Searching for new physics in high-mass ditau events at ATLAS

Ryan Reece

University of Pennsylvania

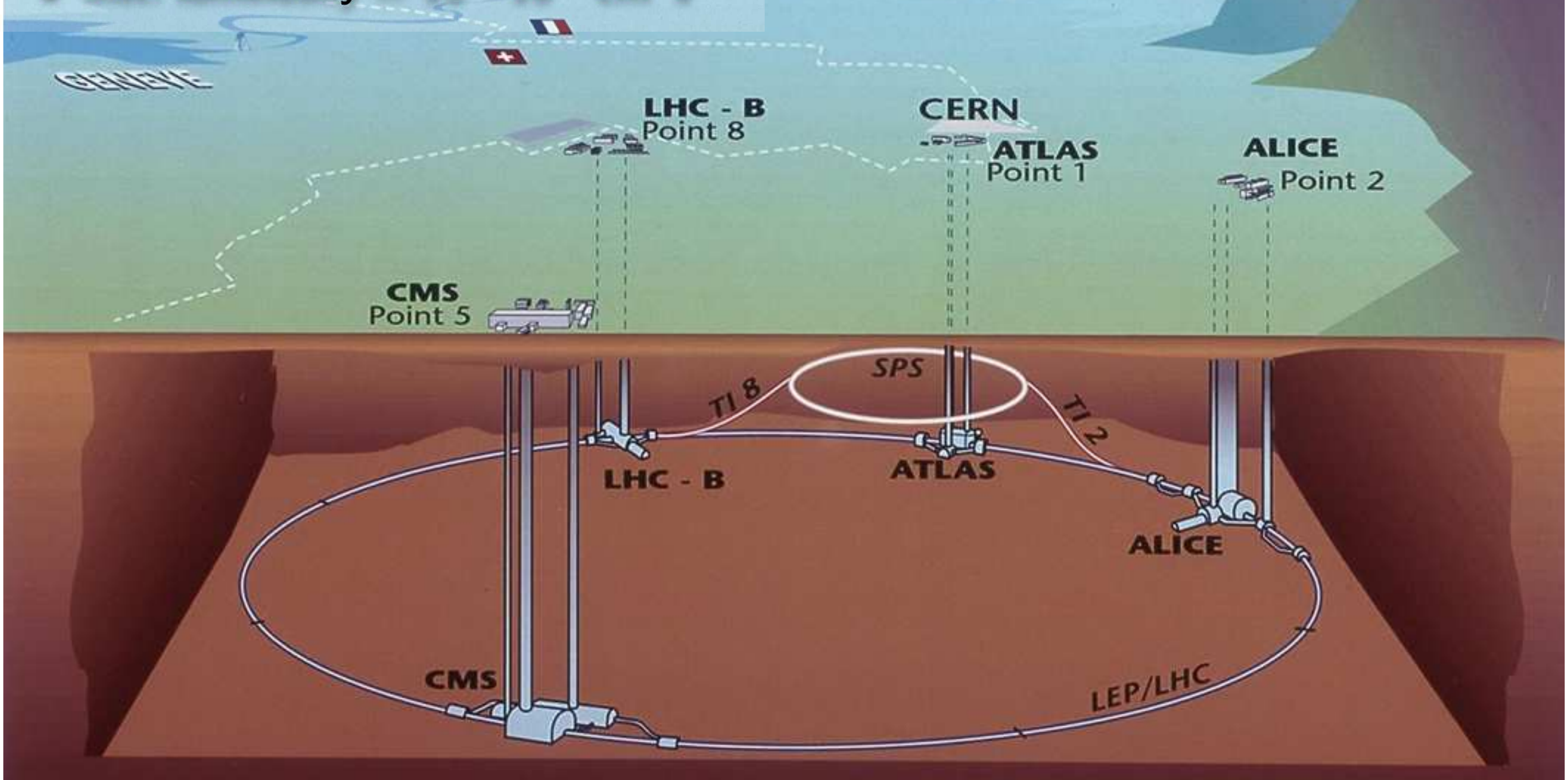
ryan.reece@cern.ch

Outline

1. Quick intro to the LHC and ATLAS
2. Tau performance
   reconstruction, cut-based ID, pile-up robustness
3. SM $Z \to \tau\tau$
   selection design, observation, cross section
4. $Z' \to \tau\tau$
   motivation, fake factor method, exclusion
5. Review of SM Higgs results

Overall view of the LHC experiments.

- 27 km circumference
- 1232 dipoles: 15 m , 8.3 T
- 96 tonnes liquid He, 1.9 K
- p-p collisions at $\sqrt{s} = 7$-8 TeV
- inst. luminosity = 10^{32}-10^{34} cm^{-2}s^{-1}
- 10^{11} protons / bunch $\times$ 1000 bunches
- 20 MHz , 50 ns bunch spacing
- 1-40 interactions / crossing
- 0.5×10^{9} interactions / sec



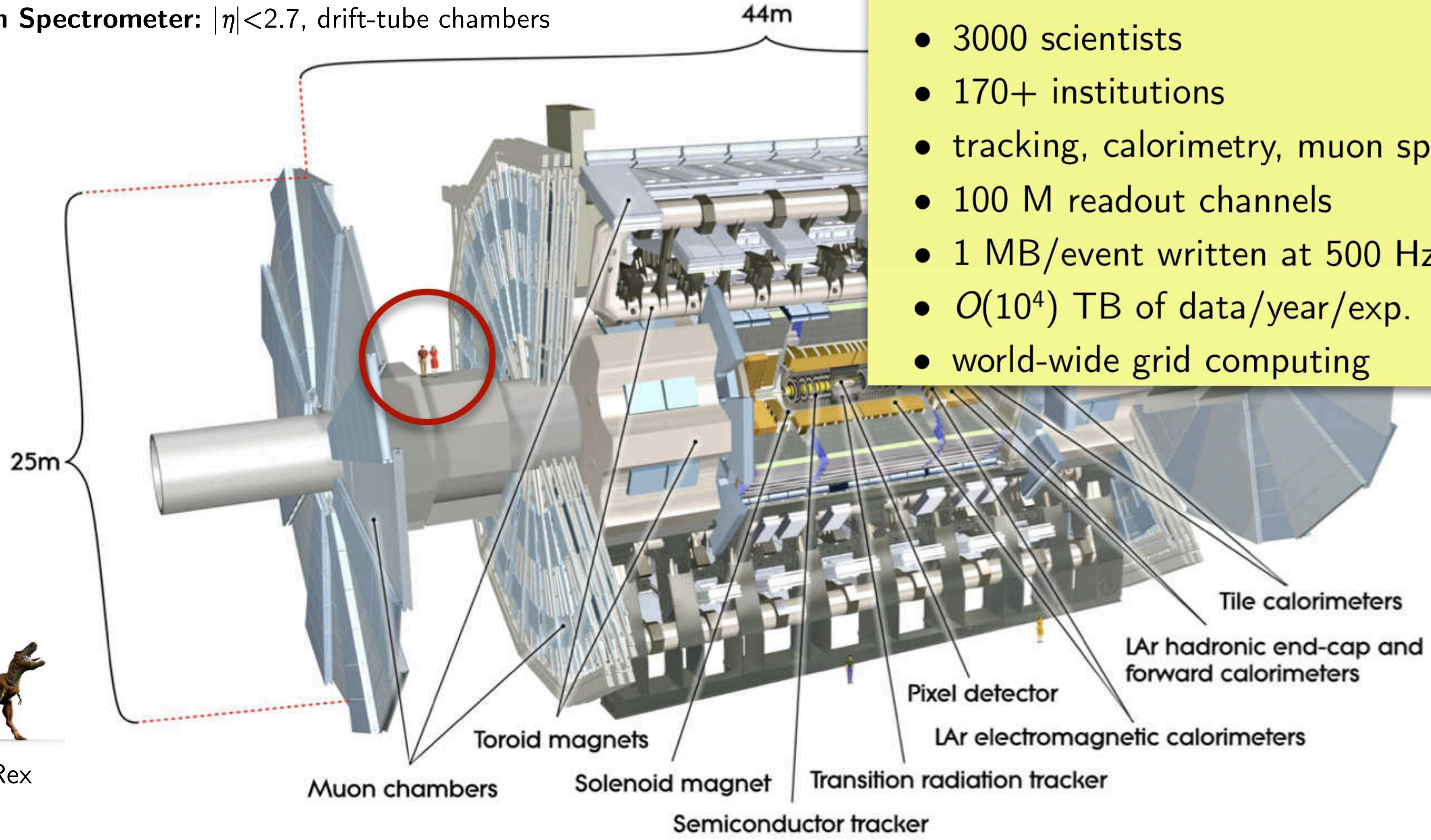
ATLAS Detector

Magnets: 5 tonne central solenoid, 2T in inner detector, 4T toroid system

Muon Spectrometer: $|\eta|<2.7$, drift-tube chambers

Both ATLAS and CMS have:

- 3000 scientists
- 170+ institutions
- tracking, calorimetry, muon spec.
- 100 M readout channels
- 1 MB/event written at 500 Hz
- $O(10^4)$ TB of data/year/exp.
- world-wide grid computing



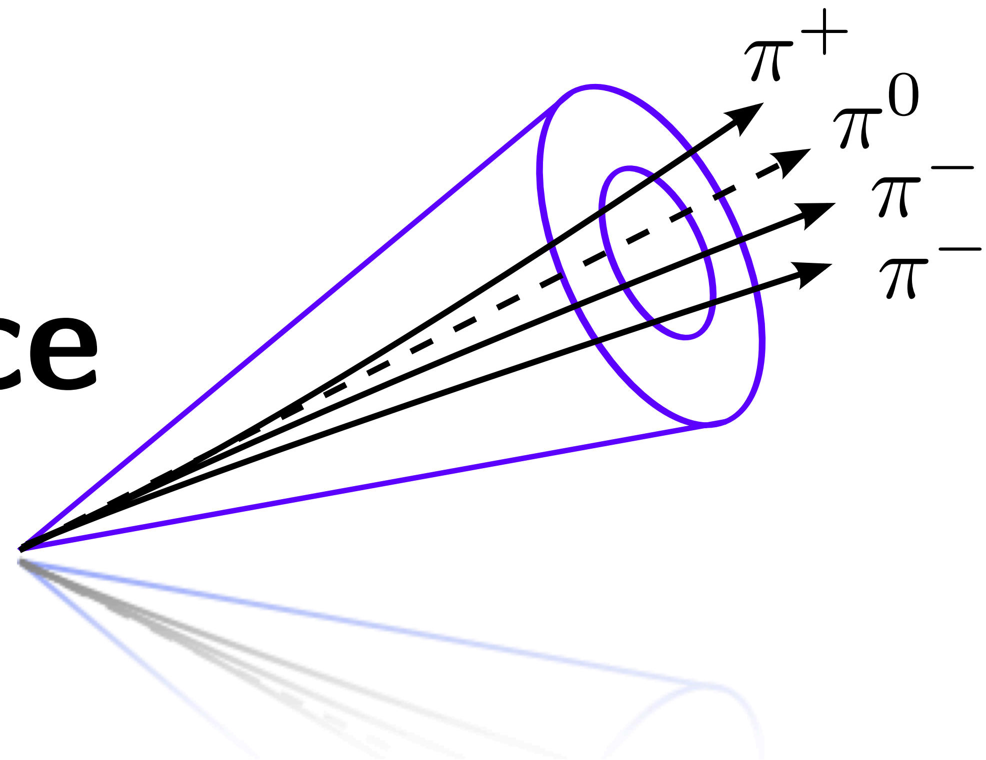
T.Rex

Tracking: $|\eta|<2.5$, B=2T, precise tracking and vertexing, Si pixels, strips, and TRT straws, TR electron ID

Electromagnetic Calorimeter: $|\eta|<3.2$, 3+1 layers corrugated layers of lead and LAr

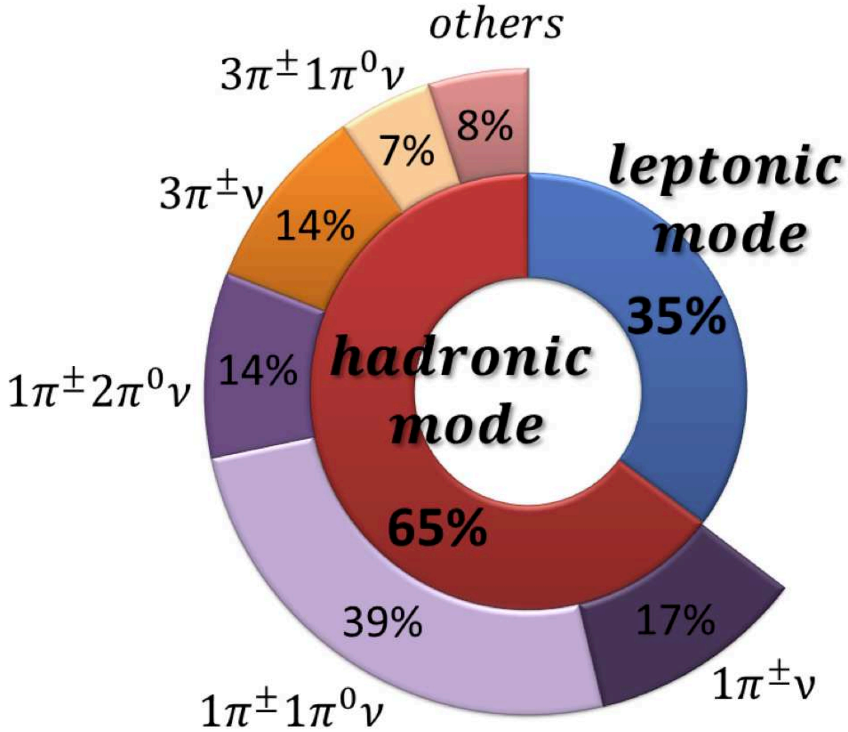
Hadronic Calorimeter: $|\eta|<5$, Central: iron/scintillator tiles, Forward: copper/tungsten-LAr

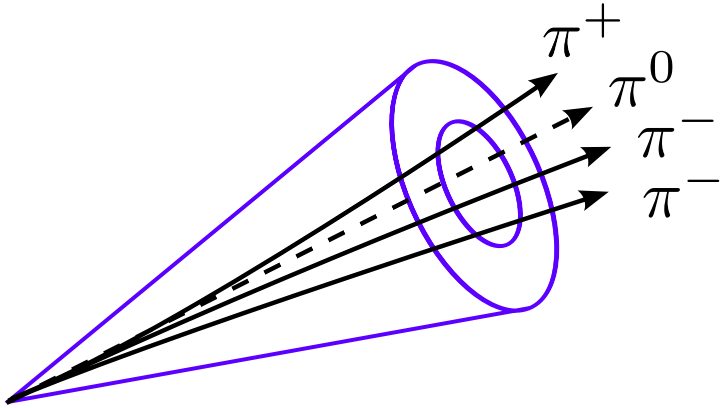
Tau performance

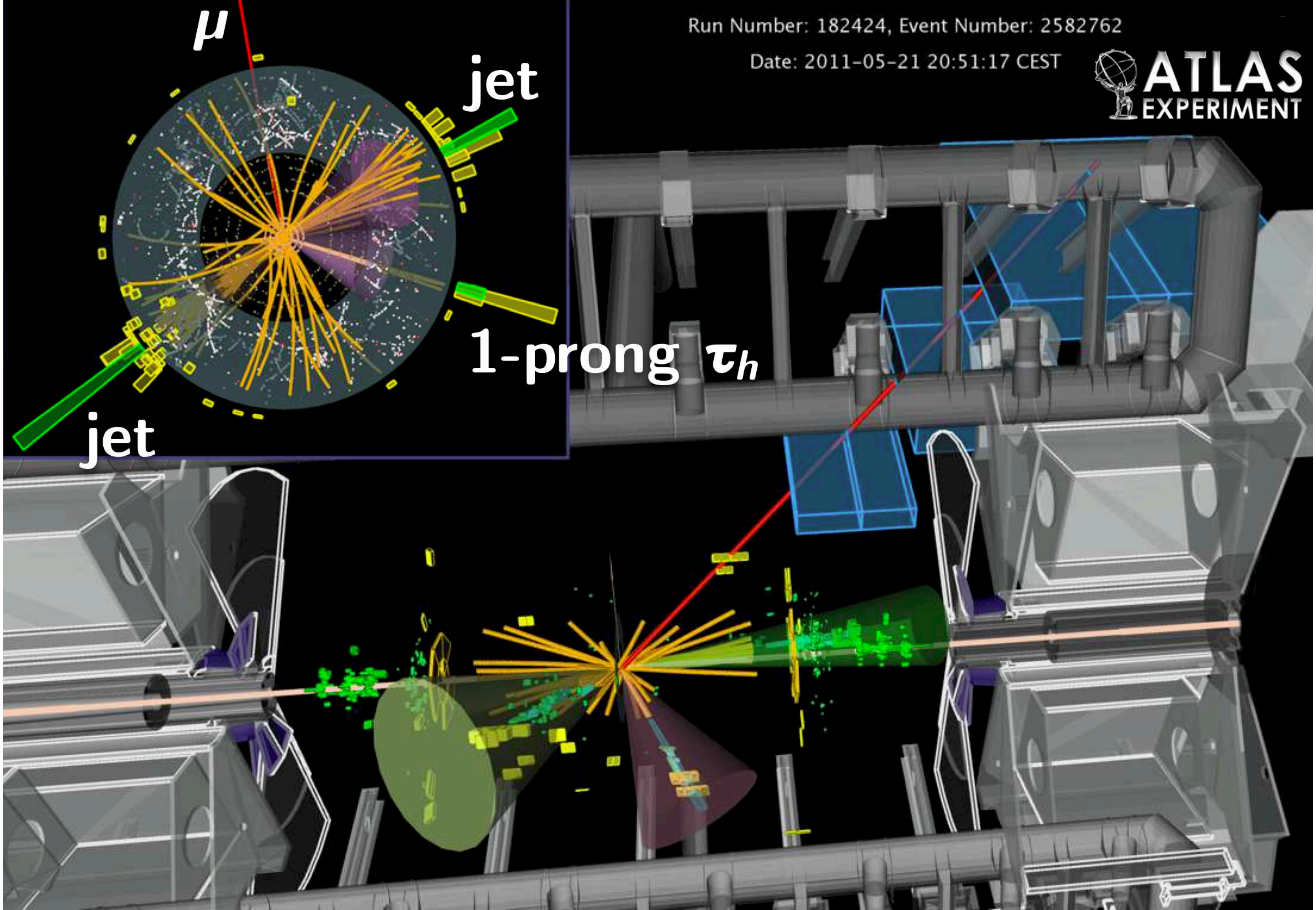


What's a tau?

- Only lepton massive enough to decay hadronically.
- Decay in beam pipe: $c\tau \approx 87$ μm
- 65% hadronic
  50% 1-prong, 15% 3-prong.
- **Signature:** narrow jet with 1 or 3 tracks, possibly additional EM clusters.
- **Challenge:** large multijet background at hadron colliders.
- **Importance:** often preferred coupling to new physics:
  SM $H \to \tau\tau$, $H^+ \to \tau^+\nu$, $Z' \to \tau\tau$,
  high-tanβ SUSY...



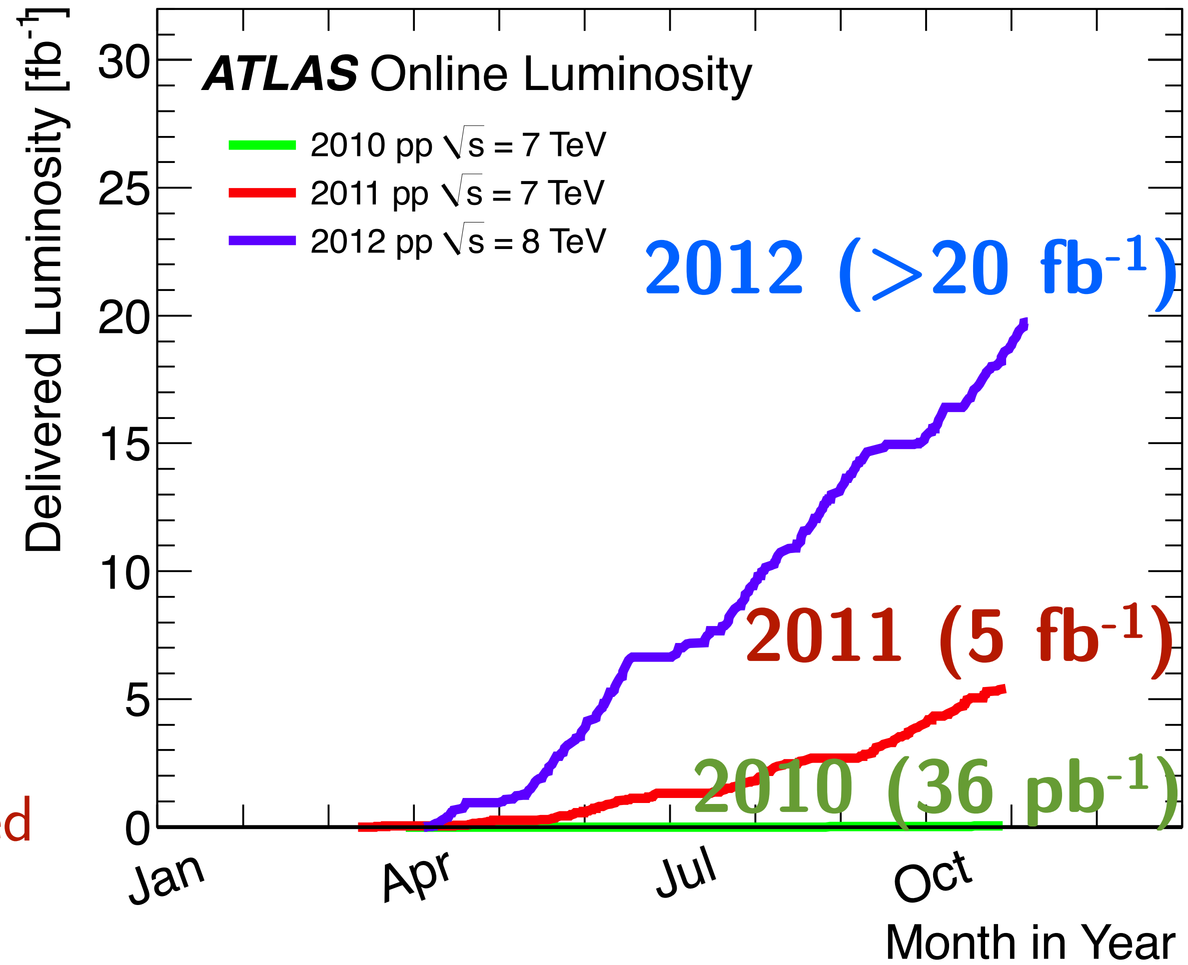




$t\bar{t} \to b\bar{b}(\mu\nu)(\tau_{\mathrm{h}}\nu)$ candidate

Timeline of taus at ATLAS

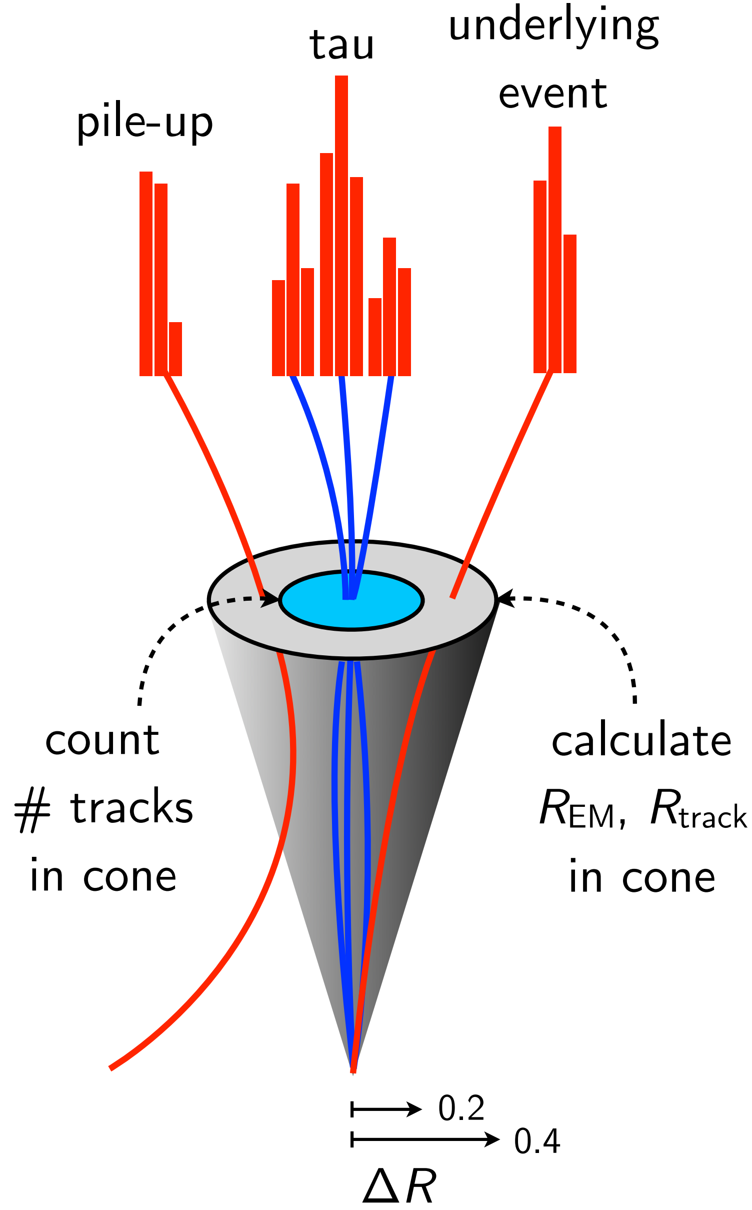
- Nov 2010: Obs. of $W\to\tau\nu$ (546 nb^{-1})
- Feb 2011: Obs. of $Z\to\tau\tau$ (8.5 pb^{-1})
- July 2011: $W\to\tau\nu$ and $Z\to\tau\tau$ cross section measurements (36 pb^{-1})
- Feb 2012: $Z\to\tau\tau$ cross section (1.5 fb^{-1})
- June 2012: SM $H\to\tau\tau$ excluded 3-4$\times$SM at $m_H\approx 125$ GeV [arXiv:1206.5971]
- Several other analyses: MSSM $H\to\tau\tau$, $t\bar{t}$ with τ, $H^+\to\tau\nu$, $Z'\to\tau\tau$, SUSY τ+MET, ...



