rm GeV}$



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q'(jet)

Identification efficiency

- **Tag-and-probe:** selecting a sample of a known composition without some ID, so one can probe its efficiency.
- For the case of tau ID, select $Z \rightarrow \tau \tau \rightarrow \mu \tau_h 3\nu$ by triggering on the muon and selecting events with muon + tau candidate.



Scale factor ≈ 1, known to a few percent, 2-3% (1-prong),
 5-6% multi-prong.

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[ATLAS-CONF-2012-142] 100

Trigger efficiency

- The same $Z \rightarrow \tau \tau \rightarrow \mu \tau_h 3\nu$ tag-and-probe sample is used to measure the efficiency of the tau triggers.
- Known to O(5%) in the turn-on.
- Efficiency Improving with EF_tau20_medium1 09 0.8 statistics in 2012. 0.7 0.6 0.5 **ATLAS** Preliminary 0.4 $\int dt L = 1.0 \text{ fb}^{-1}, \sqrt{s} = 7 \text{ TeV}$ 0.3 🗕 Data 0.2E 7→ττ 0.1E $\epsilon_{\text{Data}}/\epsilon_{\text{MC}}$ 1.2 0.8 70 0 10 20 30 40 50 60 80 90 100 $p_{\tau}(\tau)$ [GeV]

Electron veto fake rate



 Probe the *e*-veto efficiency after removing overlap with selected electrons.

data/MC	scale	facto	r and	uncerta	ainty
from $Z \rightarrow ee$	tag-and	-probe v	with 2.0	5/fb from	2011

electron BDT veto	$ \eta_{\rm trk} < 1.37$	$1.37 < \eta_{\rm trk} < 1.52$	$1.52 < \eta_{\rm trk} < 2.00$	$ \eta_{\rm trk} > 2.00$
loose	0.96 ± 0.22	0.8±0.3	0.47 ± 0.14	1.7 ± 0.4
medium	1.3 ±0.5	-	0.5 ± 0.4	2.8 ± 1.3

- Statistically limited by the sample that pass the veto, giving uncertainties \approx 50-100%.
- Improving with the data added in 2012. Ryan Reece (Penn)

[ATLAS-CONF-2012-142] 102

Energy scale



- Tau candidates are first brought from the EM to the Jet Energy Scale with LC calibration of the clusters within $\Delta R < 0.2$ (from 0.4 to be pile-up robust).
- Then response functions are calibrated with tau Monte Carlo to make final corrections of a few percent.
- Uncertainties are determined by smearing the Monte Carlo truth according the tau decays true composition, using uncertainties constrained by single particle response measurements (CTB, E/p, $Z \rightarrow ee/\pi^0$ -resp.) Ryan Reece (Penn) [ATLAS-CONF-2012-054] 103

Energy scale cross check

- Tau energy scale is manually shifted in the modeling.
- Median of the visible mass peak is used to decide which scale matches the data.
- Toy experiments are used to estimate the uncertainty.

$ \eta $	best scale	uncert.
0.0-0.8	-1.5%	3.3%
0.8-2.5	+1.5%	2.8%

- Scale consistent with 1 within singleparticle-response uncertainties $\approx 3\%$.
- May become primary method with more data.







[ATLAS-CONF-2012-054] 104

CMS Tau ID



Hadronic decays dominantly to 1 or 3



• τ_h reco seeded by calorimeter jets

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- associate tracks in $\Delta R < 0.2$, select 1 or 3
- combine calorimeter and tracking information in a BDT or likelihood discriminant, preferring narrow clustering, hadronic activity
- particle-flow reconstructs constituent 4-vectors
- τ_h reco seeded by particle-flow hadrons
- Hadron Plus Strip (HPS) algorithm for counting $\pi^0 s$
- isolation cone for rejecting QCD jets

[ATLAS-CONF-2011-152, CMS PAS TAU-11-001] 106

CMS Particle Flow



- Matches track to clusters to form charged and neutral PF objects.
- PF objects are used as input for all CMS tau reconstruction.

CMS: Hadron Plus Strip (HPS)



```
Build all possible taus
that have a 'tau-like'multiplicity
from the seed jet
\pi^+
\pi^+ \pi^0
\pi^+ \pi^+ \pi^-
```

tau that is 'most isolated' with compatible m_{vis} is the final tau candidate associated to the seed jet

Discrimination with calorimeter based isolation $\Delta R < 0.5$.

[CMS PAS TAU-11-001]
CMS: Tau Neural Classifier (TaNC)

- Uses a *shrinking core-cone*:
 - $\Delta R(\text{photons}) < 0.15$ for photons
 - $\Delta R(\text{charged}) < (5 \text{ GeV})/E_{T}$ for charged hadrons
 - $\Delta R(\text{charged}) < \Delta R(\text{isolation}) < 0.5$



Decay mode	Resonance	Mass (MeV/ c^2)	Branching fraction (%)
$ au^- o h^- u_ au$			11.6%
$ au^- o h^- \pi^0 u_ au$	$ ho^-$	770	26.0%
$ au^- o h^- \pi^0 \pi^0 u_ au$	a_1^-	1200	9.5%
$ au^- o h^- h^+ h^- {m v}_ au$	a_1^-	1200	9.8%
$ au^- o h^- h^+ h^- \pi^0 u_ au$	-		4.8%

• Dedicated Neural-net classifier for each decay mode [CMS PAS TAU-11-001]

CMS Performance

 Not trivial to compare ATLAS and CMS tau
 performance because
 we bin fake-rates in
 N(track) instead of
 categorizing the
 decay mode.



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CMS decay mode ID



[CMS PAS TAU-11-001] 111

Calorimeter granularity

ATLAS

- *B* = 2.0 T
- $\Delta \eta \times \Delta \phi =$ 0.025×0.0245
- R = 0.4 anti- k_T topo-jets

CMS

- *B* = 3.8 T
- $\Delta\eta \times \Delta\phi = 0.0174 \times 0.0174$
- R = 0.5 anti- k_T PF-jets



ATLAS Barrel EM Calorimeter

Granularity could fundamentally limit our capacity to reconstruct sub-structure / π^0 s.

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