- Now eagerly waiting to see if $H\to\tau\tau$ will be excluded at 1$\times$SM this year?

Tau reconstruction

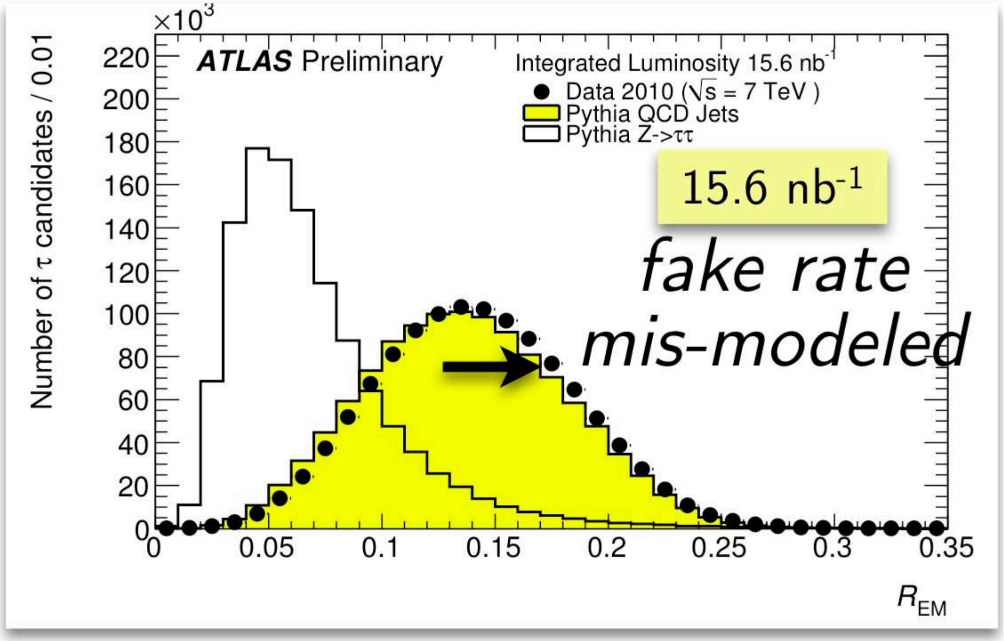
1. **Seeded by anti-k_t jets** (R=0.4) of 3-D topological calorimeter clusters.
2. **Define the four-momentum** as the jet-axis with a tau-specific calibration.
3. **Associate tracks** with the jet that are consistent with the chosen vertex.
4. **Calculate discriminating variables** from the combined calorimeter and tracking information, later used to identify hadronic tau decays with BDT and likelihood based discriminants.

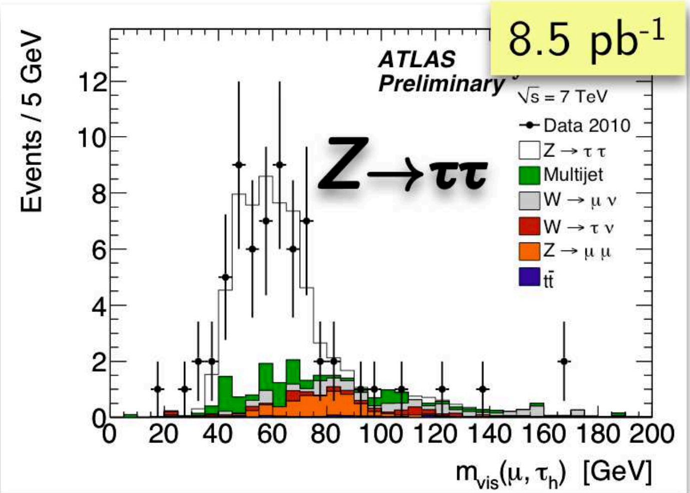


Cut-based tau ID

2010

- My timing with the development of tau reconstruction software and the arrival of first collision data allowed me to contribute to the ***first data/MC comparisons*** of tau ID variables and ***develop the first cut-based ID*** used with ATLAS data.
- Prefers narrow calorimeter deposits and closely associated tracks.
- Used in the ***first observation*** of $W \to \tau\nu$ and $Z \to \tau\tau$.





Safe Cut Variables

1. $R_{\mathrm{EM}} = \frac{\sum \Delta R\, E_{\mathrm{T}}}{\sum E_{\mathrm{T}}}$, summed over cells from the first 3 layers of the EM calorimeter within $\Delta R < 0.4$.
2. $R_{\mathrm{track}} = \frac{\sum \Delta R\, p_{\mathrm{T}}}{\sum p_{\mathrm{T}}}$, summed over tracks associated to the tau within $\Delta R < 0.2$.
3. $f_{\mathrm{trk},1} = \frac{p_{\mathrm{T}}(\mathrm{lead\ track})}{p_{\mathrm{T}}(\tau_{\mathrm{h}})}$, the transverse momentum fraction of the leading track.

E_T-parametrized ID

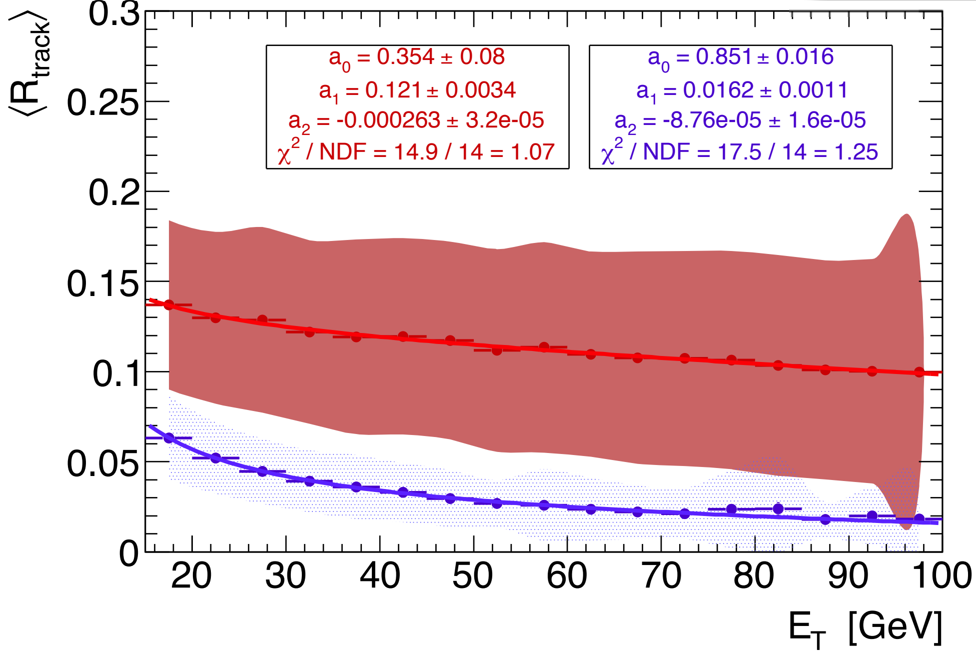
2010

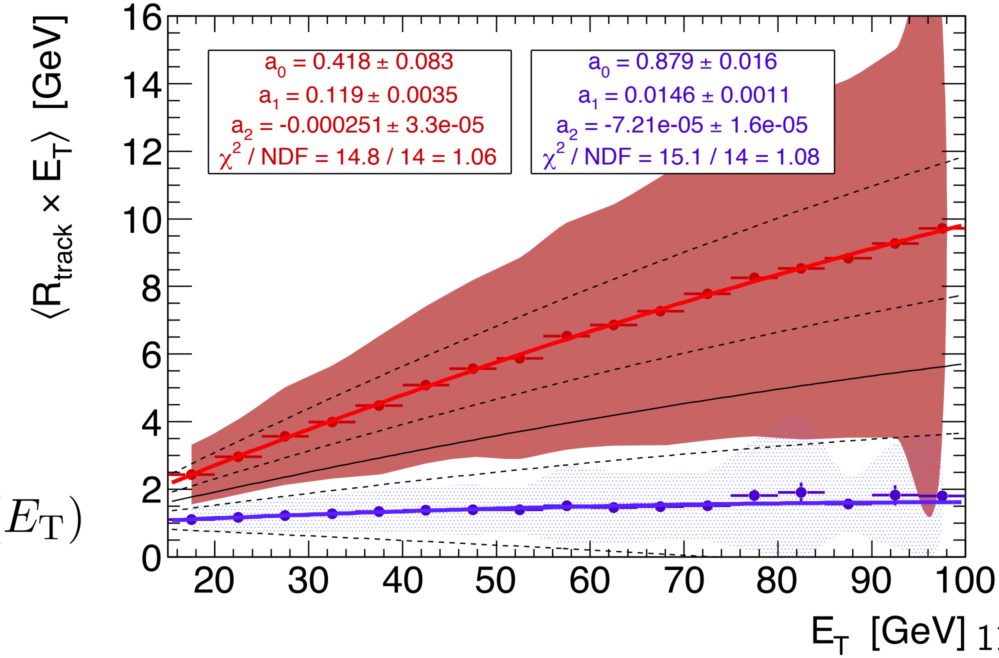
- Lorentz boost implies R should collimate as

$$R(E_{\mathrm{T}}) \propto \frac{1}{E_{\mathrm{T}}}$$

- Multiplying by E_{T} flattens out E_{T}-dependence.
- Example plots here are for 3-prong, but the π^0s of 1-prong taus also collimate.
- Construct a smooth family of curves between the signal and background that have efficiency that is approximately flat in E_{T}.

$$R_{\mathrm{cut}}(E_{\mathrm{T}}; x_{\mathrm{cut}})\, E_{\mathrm{T}} = (1 - x_{\mathrm{cut}})\, f_{\mathrm{sig}}(E_{\mathrm{T}}) + \; x_{\mathrm{cut}}\, f_{\mathrm{bkg}}(E_{\mathrm{T}})$$



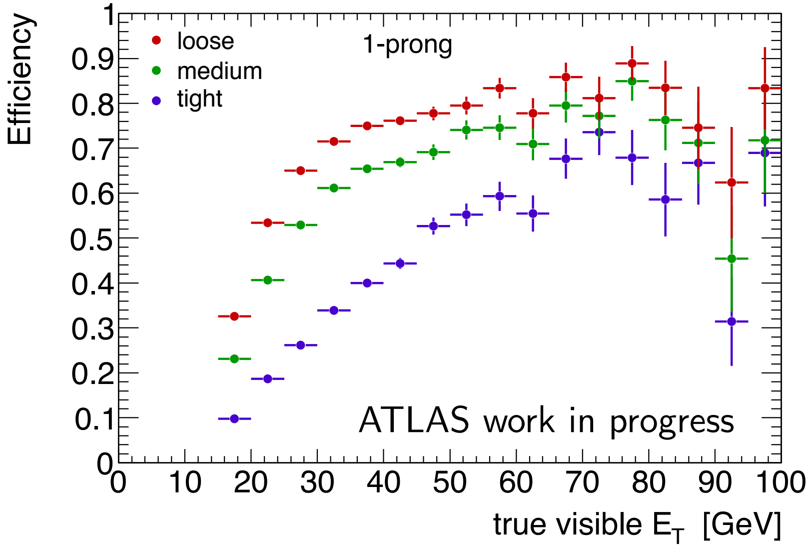


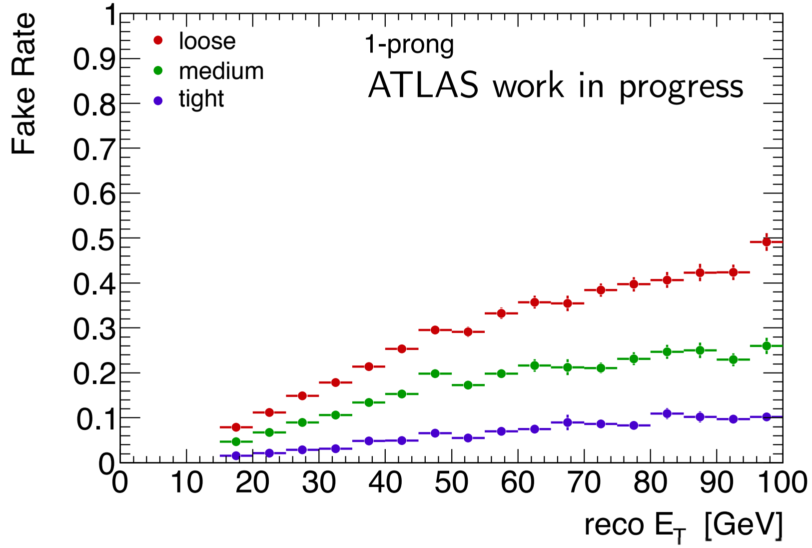
Efficiency / Rejection

2010

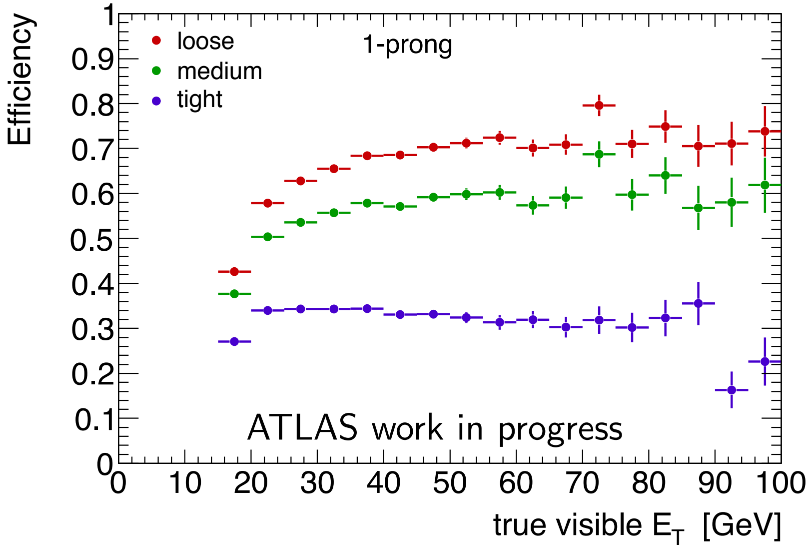
Simple cuts

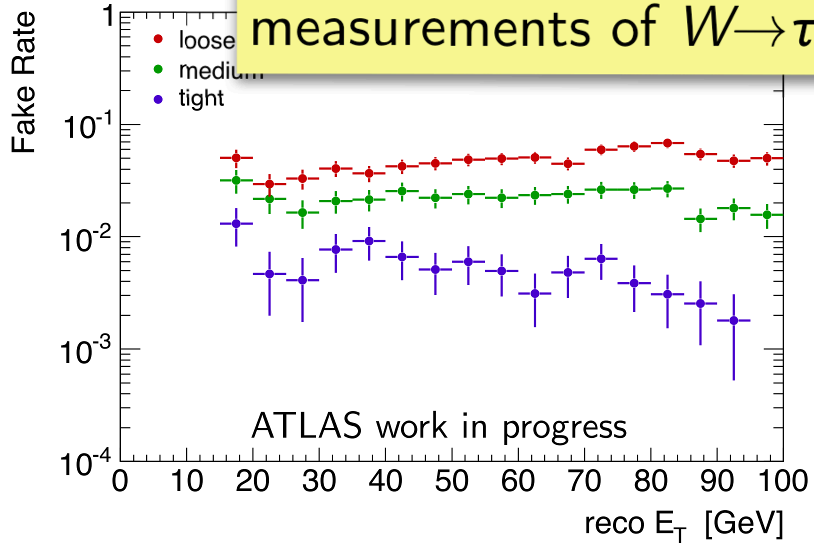
1-prong





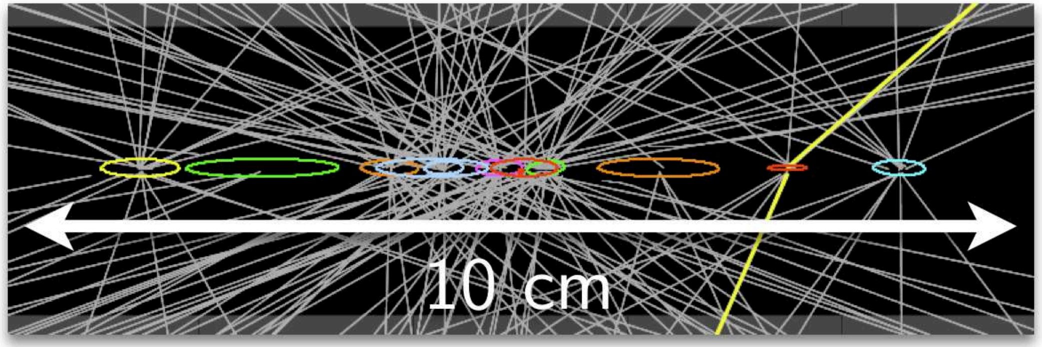
E_{T}-parametrized cuts

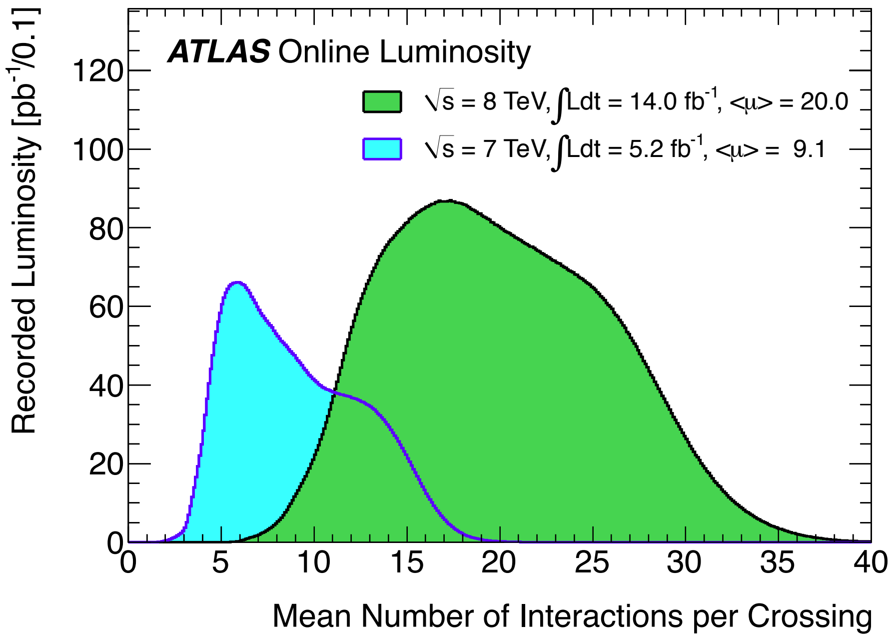




used for first ATLAS cross section measurements of $W \to \tau\nu$ and $Z \to \tau\tau$

Pile-up





- 1-40 pile-up interactions / crossing
- The additional tracks and clusters from pile-up are especially challenging for tau identification, which discriminates hadronic tau decays from jets with isolation-related track and calorimeter quantities.
- Efforts in 2011-2012 involved re-defining or adding corrections to identification variables to be more robust against the increasing pile-up.

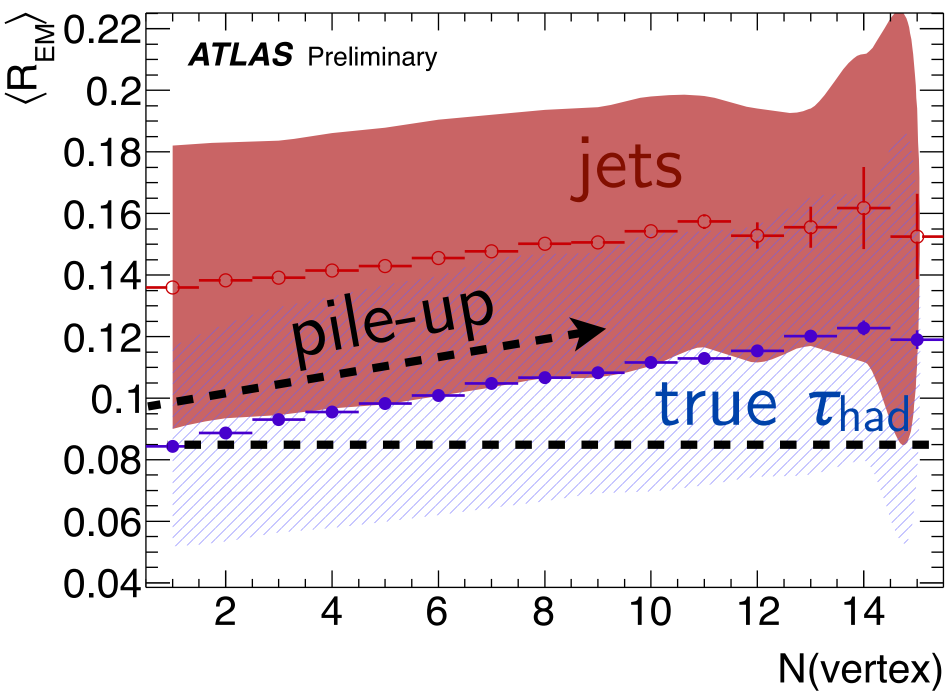
Identification and pile-up

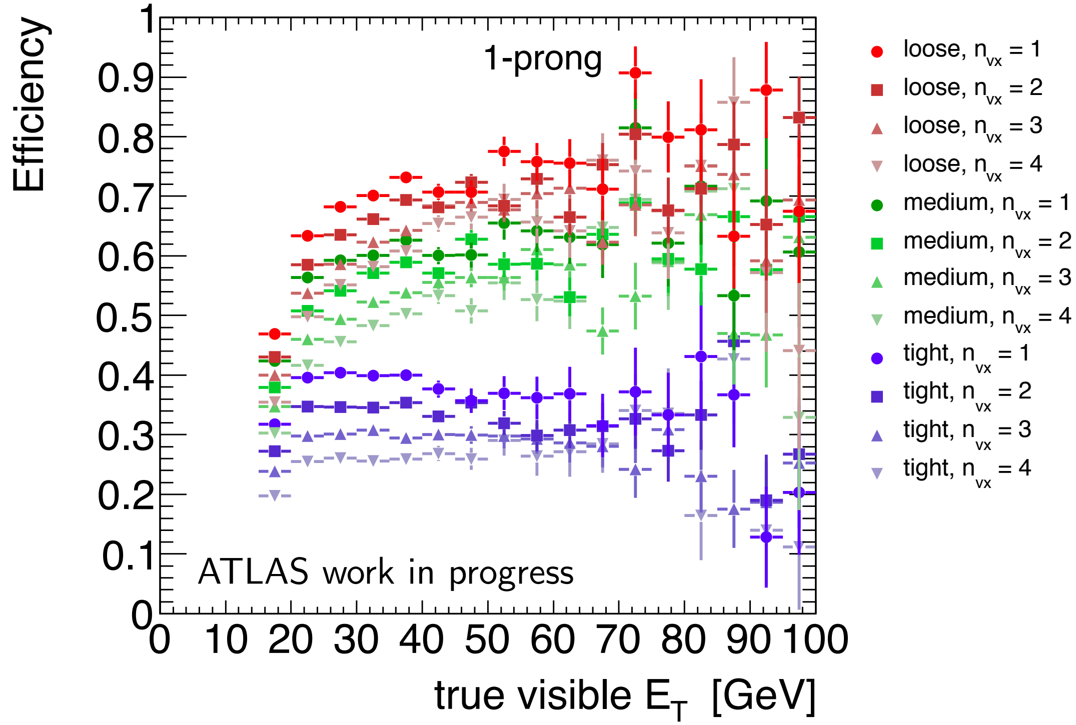
2011

- Important offline variable in 2010-2011: ***EM radius*** - "width of jet in calorimeter"

$$R_{\rm EM} = \frac{\sum_{\{\Delta R<0.4\}} E_{\rm T}^{\rm EM}({\rm cell})\,\Delta R({\rm cell,\ jet})}{\sum_{\{\Delta R<0.4\}} E_{\rm T}^{\rm EM}({\rm cell})}$$

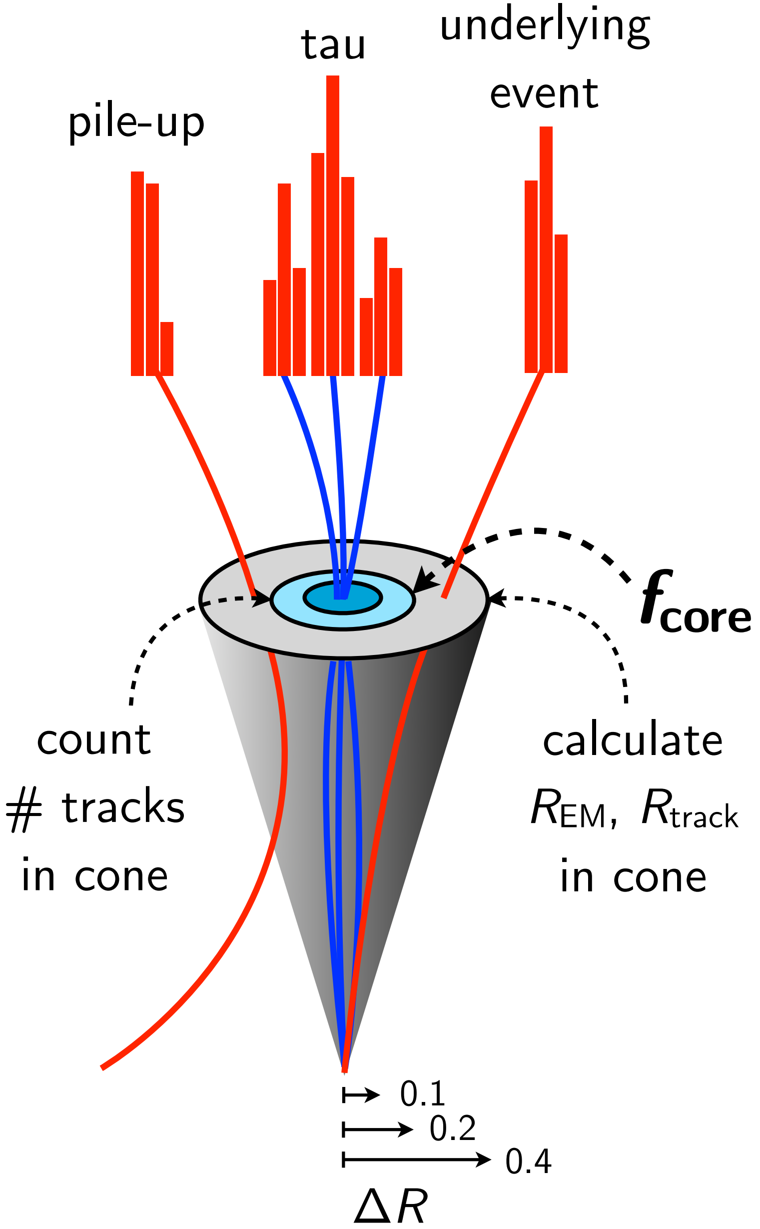
- ***Strong pile-up dependence*** due to using calorimeter deposits in the wide cone: $\Delta R < 0.4$.





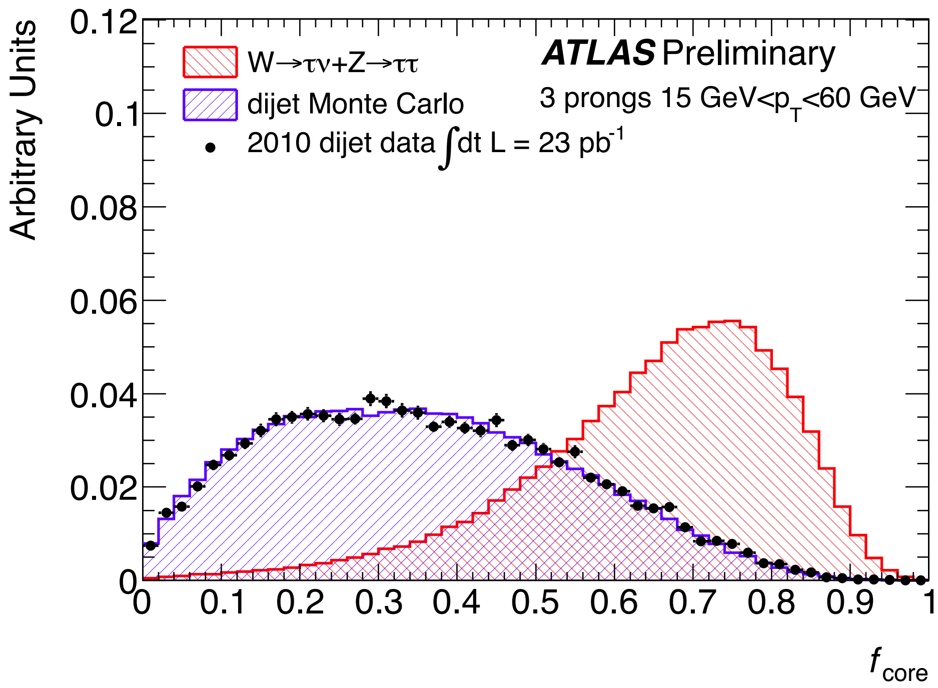
Pile-up robust variables

2011



- Beginning in 2012, the core energy fraction is used instead of R_{EM}, which has less pile-up dependence by using the ratio of energies in smaller ΔR cones of 0.1 and 0.2.

$$f_{\text{core}} = \frac{\sum_{\{\Delta R<0.1\}} E_{\text{T}}^{\text{EM}}(\text{cell})}{\sum_{\{\Delta R<0.2\}} E_{\text{T}}^{\text{EM}}(\text{cell})}$$



[ATLAS-CONF-2011-152]

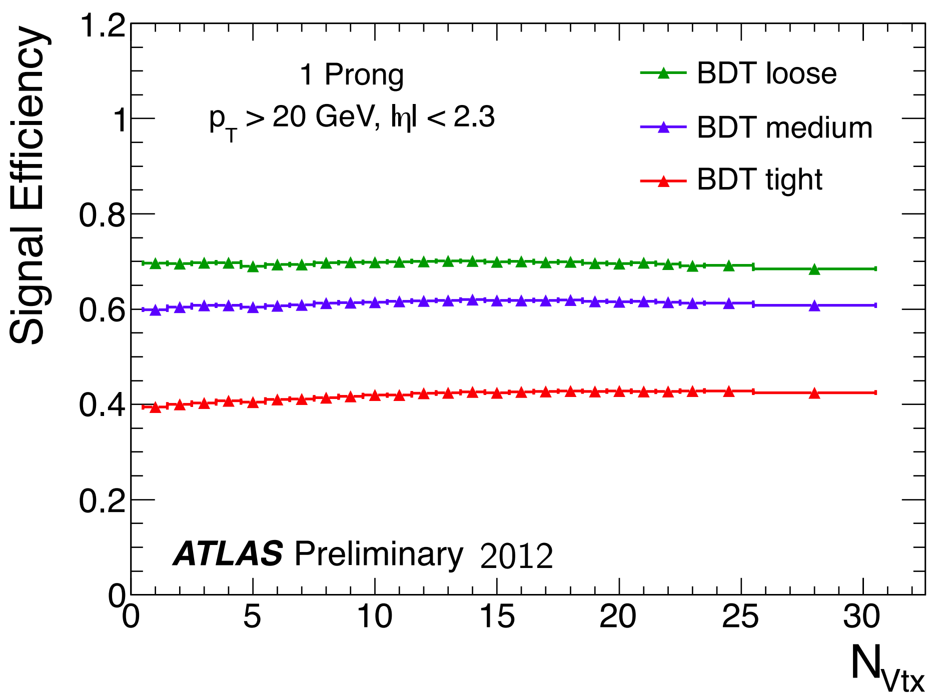
Pile-up corrections

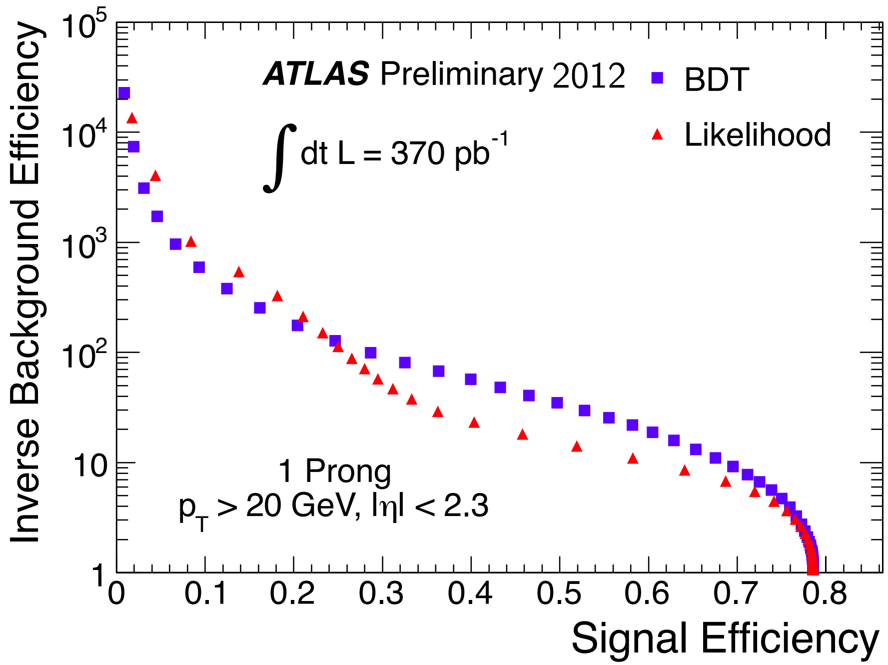
2011

- Also beginning in 2012, the variables with the largest pile-up dependence (f_{core} and f_{track}) are corrected with terms that are linear in the number of reconstructed vertices.

$$f_{\text{core}} = \frac{\sum_{\{\Delta R<0.1\}} E_{\text{T}}^{\text{EM}}(\text{cell})}{\sum_{\{\Delta R<0.2\}} E_{\text{T}}^{\text{EM}}(\text{cell})} + (0.3\%/\text{vertex}) \times N(\text{vertex})$$

- Tight/Medium/Loose working points of the BDT and LLH are defined (≈40%, 60%, 70% efficient), optimized as function of p_{T} and in separate N(vertex) categories.





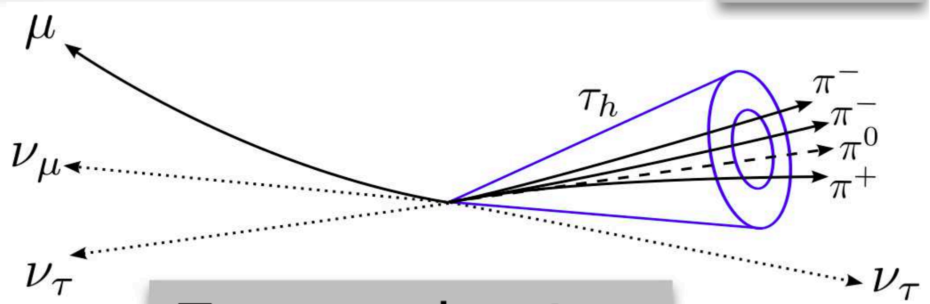
[https://twiki.cern.ch/twiki/bin/view/AtlasPublic/TauPublicCollisionResults]

SM $Z \to \tau\tau$

$Z\to\tau\tau$ studies

'09-'11

- Focus on lep-had final state. Trigger on *e* or *μ*.
- Able to select $Z\to\tau\tau$ control sample for ***studying tau ID***.
- Important to establish understanding of this ***irreducible background to new physics*** with taus: $H\to\tau\tau$ and $Z'\to\tau\tau$ searches.
- ***Complicated background composition***. Multijet, *W*+jet, and *Z*+jet backgrounds all compete at the same order.

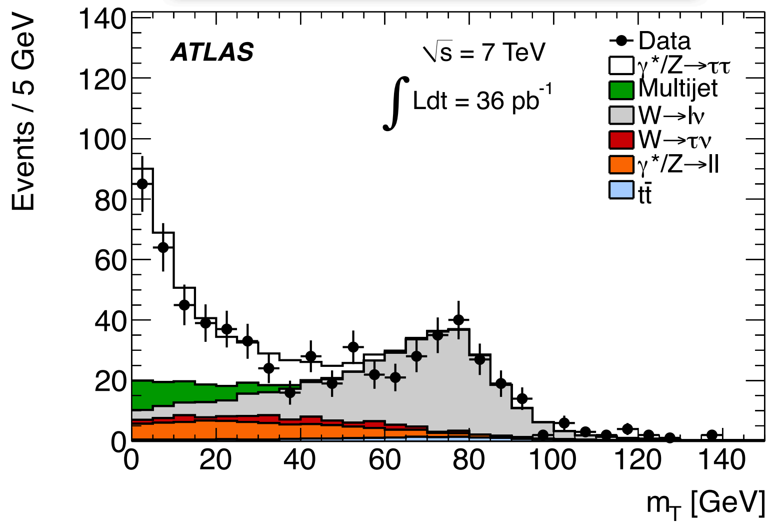


Event selection

- dilepton veto
- $\sum \cos\Delta\phi < -0.15$
- $m_{\mathrm{T}}(\ell, E_{\mathrm{T}}^{\mathrm{miss}}) < 50$ GeV
- $m_{\mathrm{vis}} = 35 - 75$ GeV
- $N_{\mathrm{track}}(\tau_{\mathrm{h}}) = 1$ or 3
- $|\mathrm{charge}(\tau_{\mathrm{h}})| = 1$
- Opposite sign τ_{h} and ℓ

Data-driven background estimates:

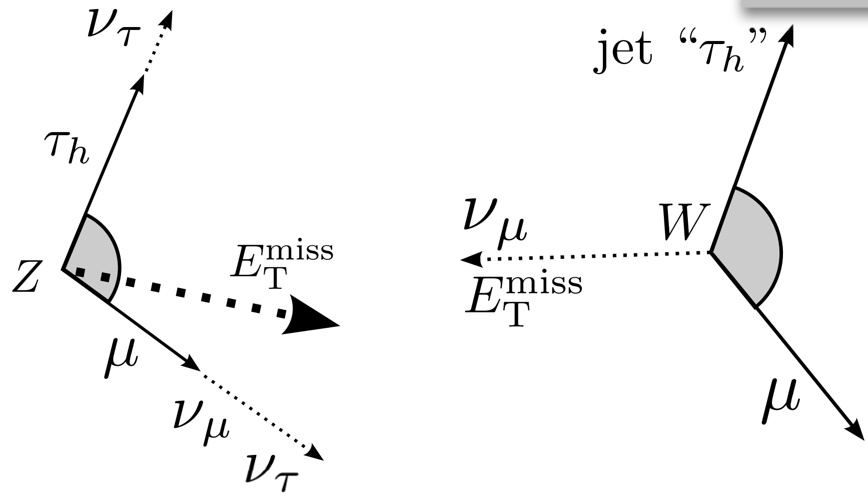
- Jet fake rate mis-modeled in ATLAS Monte Carlo. W+jet background normalized with high-m_T control region.
- Multijet background modeled from same-sign data.



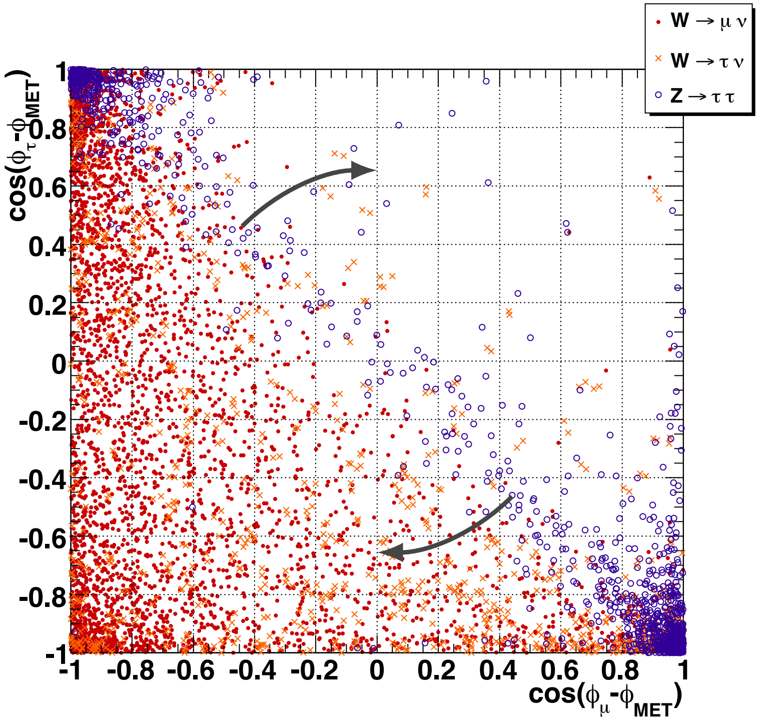
$\sum \cos\Delta\phi$ for suppressing W+jet

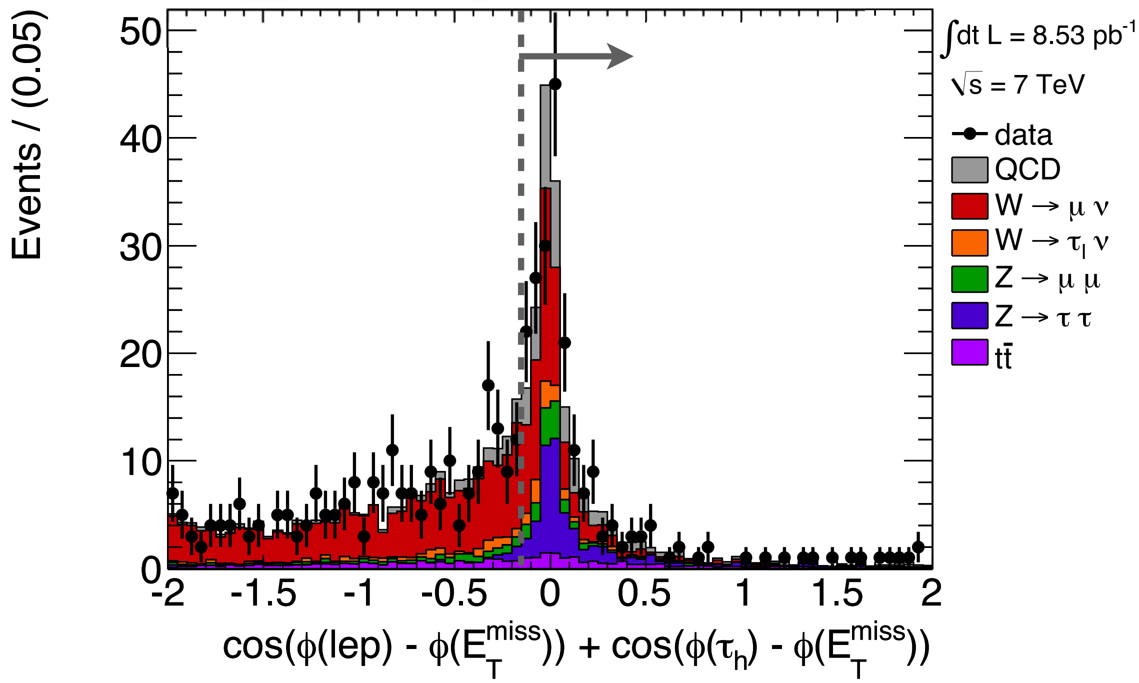
'09-'10

- ***I introduced a new variable*** for rejecting the *W*+jet background.
- Exploits the correlation of the direction of the E_T^{miss} and the decay products.
- Essentially require the E_T^{miss} to be between instead of away.
- Only dependent on E_T^{miss} direction and not the magnitude ⇒ less systematics.



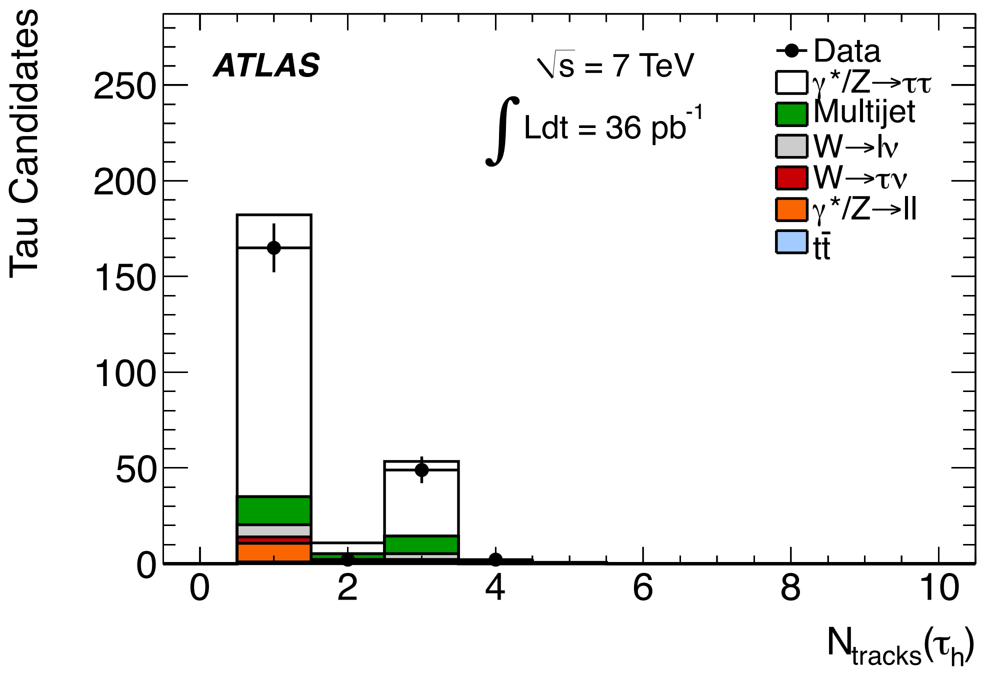
$\sum \cos\Delta\phi$ > -0.15

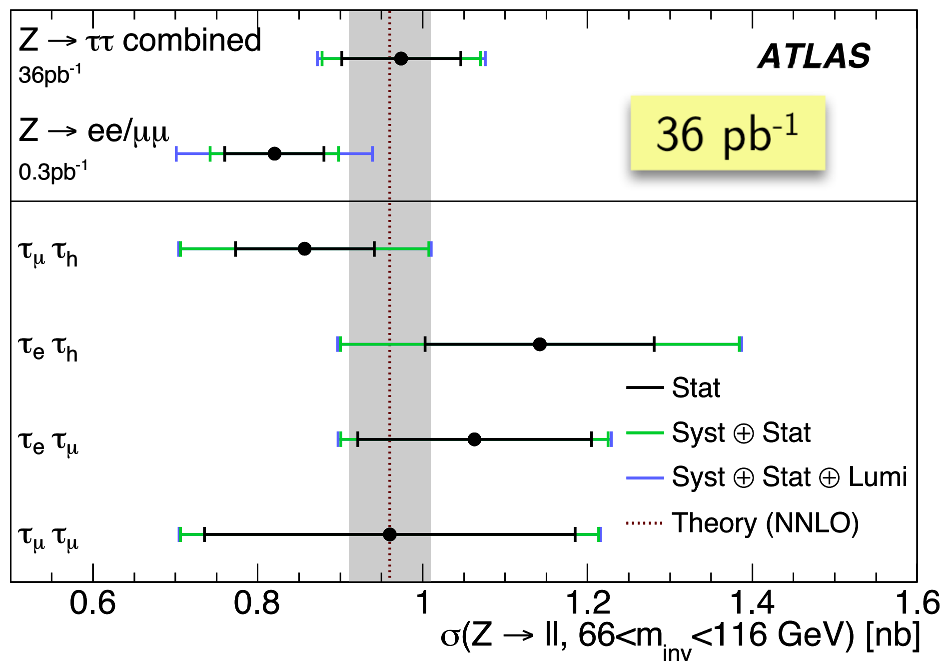




$Z\rightarrow\tau\tau$ cross section

'10-'11





- Claimed observation of $Z\rightarrow\tau\tau$ with 8.5 pb^{-1}.
- Measured cross section to 10% with 36 pb^{-1}, consistent with SM.
- ***Published in PRD***.

$$\sigma_{\text{combined}} = 0.97 \pm 0.07(\text{stat.}) \pm 0.07(\text{sys.}) \pm 0.03(\text{lumi.})\ \text{nb}$$

$$\sigma_{\text{theory}} = 0.96 \pm 0.05\ \text{nb at NNLO}$$

[Phys. Rev. D 84, 112006 (2011) arxiv:1108.2016v2]

New physics:

$Z' \to \tau\tau$

Why the Standard Model?

- Why the **gauge group** $SU(3)_C \times SU(2)_L \times U(1)_Y$?
- Why are there **3 generations** of quarks and leptons?
- Why are lepton and hadron charges quantized in the same units? Why the existing **hypercharges**?

$$Q_{\mathrm{EM}} = T_{3\mathrm{L}} + Y/2$$

Is it because...

- ***the gauge group of Nature is actually bigger?***

$SO(10) \rightarrow SU(5) \times U(1)$ Georgi-Glashow

$SO(10) \rightarrow SU(4)_C \times SU(2)_L \times SU(2)_R$ Pati-Salam

1974

- *e.g.* Pati-Salam SO(10): $Q_{\mathrm{EM}} = T_{3\mathrm{L}} + T_{3\mathrm{R}} + 1/2(B - L)$
- lifetime of the proton $> 10^{33}$ years $\Rightarrow$ if unification happens it must suppress proton decay, *e.g.* it happens at a high energy scale

GUT Motivations

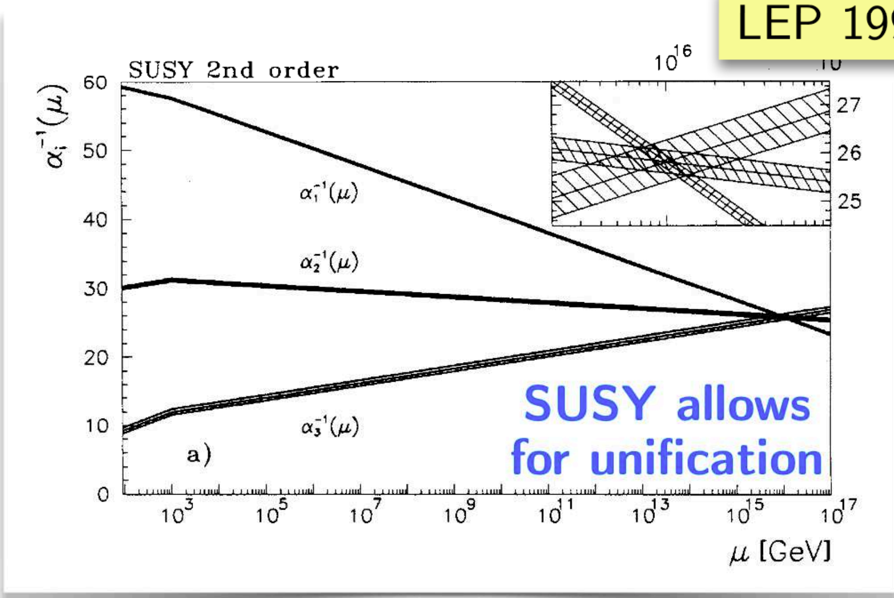
It is interesting that:

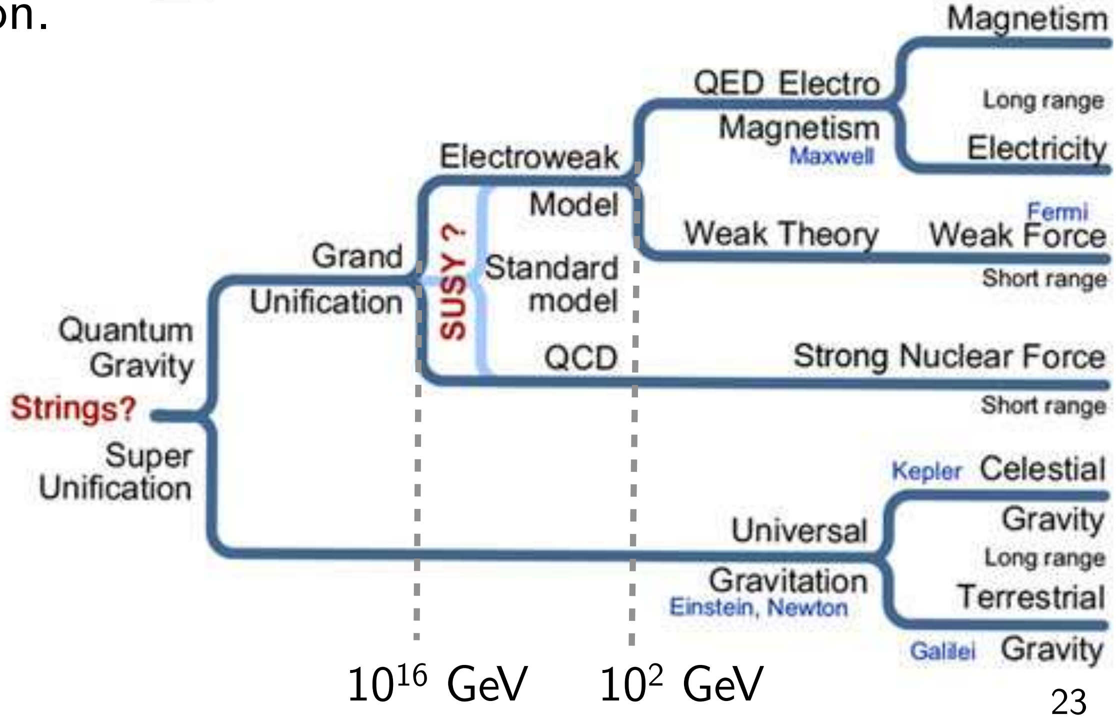
- the structure of SM can be embedded in larger groups, and this ***could explain the SM hypercharges***,
- the SM ***couplings apparently converge*** when run to very high energies.

Motivates the possibility of grand unification.

- ***Z′* bosons occur in theories with additional U(1) symmetries.**
- **Best limits on $Z' \to ee/\mu\mu$ are**
  M > 2.3 TeV CMS [arxiv:1206.1849]
  M > 2.2 TeV ATLAS [arxiv:1209.2535]
- **Important to test the couplings to all lepton flavors.**

LEP 1991





$Z' \rightarrow \tau\tau \rightarrow \mu\tau_h$

'11-'12

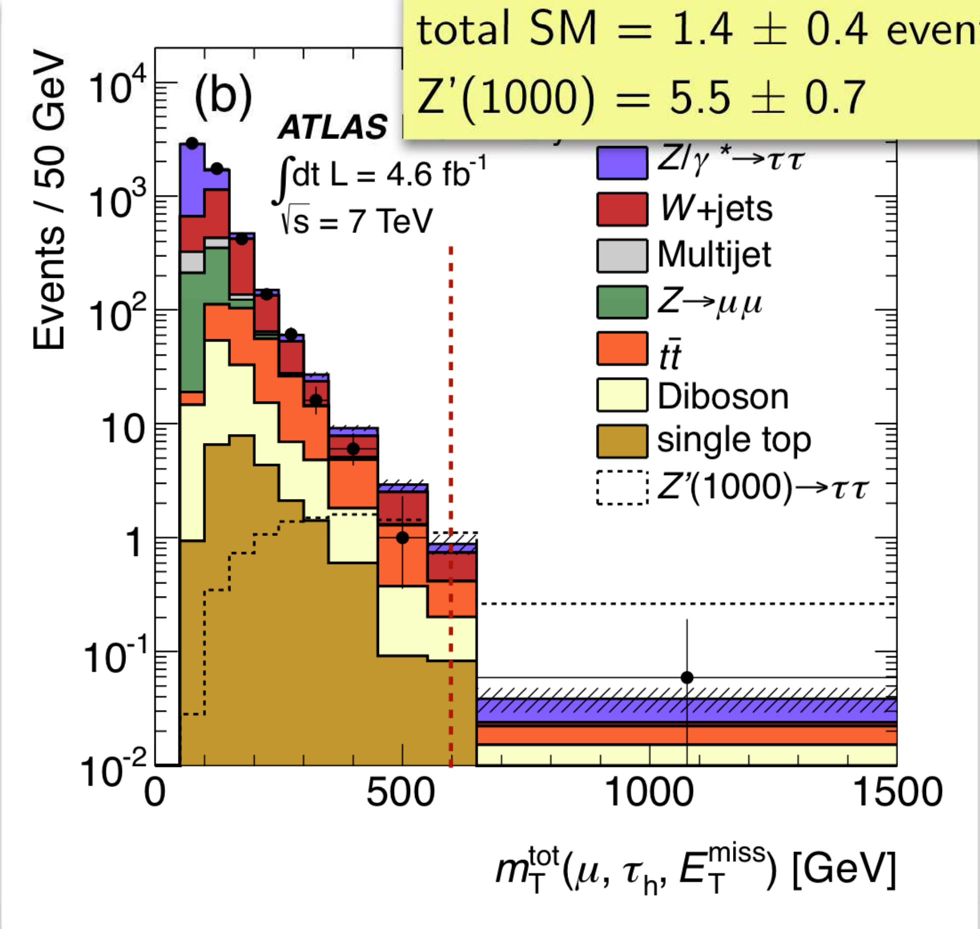
- Select OS back-to-back tau decays.
- Count high-mass events.

Event selection

- $p_T(\mu) > 25$, $p_T(\tau_h) > 35$ GeV
- 1-prong τ_h
- $|\Delta\phi(e, \tau_h)| > 2.7$
- opposite sign μ and τ_h
- $m_T(e, \tau_h, E_T^{miss}) > 600$ GeV

total SM = 1.4 ± 0.4 events

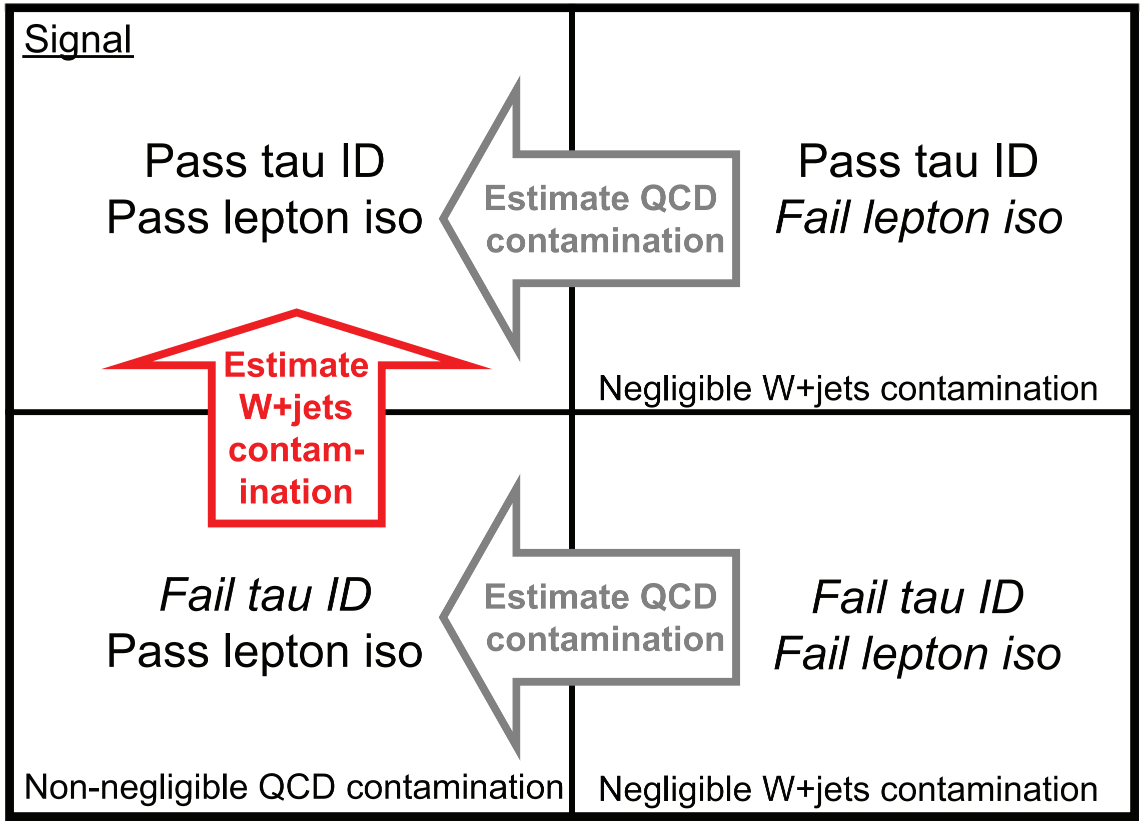
Z'(1000) = 5.5 ± 0.7



- Fake factor methods used to model multijet and W+jet backgrounds
- Need to be modeled in data-driven ways for two reasons:
  1. populate the model in the high-mass tail
  2. jet$\rightarrow$ τ_h fake rate is mis-modeled in Monte Carlo.

Double fake factor procedure

'11-'12



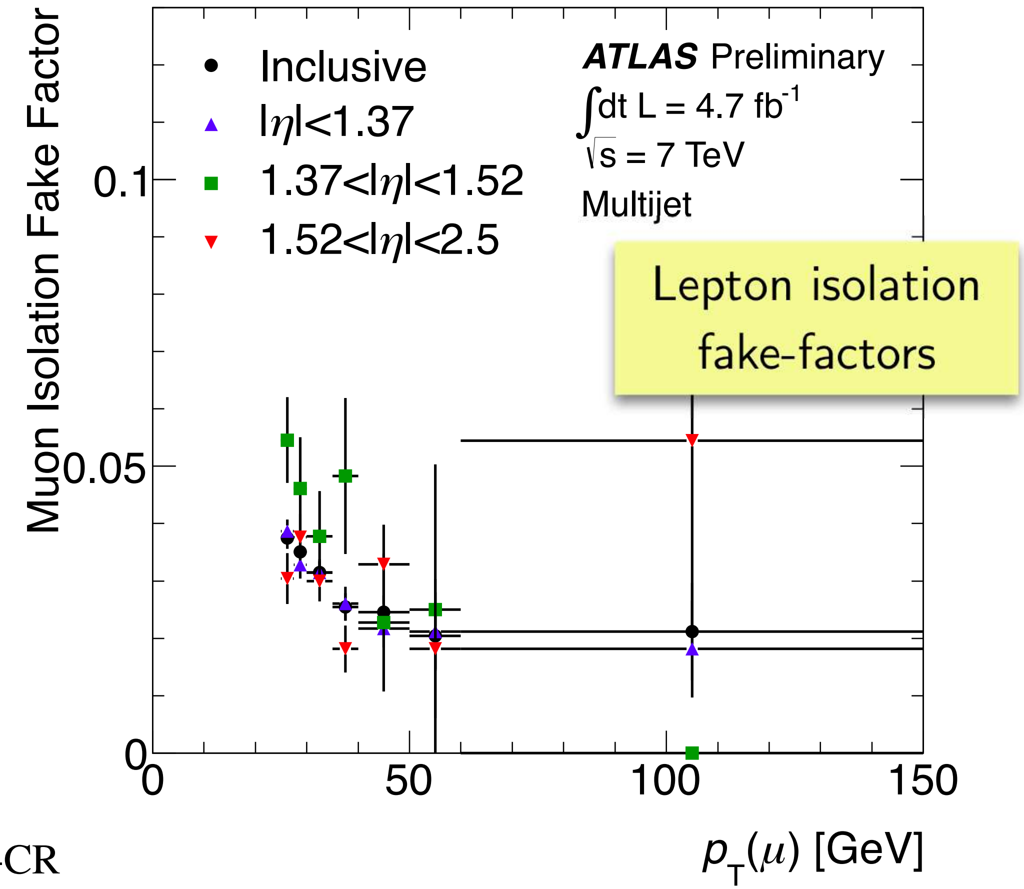
- The multijet contamination is estimated from the rate of non-isolated leptons, in both the signal region that passes tau ID, and the sample that fails.
- Then, the corrected number of tau candidates failing ID are weighted to predicted the W+jet background.
- This way, *the corrections are small at each step*.

Multijet background estimation

'11-'12

Multijet control region

- no isolation
- $E_T^{miss} < 30$ GeV
- $m_T(\mu, E_T^{miss}) < 30$ GeV



- In the control region, divide leptons into pass and fail isolation.
- Define fake factor:

$$f_{\mu\text{–iso}}(p_T, \eta) \equiv \left. \frac{N^{\text{pass } \mu\text{–iso}}(p_T, \eta)}{N^{\text{fail } \mu\text{–iso}}(p_T, \eta)} \right|_{\text{multijet–CR}}$$

- Predict the number of multijet events:

$$N_{\text{multijet}}(p_T, \eta, x) = f_{\mu\text{–iso}}(p_T, \eta) \cdot N^{\text{fail } \mu\text{–iso}}_{\text{multijet}}(p_T, \eta, x)$$

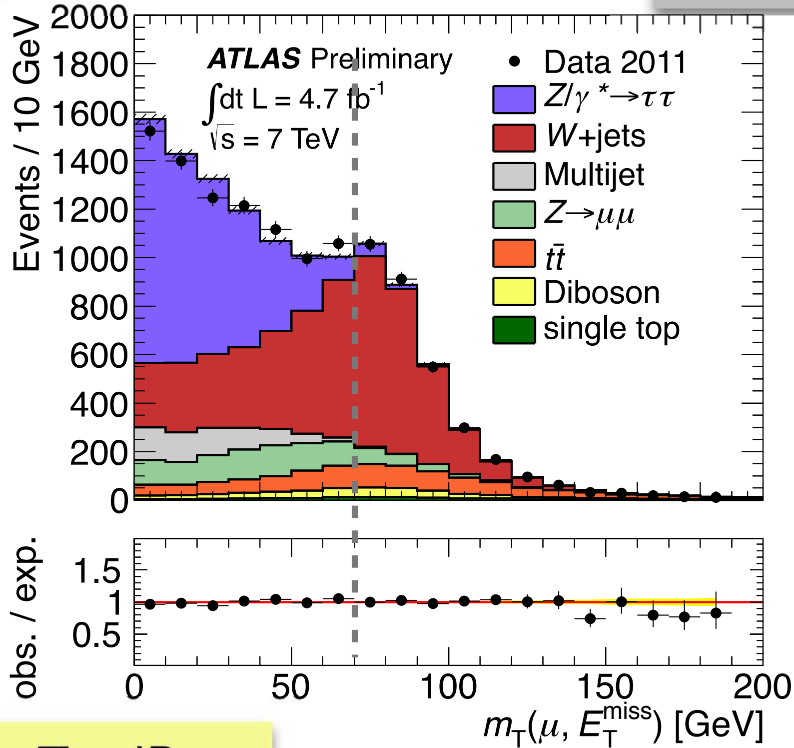
$$= f_{\mu\text{–iso}}(p_T, \eta) \cdot \left(N^{\text{fail } \mu\text{–iso}}_{\text{data}}(p_T, \eta, x) - N^{\text{fail } \mu\text{–iso}}_{\text{MC}}(p_T, \eta, x) \right)$$

W+jet background estimation

'11-'12

W+jet control region

- $m_T(\mu,\ E_T^{miss}) = 70\text{-}200$ GeV
- isolated lepton

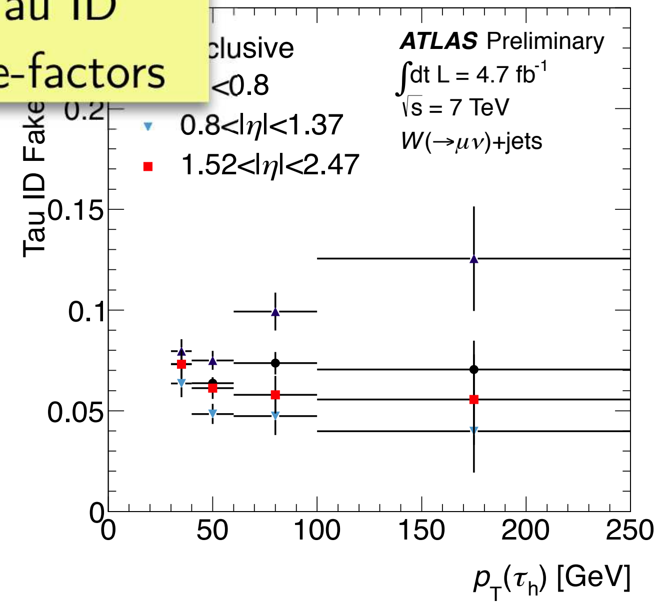


- In a W+jet control region, divide tau candidates into pass and fail identification.

- Define fake factor:

$$f_\tau(p_T, \eta) \equiv \left.\frac{N^{\text{pass }\tau\text{–ID}}(p_T, \eta)}{N^{\text{fail }\tau\text{–ID}}(p_T, \eta)}\right|_{\text{W–CR}}$$

Tau ID fake-factors



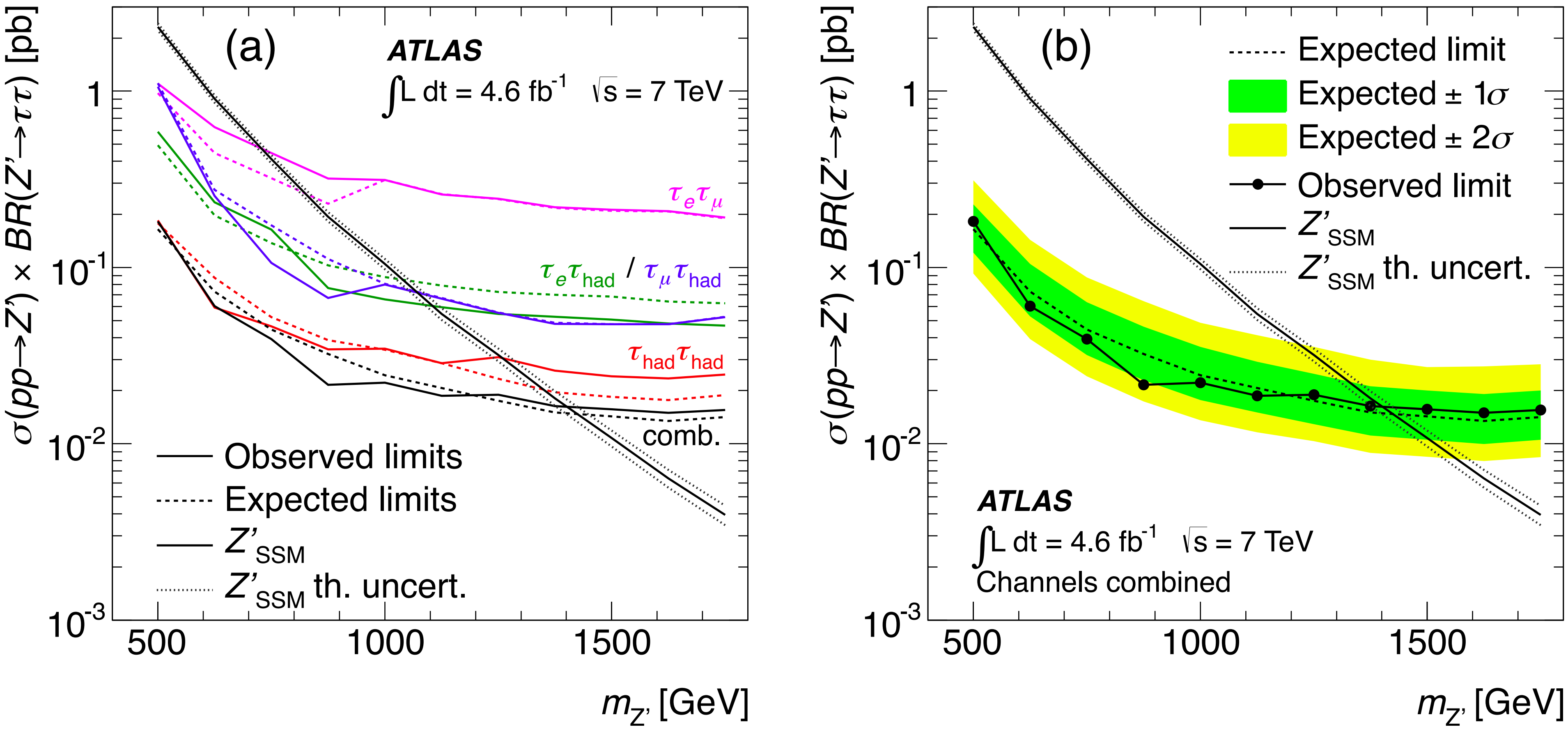
- Predict the number of W/Z+jet events:

$$N_{W/Z\text{+jet}}(p_T, \eta, x) = f_\tau(p_T, \eta) \cdot N^{\text{fail }\tau\text{–ID}}_{W/Z\text{+jet}}(p_T, \eta, x)$$

$$= f_\tau(p_T, \eta) \cdot \left(N^{\text{fail }\tau\text{–ID}}_{\text{data}}(p_T, \eta, x) - N^{\text{fail }\tau\text{–ID}}_{\text{multijet}}(p_T, \eta, x) - N^{\text{fail }\tau\text{–ID}}_{\text{MC}}(p_T, \eta, x)\right)$$

Combined limit

'11-'12



ATLAS Z' SSM Exclusions: observed (expected) @ 95% CL

- $\tau_h\tau_h$: 1.26 (1.35) TeV
- $\mu\tau_h$: 1.07 (1.06) TeV
- $e\tau_h$: 1.10 (1.03) TeV
- $e\mu$: 0.72 (0.82) TeV

combined: 1.40 (1.42) TeV

Preprint submitted to PLB

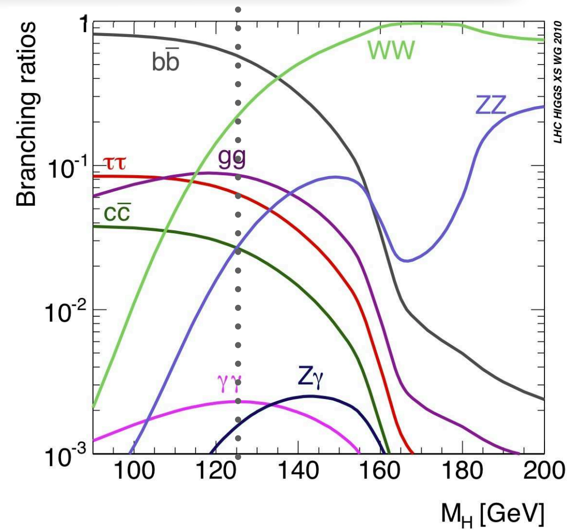
[arxiv:1210.6604]

Review of Higgs search results

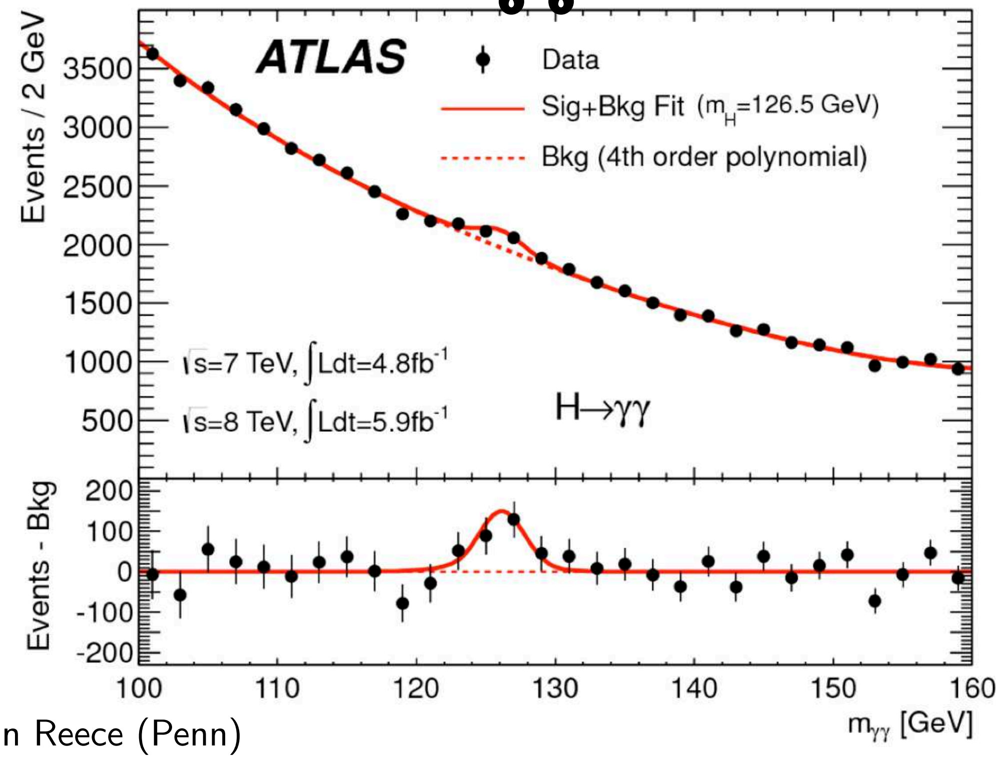
ATLAS

Current Higgs results

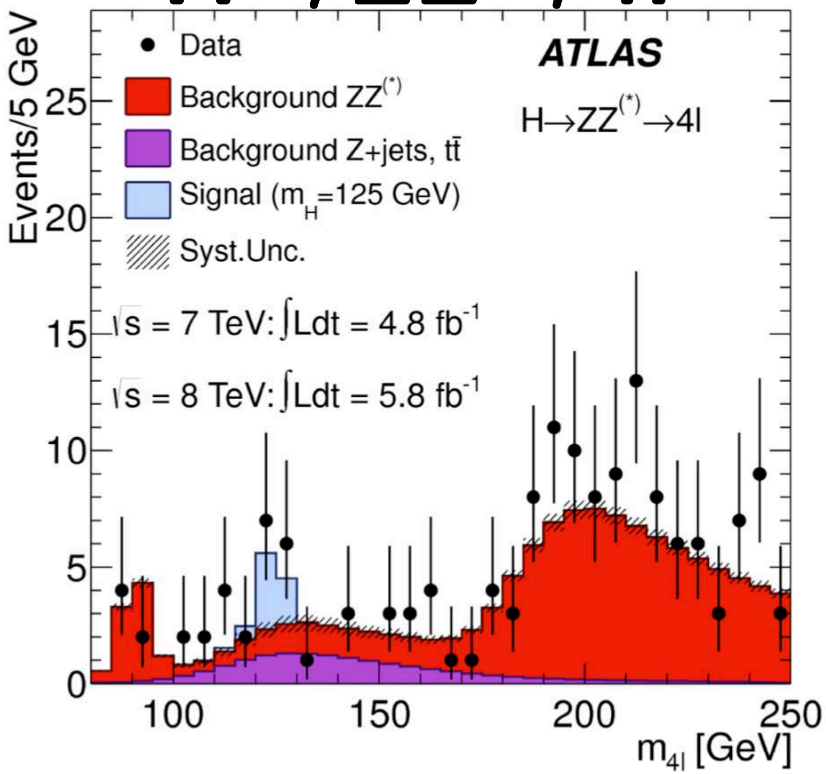
- Observation claimed July 4th 2012.
- Two channels with precise mass measurements: $H \to \gamma\gamma$, $H \to ZZ \to 4l$.
- $H \to WW$ also observes a broad excess.
- $H \to \tau\tau$ and $H \to bb$ are approaching sensitivity.



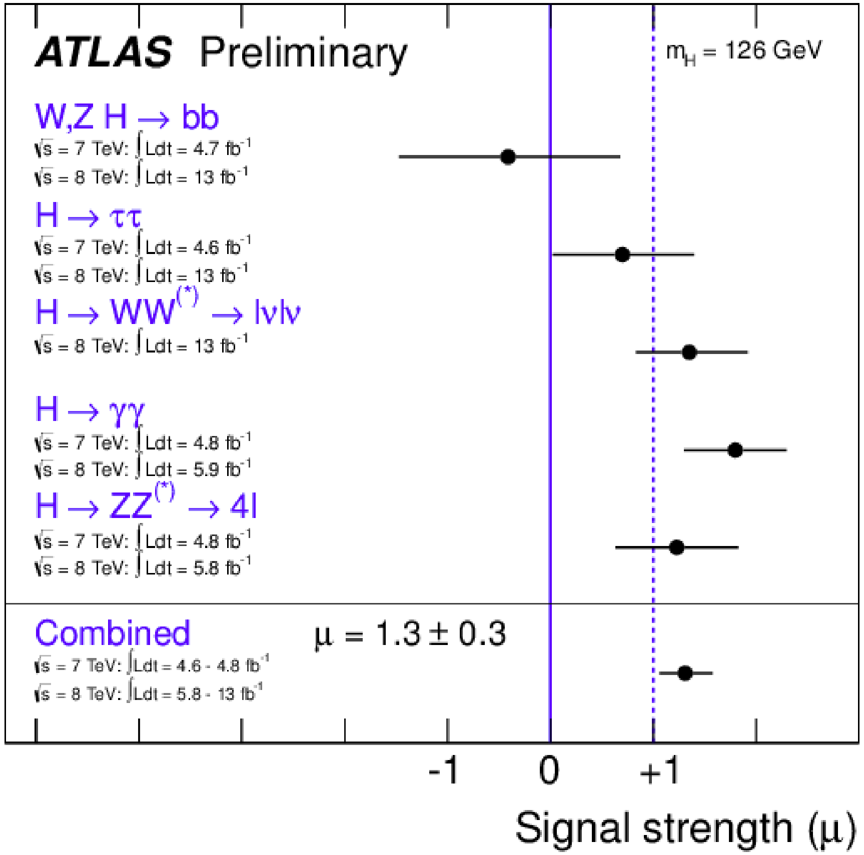
$H \to \gamma\gamma$

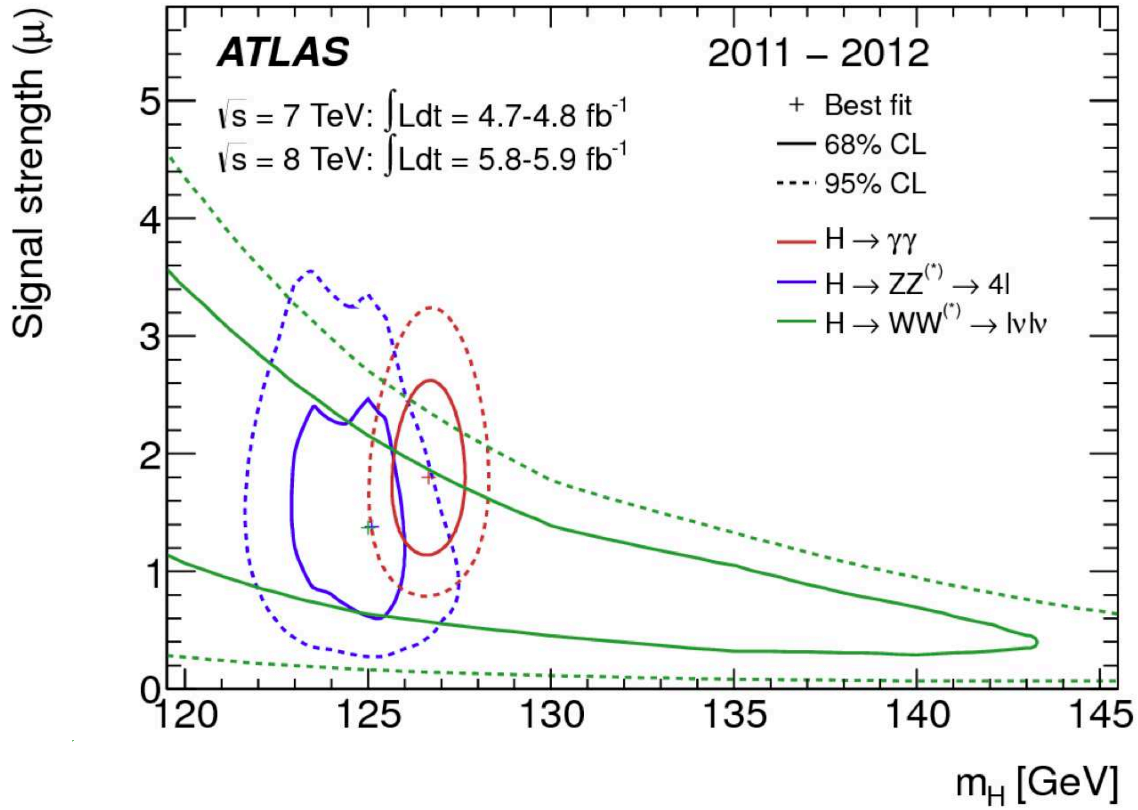


$H \to ZZ \to 4l$

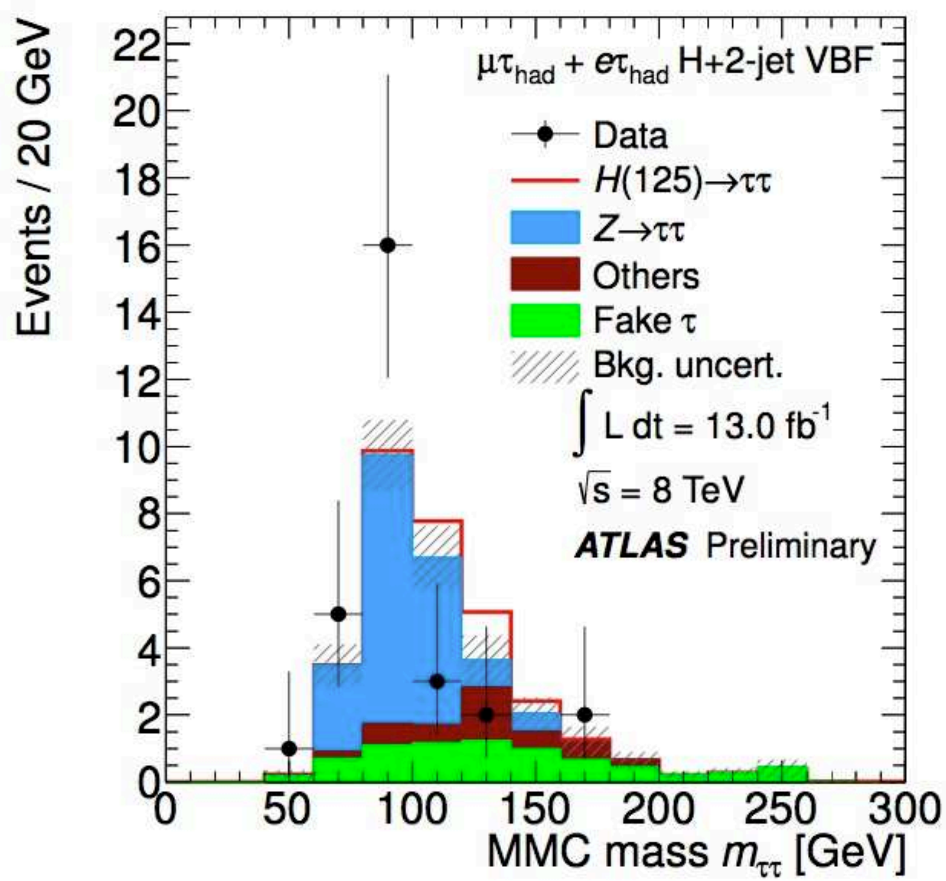


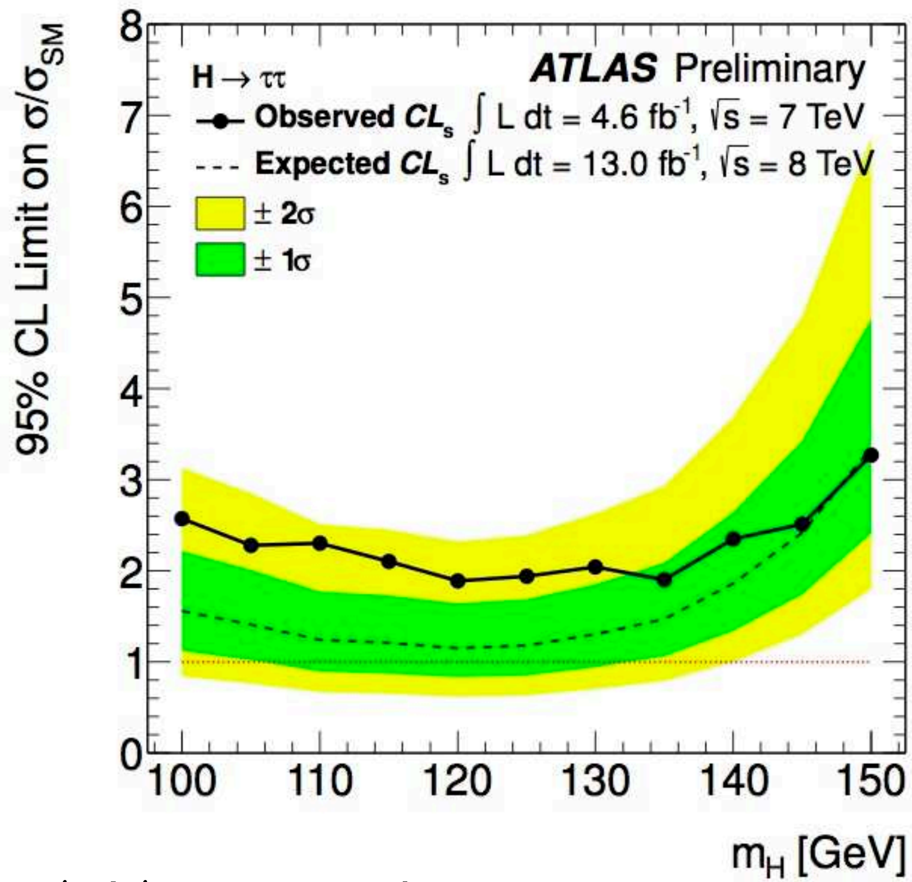
Current Higgs results





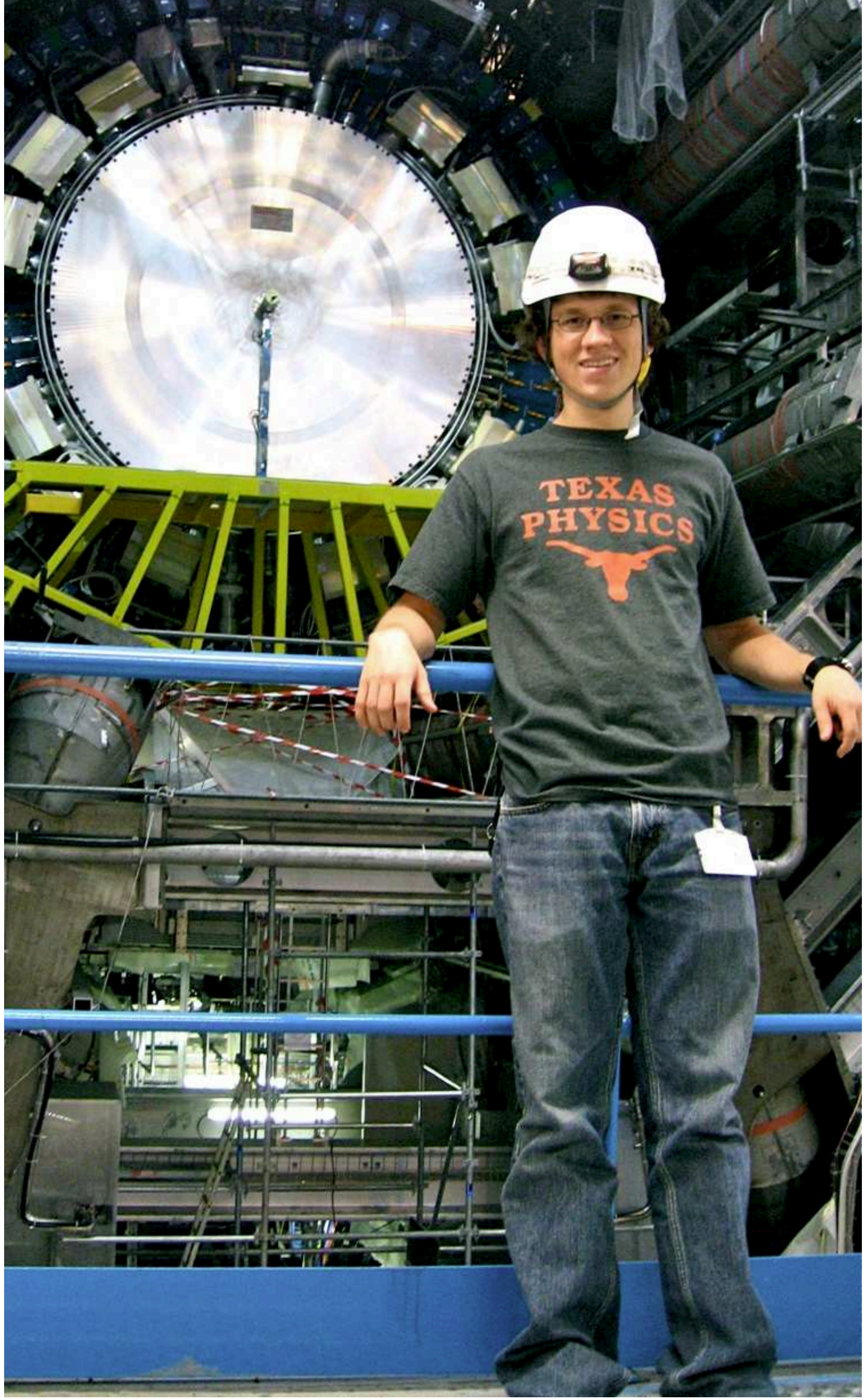
Current $H \to \tau\tau$ result





- A lot of shared experience between $Z/Z'/H \to \tau\tau$ analyses.
- Uses similar $\sum \Delta\phi$ cut for suppressing W+jet.
- Uses fake factor method for predicting fake backgrounds.
- Eagerly approaching sensitivity to $1\times$SM $H \to \tau\tau$.
- 21.7 fb^{-1} collected this year.

Conclusions



- It's been an exciting time to be a student.
- Had a lot of fun learning experimental skills during the turn-on of the LHC.
- Developed cut-based tau identification.
- Studied pile-up robustness for taus.
- Measured SM $Z \to \tau\tau$ cross section.
- Searched for new physics in high-mass ditau events.
- *Excited that discoveries at the LHC could just be beginning.*

Thanks so much for your interest and attention!

Back up

ATLAS

$$Z' \rightarrow \tau\tau$$

ATLAS

Event summary

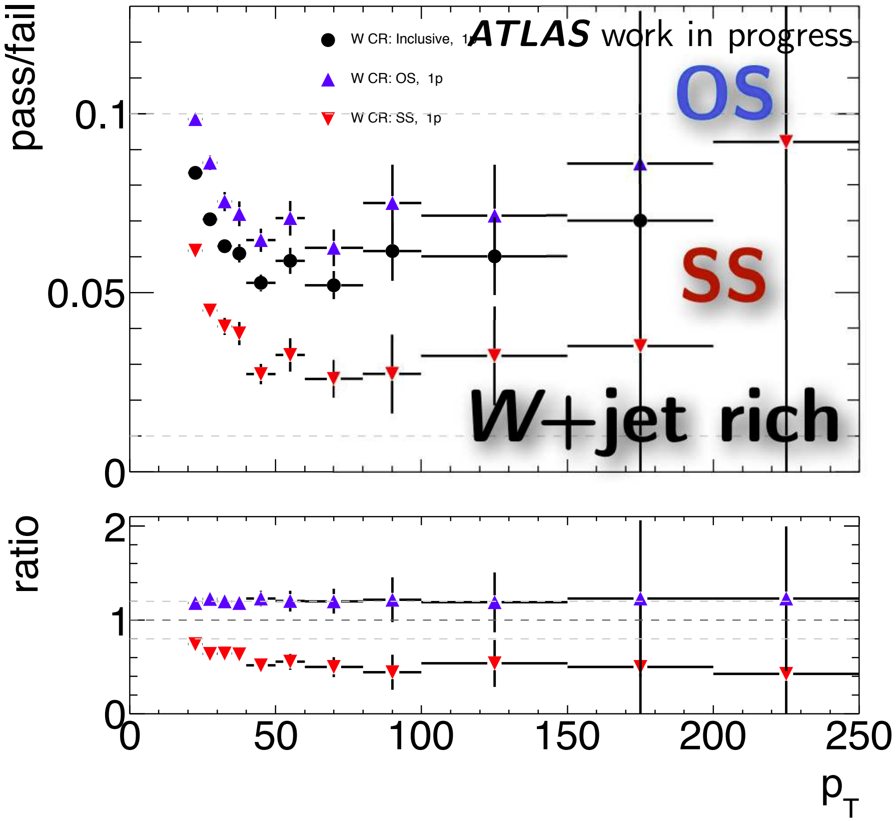
	$\tau_{\rm had}\tau_{\rm had}$	$\tau_\mu\tau_{\rm had}$	$\tau_e\tau_{\rm had}$	$\tau_e\tau_\mu$
$m_{Z'}$ [GeV]	1250	1000	1000	750
$m_{\rm T}^{\rm tot}$ threshold [GeV]	700	600	500	350
$Z/\gamma^* \to \tau\tau$	0.73±0.23	0.36±0.06	0.57±0.11	0.55±0.07
W+jets	< 0.03	0.28±0.22	0.8 ±0.4	0.33±0.10
$Z(\to \ell\ell)$+jets	< 0.01	< 0.1	< 0.01	0.06±0.02
$t\bar{t}$	< 0.02	0.33±0.15	0.13±0.09	0.97±0.22
Diboson	< 0.01	0.23±0.07	0.06±0.03	1.69±0.24
Single top	< 0.01	0.19±0.18	< 0.1	< 0.1
Multijet	0.24±0.15	< 0.01	< 0.1	< 0.01
Total expected background	0.97±0.27	1.4 ±0.4	1.6 ±0.5	3.6 ±0.4
Events observed	2	1	0	5
Expected signal events	6.3 ±1.1	5.5 ±0.7	5.0 ±0.5	6.72±0.26
Signal efficiency (%)	4.3	1.1	1.0	0.4

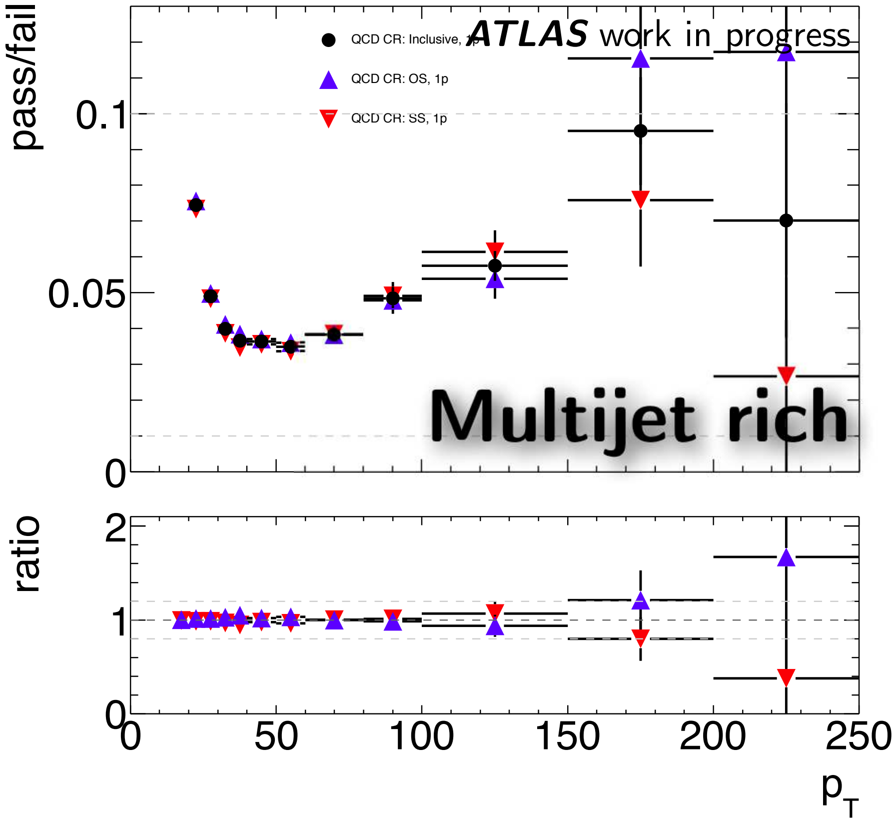
Systematics

Uncertainty [%]	Signal				Background			
	hh	μh	eh	$e\mu$	hh	μh	eh	$e\mu$
Stat. uncertainty	1	2	2	3	5	20	23	7
Eff. and fake rate	16	10	8	1	12	16	4	3
Energy scale and res.	5	7	6	2	$^{+22}_{-11}$	3	8	5
Theory cross section	8	6	6	5	9	4	4	5
Luminosity	4	4	4	4	2	2	2	4
Data-driven methods	–	–	–	–	$^{+21}_{-11}$	6	16	–

Table 2: Uncertainties on the estimated signal and total background contributions in percent for each channel. The following signal masses, chosen to be close to the region where the limits are set, are used: 1250 GeV for $\tau_{\rm had}\tau_{\rm had}$ (hh); 1000 GeV for $\tau_\mu\tau_{\rm had}$ (μh) and $\tau_e\tau_{\rm had}$ (eh); and 750 GeV for $\tau_e\tau_\mu$ ($e\mu$). A dash denotes that the uncertainty is not applicable. The statistical uncertainty corresponds to the uncertainty due to limited sample size in the MC and control regions.

Observed variance in fake-rates





(BDTMedium)

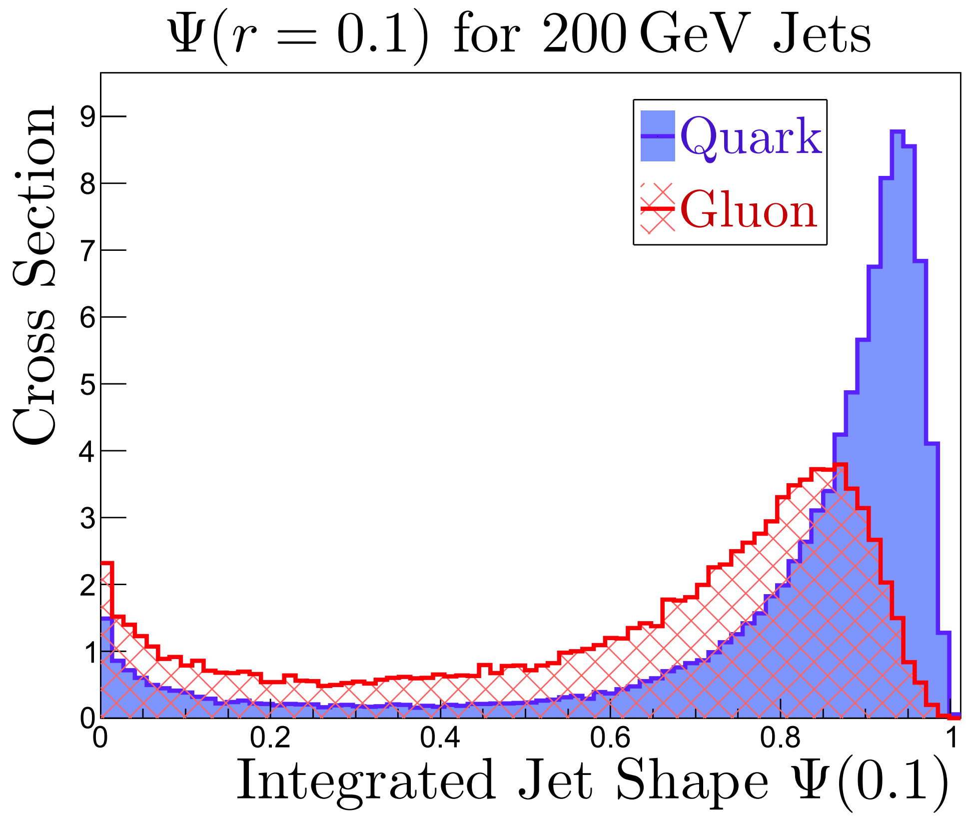
- Hypothesis: quarks vs gluons
- Divide the issue into two questions:

  1. *Why do quarks and gluons have different tau fake-rates?*
  2. *How does the quark/gluon fraction vary among samples?*

Jet width for quark/gluons

Why do quarks and gluons have different tau fake-rates?

- $\Psi(r)$ = fraction of jet energy within $\Delta R < r$.
- Quark jets are more narrow than gluon jets of the same energy.
- Tau identification prefers narrow candidates.



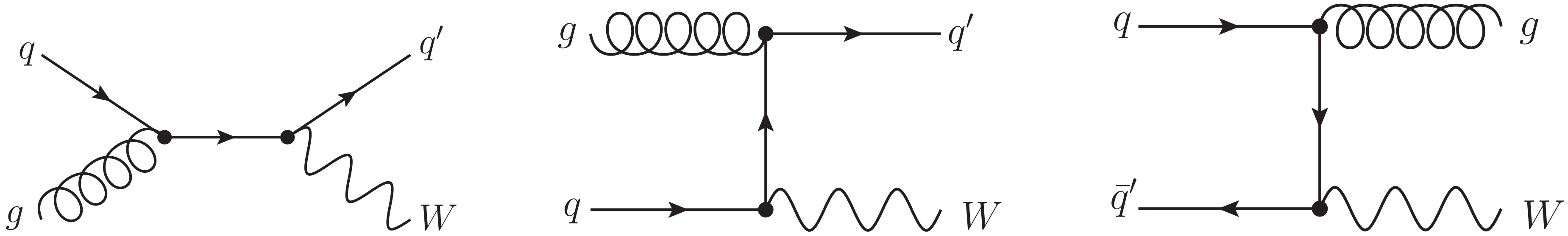
- This is consistent with samples of quark-enriched jets, like W+jet, having higher fake-rates.

J. Gallicchio, M. Schwartz. "Quark and Gluon Tagging at the LHC". arXiv:1106.3076.

OS vs SS *W*+jet

How does the quark/gluon fraction vary among samples?

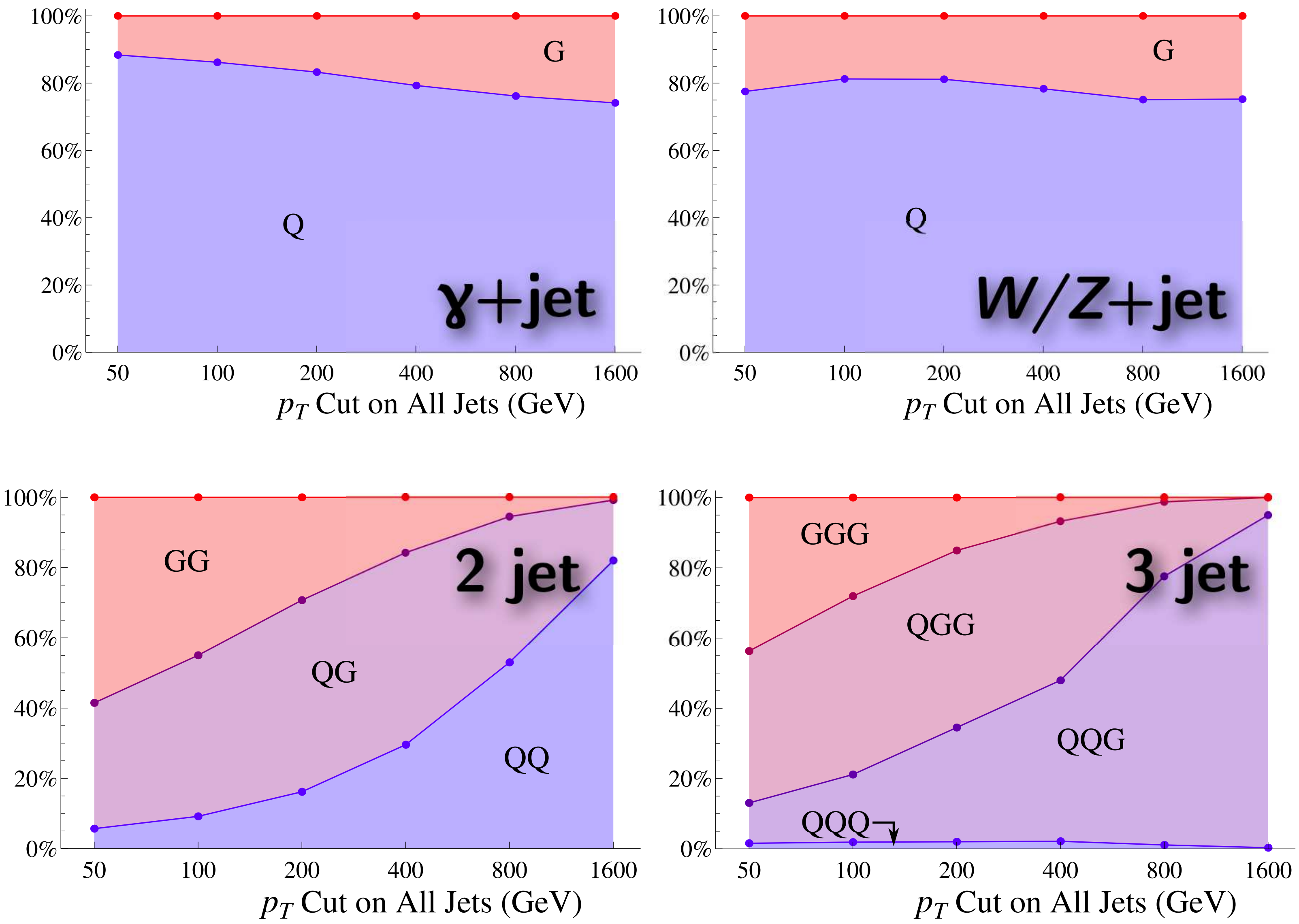
Leading order *W*+jet production:



(a) (b) (c)

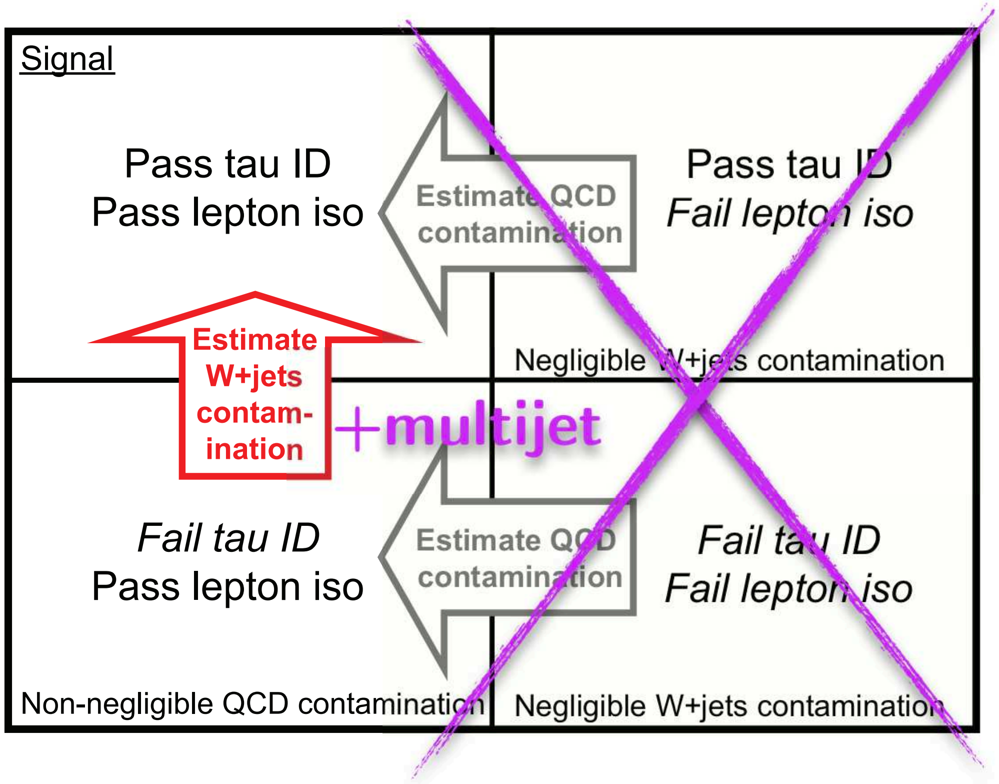
- The charge of the quark should correlate with the reconstructed charge of the tau candidate, therefore (a) and (b) preferably produce opposite sign W+jet events.
- OS and SS will have different quark/gluon fractions.

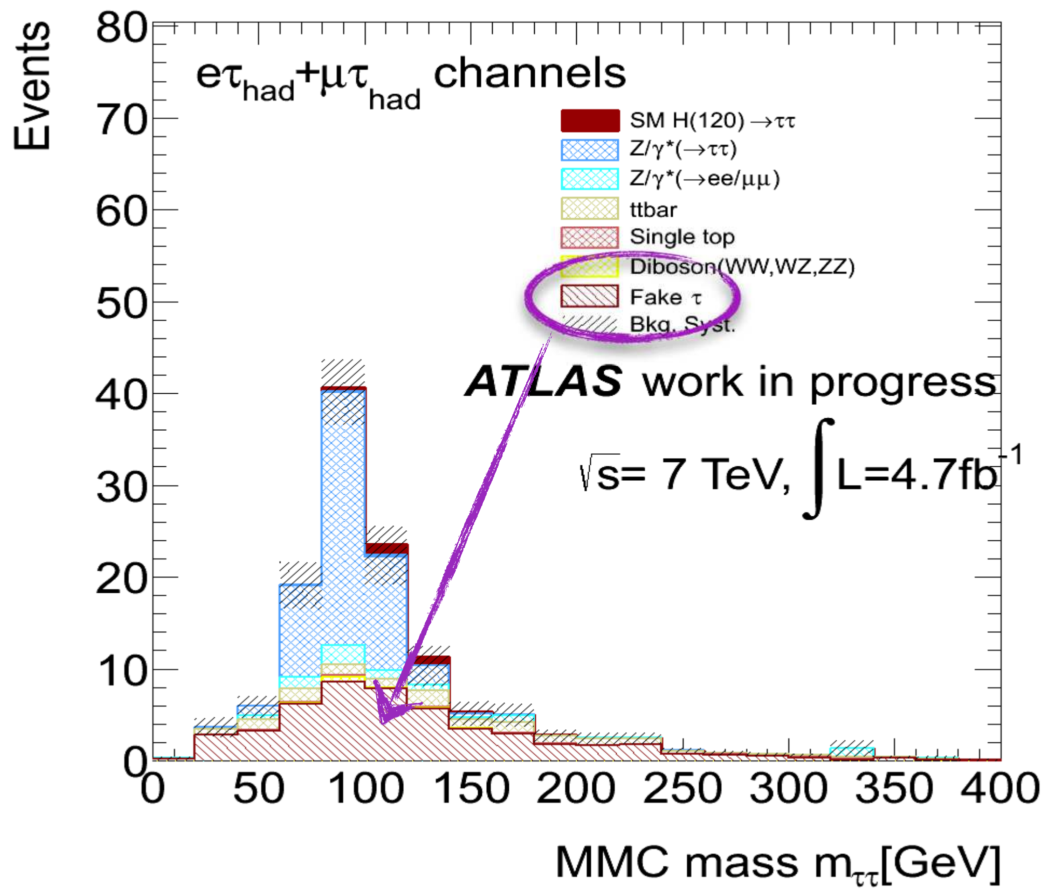
Madgraph predicted Quark/Gluon



J. Gallicchio, M. Schwartz. "Pure Samples of Quark and Gluon Jets at the LHC". arXiv:1104.1175

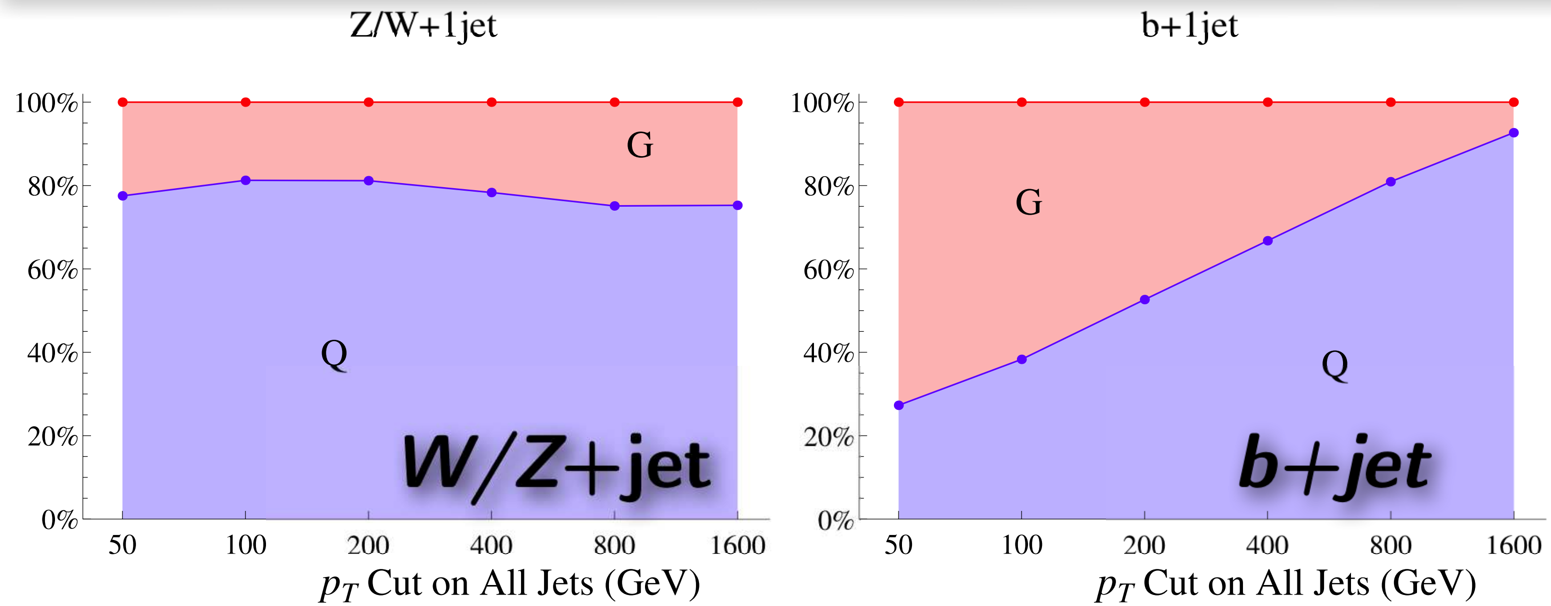
Cross-check: single fake factor method



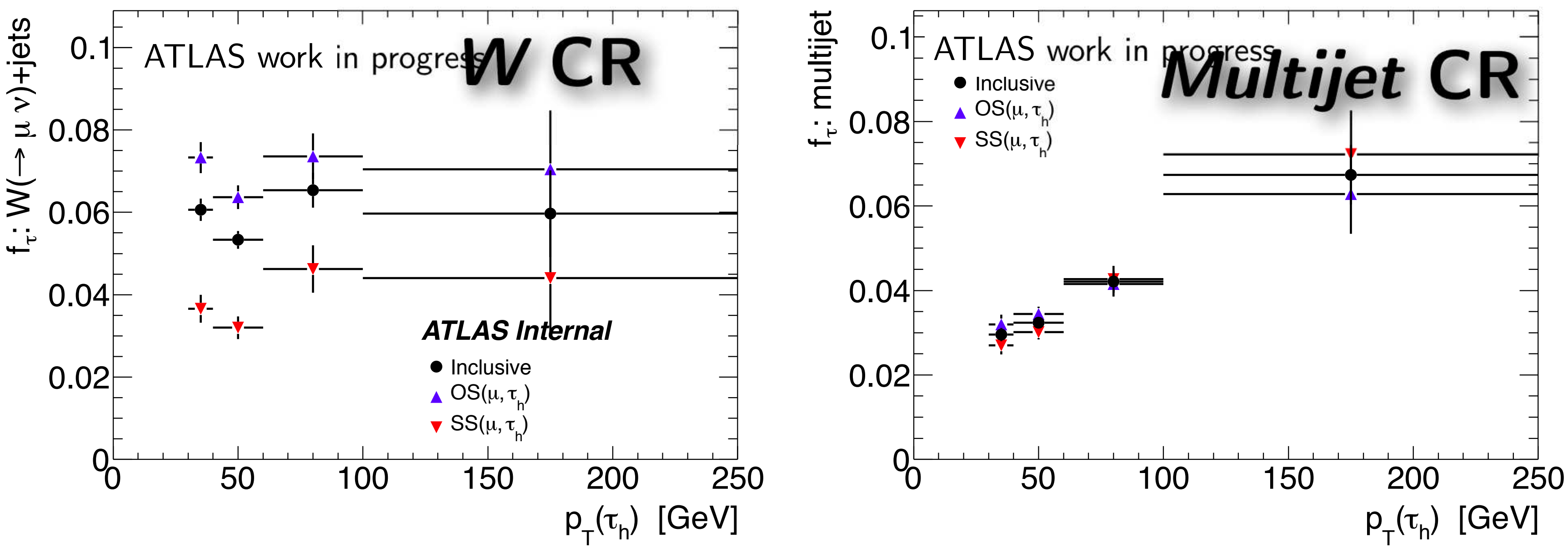


- Avoid issues of subtracting multijet from W/Z+jets background
- $H \to \tau\tau$ uses a fake factor method covering all tau fakes from W/Z+jets *and* multijet
- instead of estimating multijet independently with isolation fake factors

Fake factors more quark-like at high-p_T



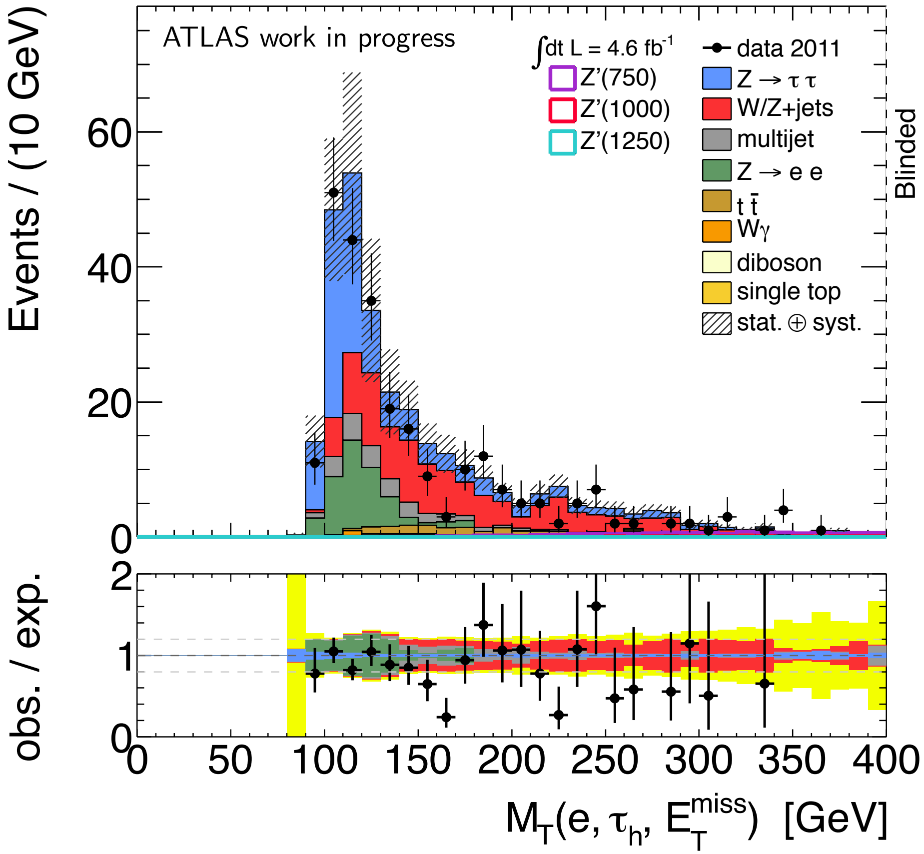
J. Gallicchio, M. Schwartz. "Pure Samples of Quark and Gluon Jets at the LHC". arXiv:1104.1175



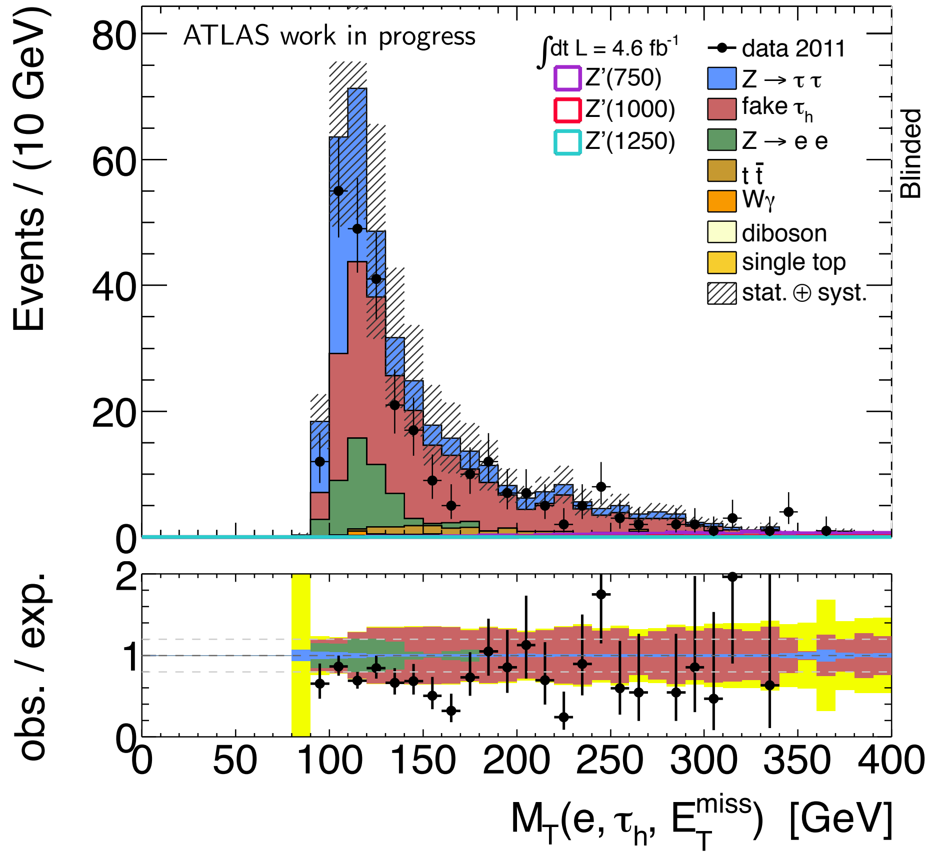
our support note: ATLAS-CONF-2012-067

Cross-check: single fake factor method

double fake factor



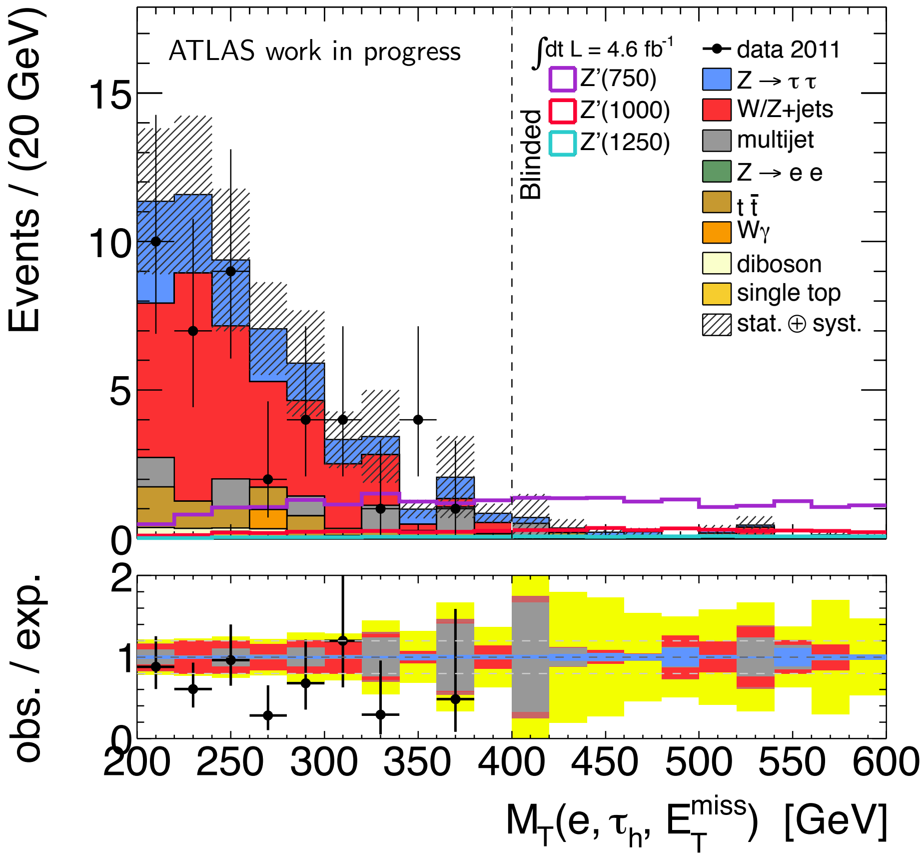
single fake factor



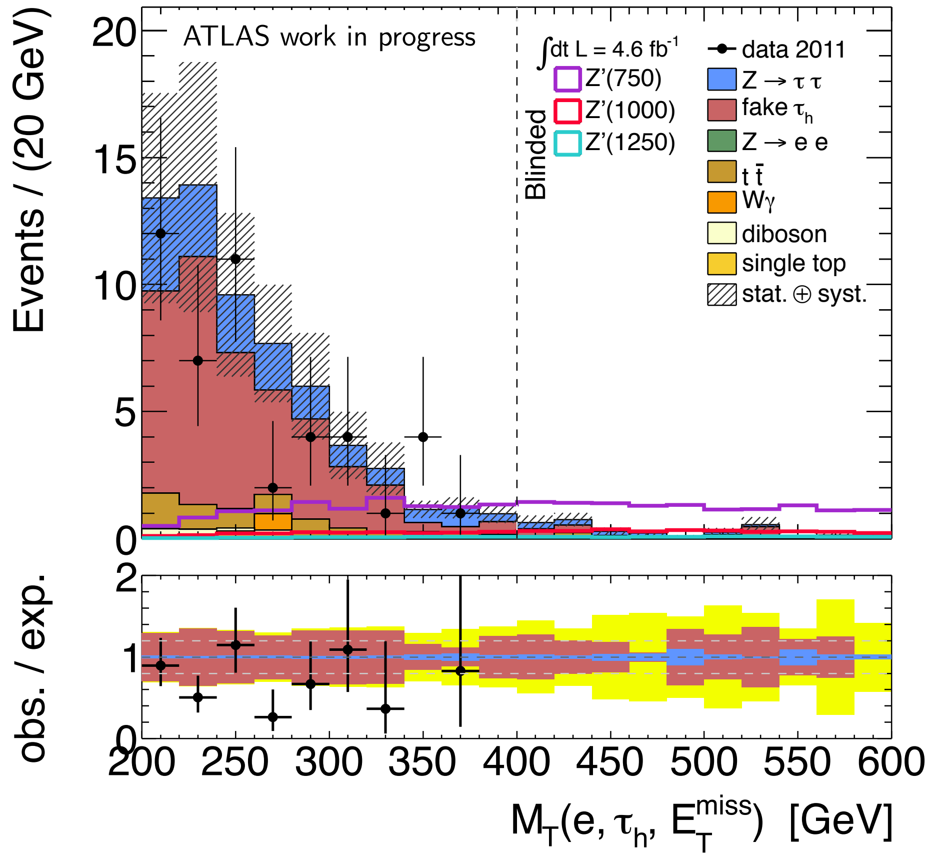
- as expected, the single fake factor method over-estimates in regions where the multijet contamination is large

Cross-check: single fake factor method

double fake factor



single fake factor

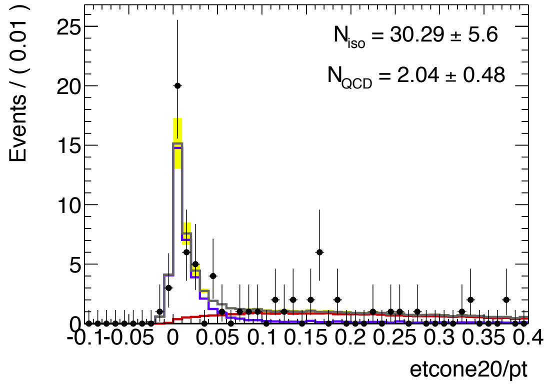


	double fake factor			single fake factor
	W/Z+jets	multijet	total	fake τ_h
$M_T > 400$ GeV	0.8(6)	0.3(3)	1.1(4)	1.3(4)
$M_T > 500$ GeV	0.8(4)	< 0.1	0.8(4)	0.9(4)

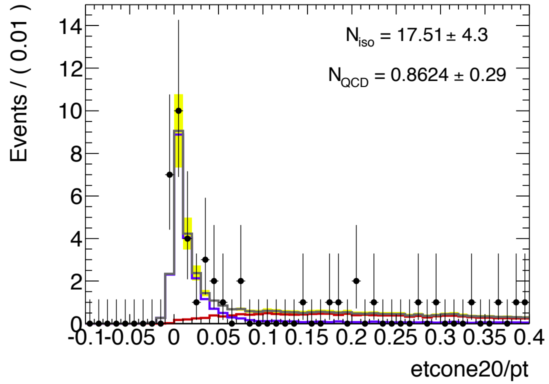
Lepton isolation fitting

- Used to cross-check the fake-factor QCD estimate.
- Similar in concept: predict normalization from the rate of non-isolated leptons.
- Fit calorimeter isolation with data-driven templates from QCD and W+jet control regions. Do fits separately in 6 p_T bins:

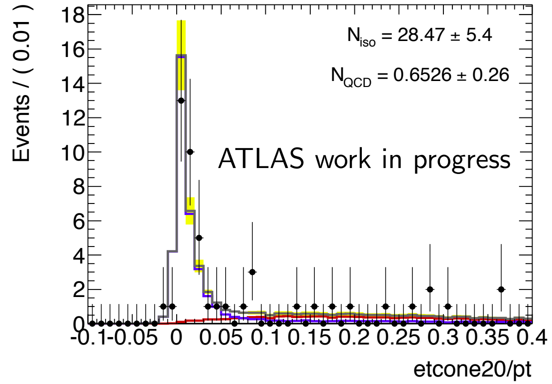
p_T = 25-27.5 GeV



p_T = 27.5-30 GeV



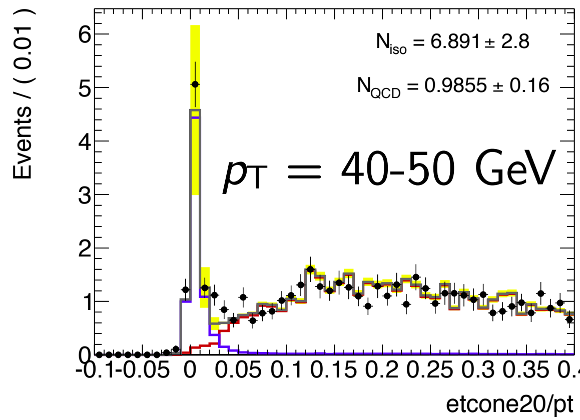
p_T = 30-35 GeV



Agrees with the lepton isolation fake-factor method:

	fake-factor	fitting
baseline	6.3(6)	4.3(4)
$p_T(\tau_h) > 80$ GeV	2.7(3)	1.8(2)
$p_T(\mu) > 40$ GeV	0.29(8)	0.31(8)
$m_{eff}(\mu, \tau_h, E_T^{miss}) > 400$ GeV	0.05(3)	0.05(3)

Fits work in failing tau ID region as well:



Control plots

Both W+jet and multijet are modeled with data.

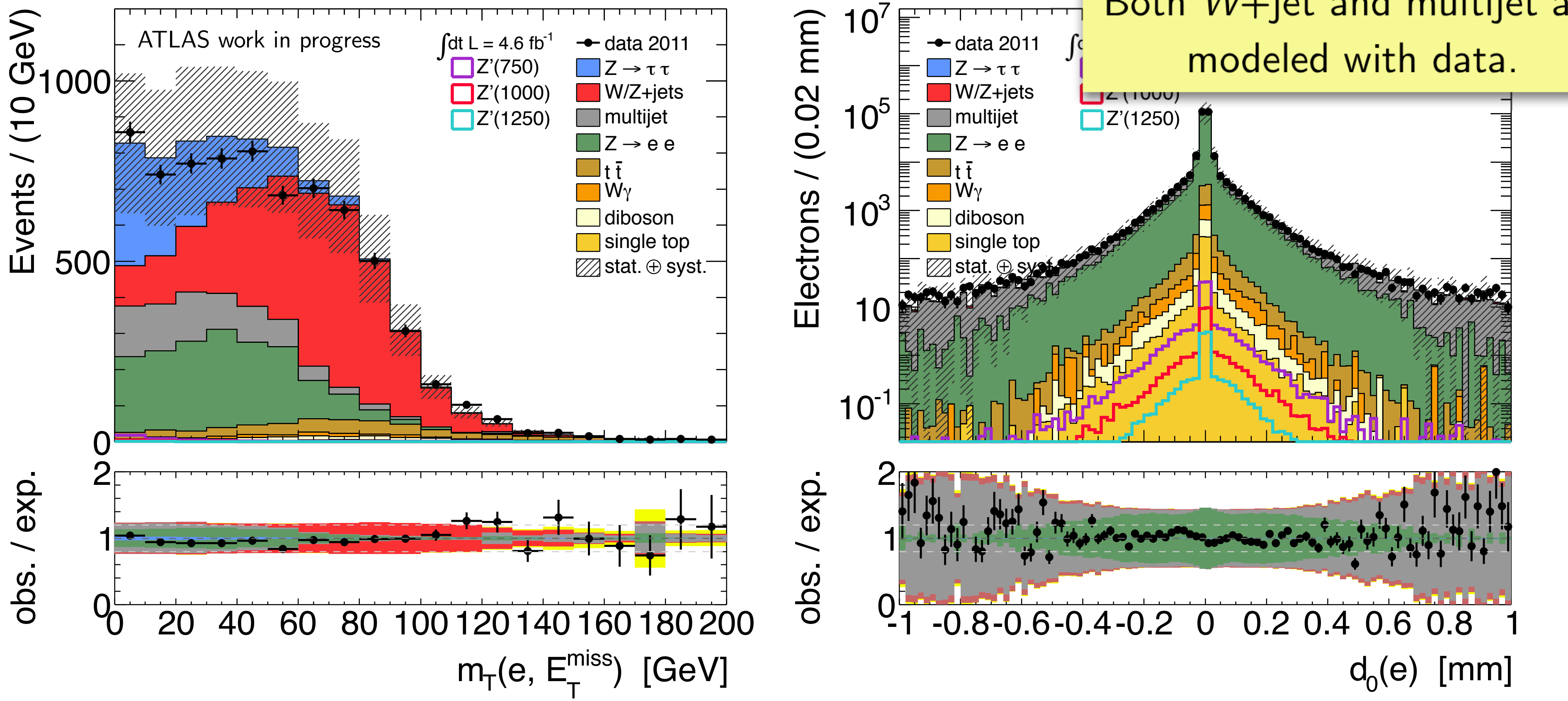


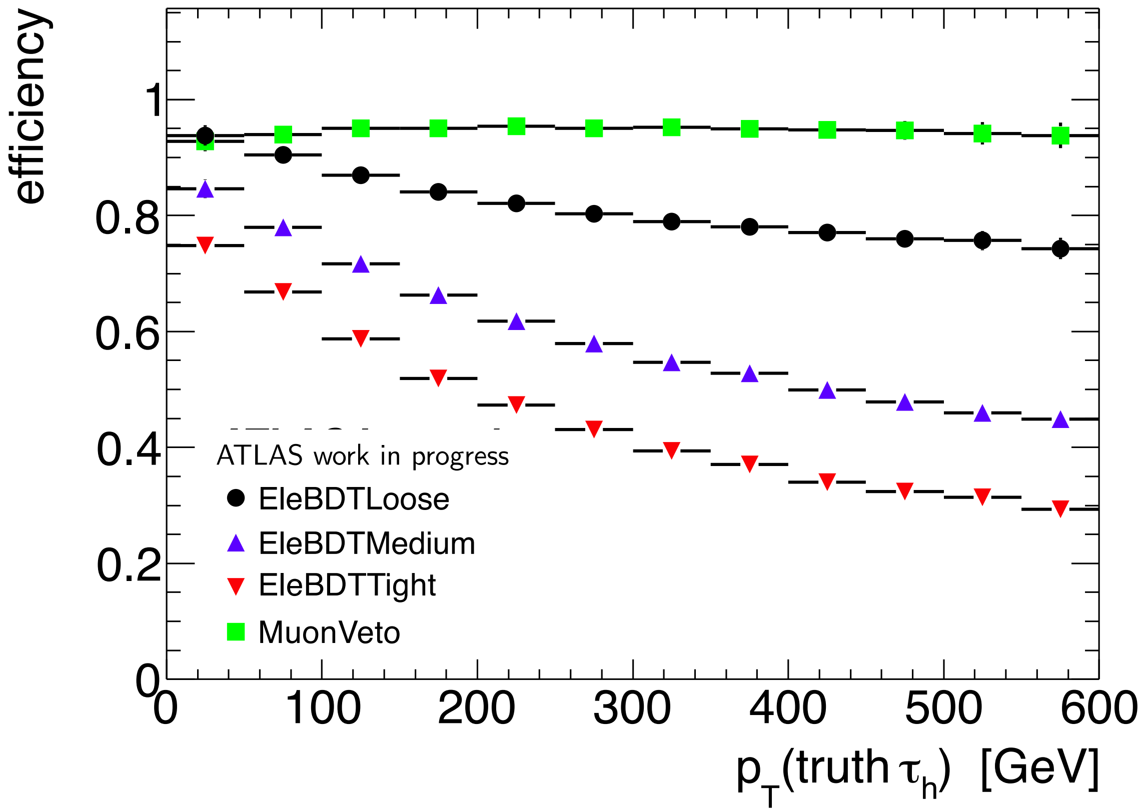
Figure 28: (left) The distribution of the transverse mass of the combination of the selected electron and the $E_{\mathrm{T}}^{\mathrm{miss}}$, $m_{\mathrm{T}}(e, E_{\mathrm{T}}^{\mathrm{miss}})$. in events with exactly one selected electron, no additional preselected electrons or muons, and exactly one selected 1-prong tau. (right) The distribution of the electron impact parameter, d_0, in events with exactly one selected electron, no additional preselected electrons or muons, and exactly one 1-prong tau candidate (without ID).

Electron veto

from tau WG Z→ee tag-and-probe with 2.6/fb from 2011

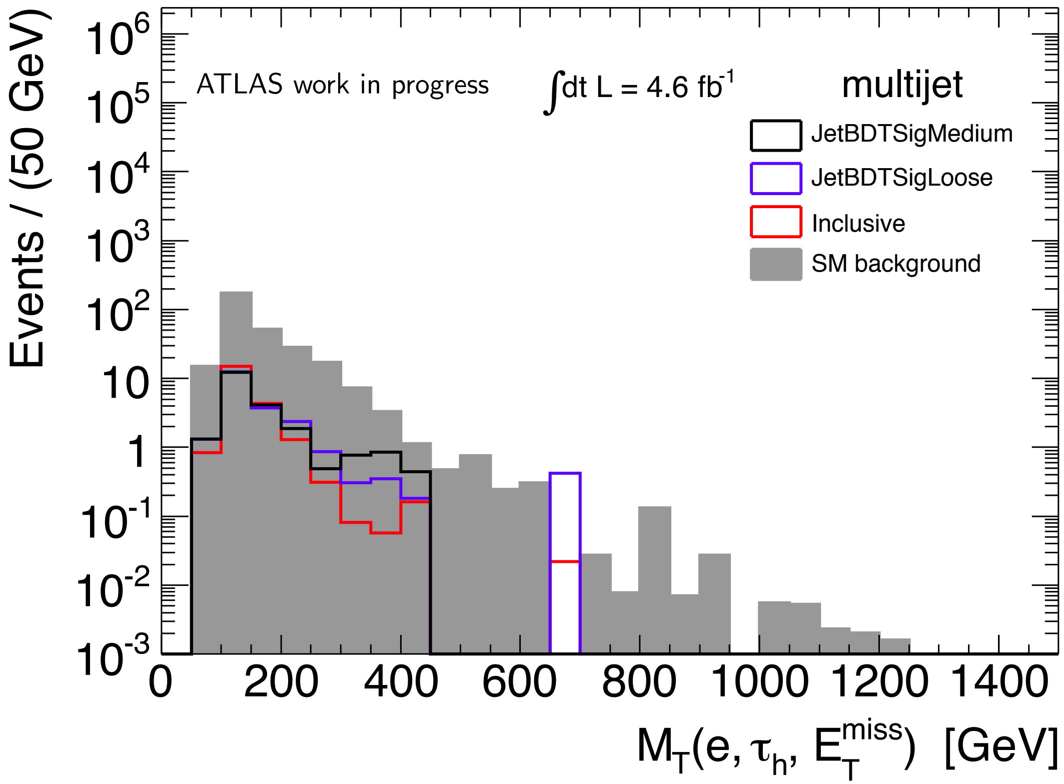
	$|\eta| < 1.37$	$1.37 < |\eta| < 1.52$	$1.52 < |\eta| < 2.0$	$2.0 < |\eta|$
BDT medium e-veto	1.64(0.81)	1.0(1.0)	0.71(0.63)	2.90(1.42)
BDT loose e-veto	1.21(0.30)	0.96(0.46)	0.59(0.21)	1.76(0.55)

- above are the scale factors and errors for the EleBDTMedium e-veto
- These SF have huge ≈100% uncertainties, and we've also previously shown that this veto is inefficient at high p_T.

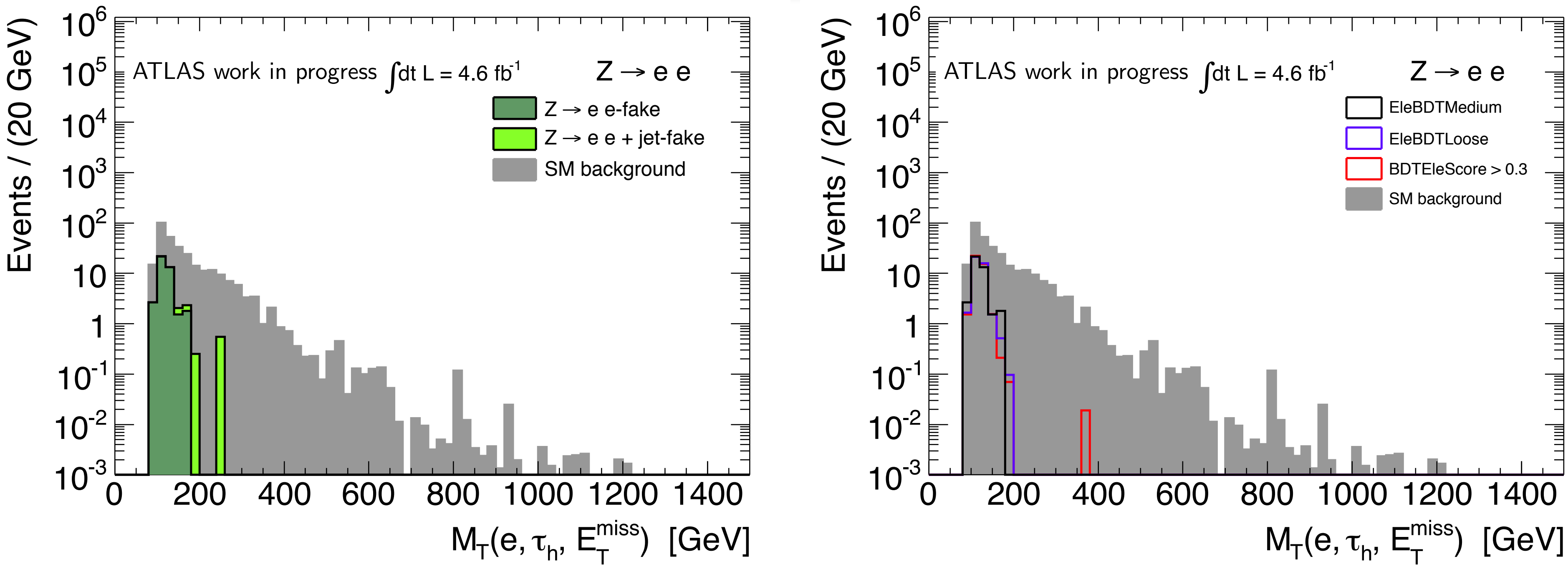


Multijets background

- Loosened tau ID requirement in the anti-isolated lepton + tau data sample used to model the multijet background.
- Sample with tau ID is already multijet dominated. Loosening tau ID improves stats.
- Shapes are statistically consistent. Inclusive shape scaled to the prediction with medium tau ID to give the estimate in the tail.



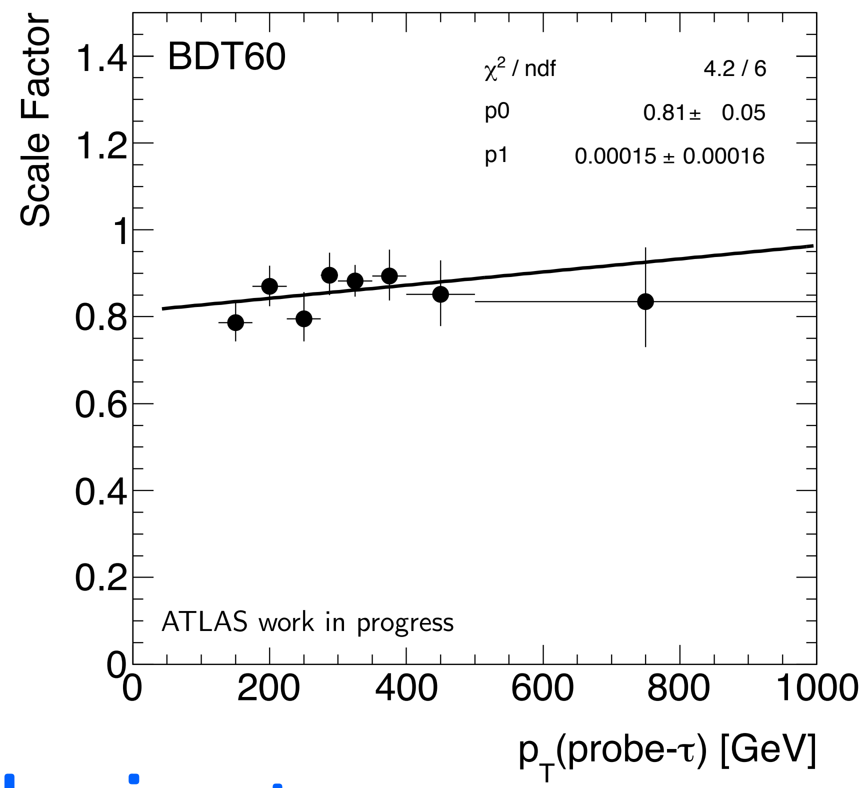
$Z \to ee$ + jets background



- Categorize Monte Carlo events by electron or jet faking tau.
- Loosen electron veto in Monte Carlo sample matched to electron fakes.
- Shapes are consistent, and only driven by the $Z \to ee$ kinematics.
- $Z \to ee$ with e-faking tau is negligible
- $Z \to ee$ + jet-fake covered with the data-driven W/Z+jet tau fake factor method.

High-p_T tau efficiency systematic

- The dominant systematic uncertainty for the Z' signal and the $Z \to \tau\tau$ background.
- Low p_T uncertainty of 4% taken from the Tau WG blessed $Z \to \tau\tau$ tag-and-probe.
- No high-p_T control sample of true taus.
- Assume mis-modeling comes from either:
  1. tau kinematics (TAUOLA)
  2. detector response to high-p_T pions ←**dominant**



- Instead of using true taus, measure the trend in the scale factor for fakes from dijet events.

> - $p_T \leq 100$ GeV: $\Delta\varepsilon = 4\%$ (taken from the $Z \to \tau\tau$ measurement)
> - $p_T > 100$ GeV: $\Delta\varepsilon = 4 + 0.016(p_T - 100)\%$, with p_T in GeV (taken from the dijets measurement).

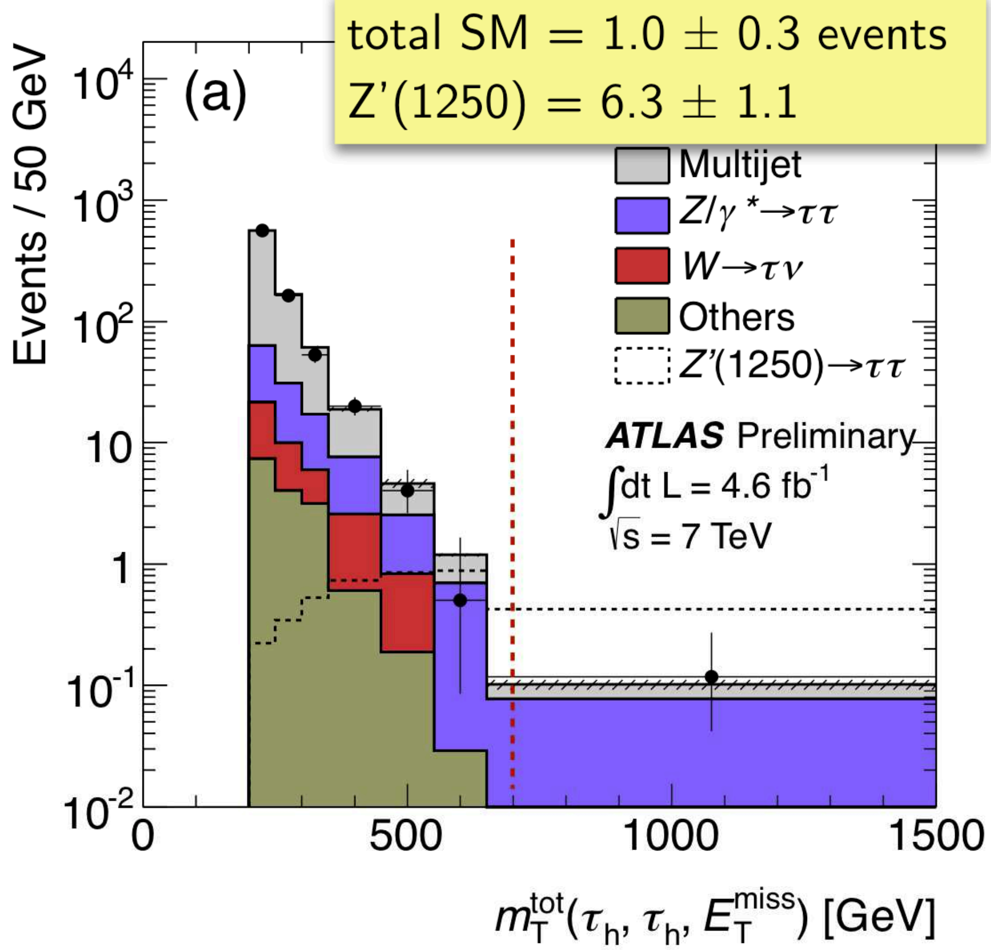
$Z' \rightarrow \tau\tau \rightarrow \tau_h\tau_h$

- New gauge bosons predicted in many GUTs with additional U(1).
- Best limit on $m(Z' \rightarrow ee/\mu\mu) > 2.3$ TeV from CMS [arxiv:1206.1849].
- Important to test the couplings to all lepton flavors.

Event selection

- 2 BDT loose 1 or 3-prong taus with $p_T(\tau_h) > 50$ GeV
- opposite sign
- $|\Delta\phi(e, \tau_h)| > 2.7$
- $m_T(\tau_h, \tau_h, E_T^{miss}) > 700$ GeV

total SM = 1.0 ± 0.3 events
Z'(1250) = 6.3 ± 1.1



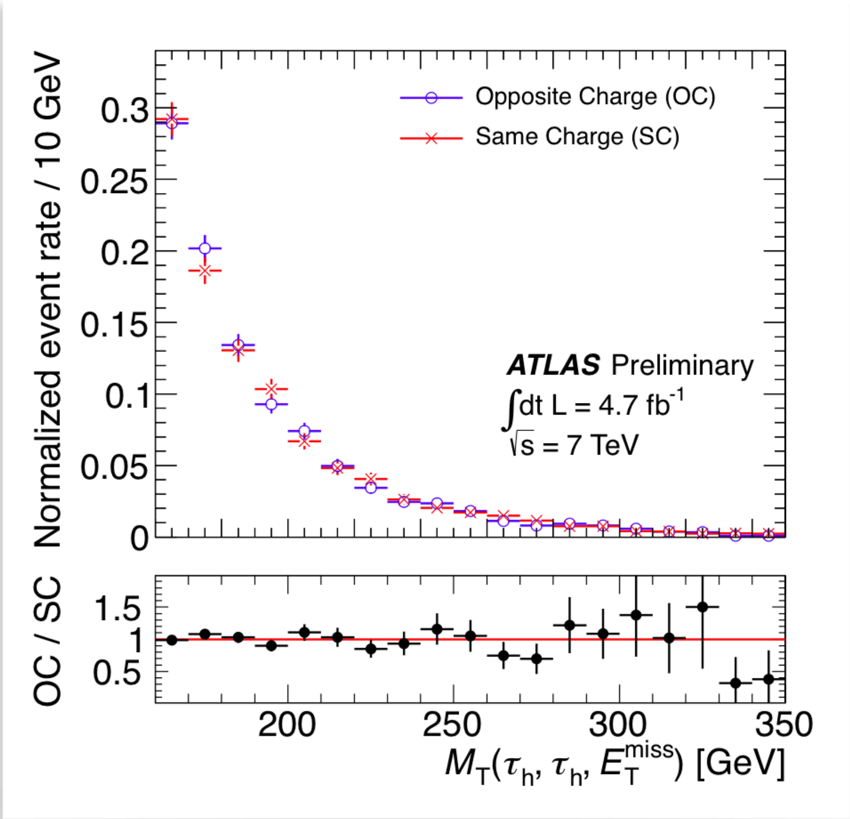
- Tau ID efficiency uncert. ≈ 11% on the signal. (4% from $Z \rightarrow \tau\tau$ tag-and-probe)
- Jet/tau energy scale uncert. ≈ +22/-11%
- Multijet modeled by fitting the shape of the SS data. uncert. ≈ +21/-11%

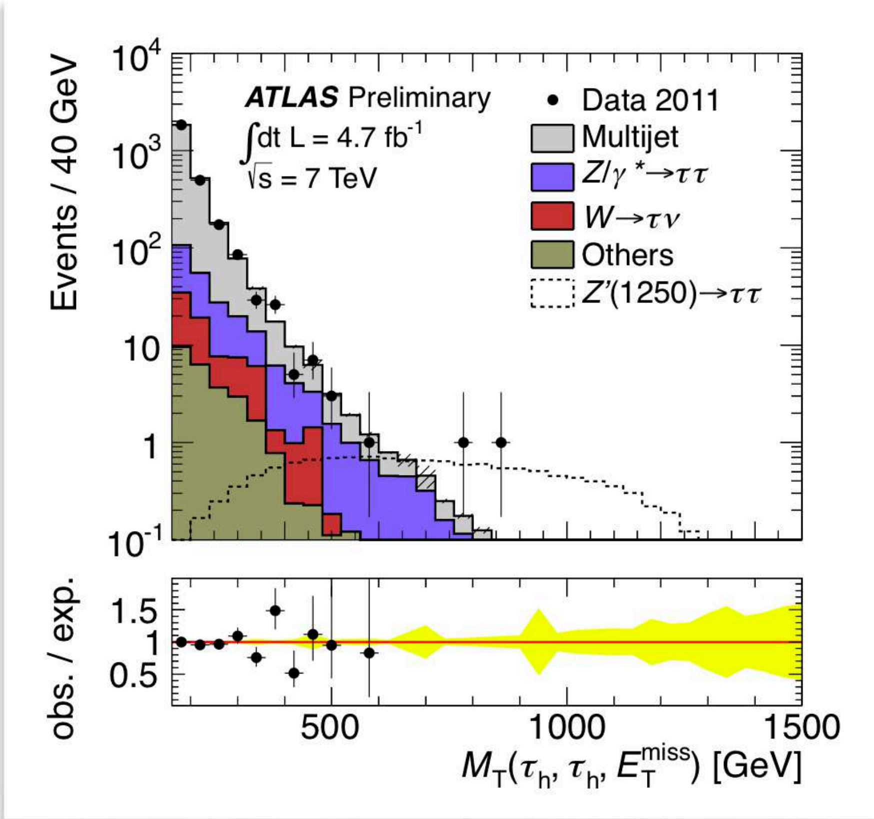
[ATLAS-CONF-2012-067]

Data-driven multijet

Fit same-sign (SS) data with dijet function:

$$f(M_T|p_0, p_1, p_2) = p_0 \cdot M_T{}^{p_1 + p_2 \log M_T}.$$

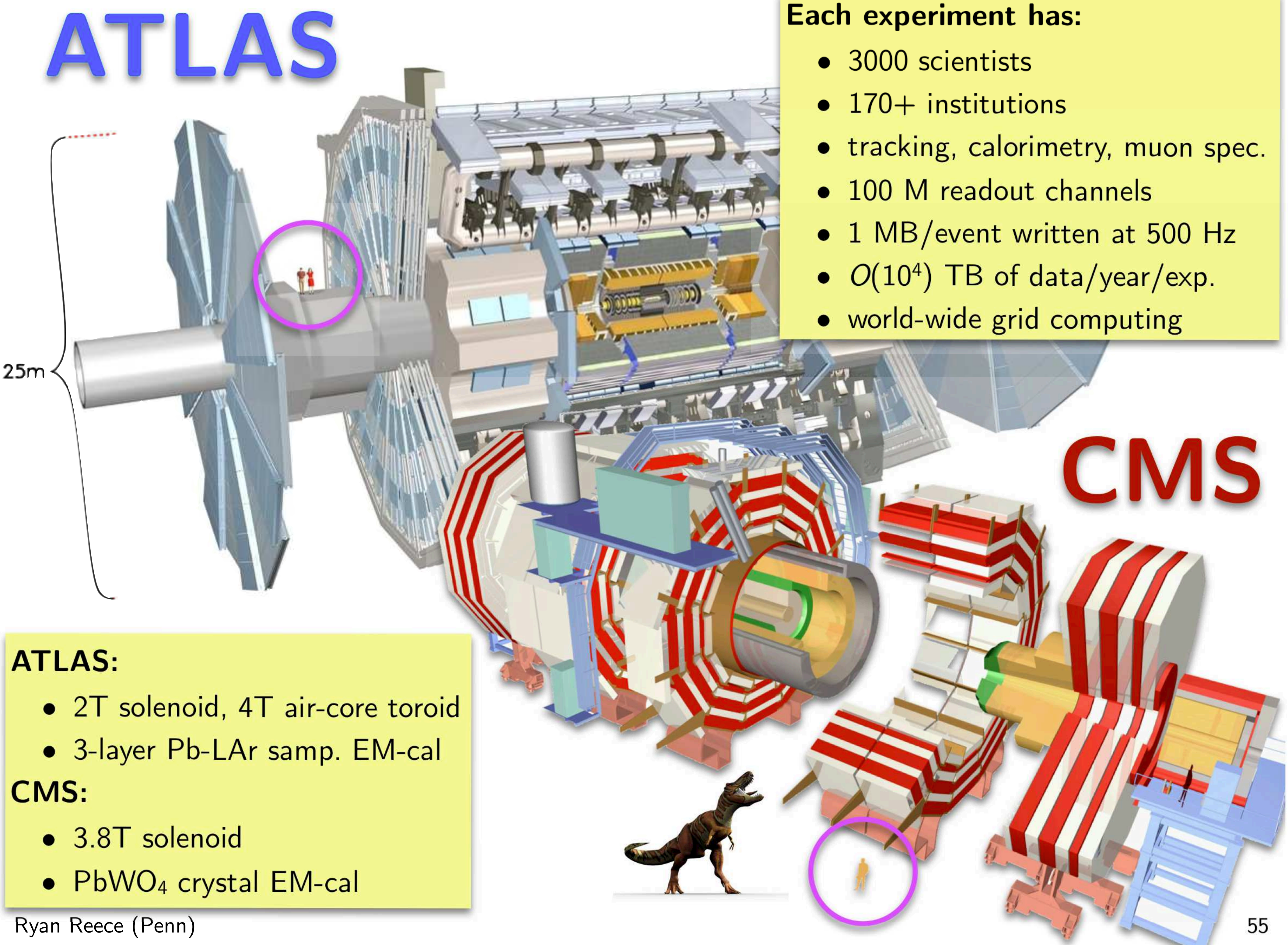




- OS/SS shapes agree well
- normalize in OS sideband with $200 < M_T < 250$ GeV

[ATLAS-CONF-2012-067]

The LHC,
ATLAS, and CMS



Each experiment has:

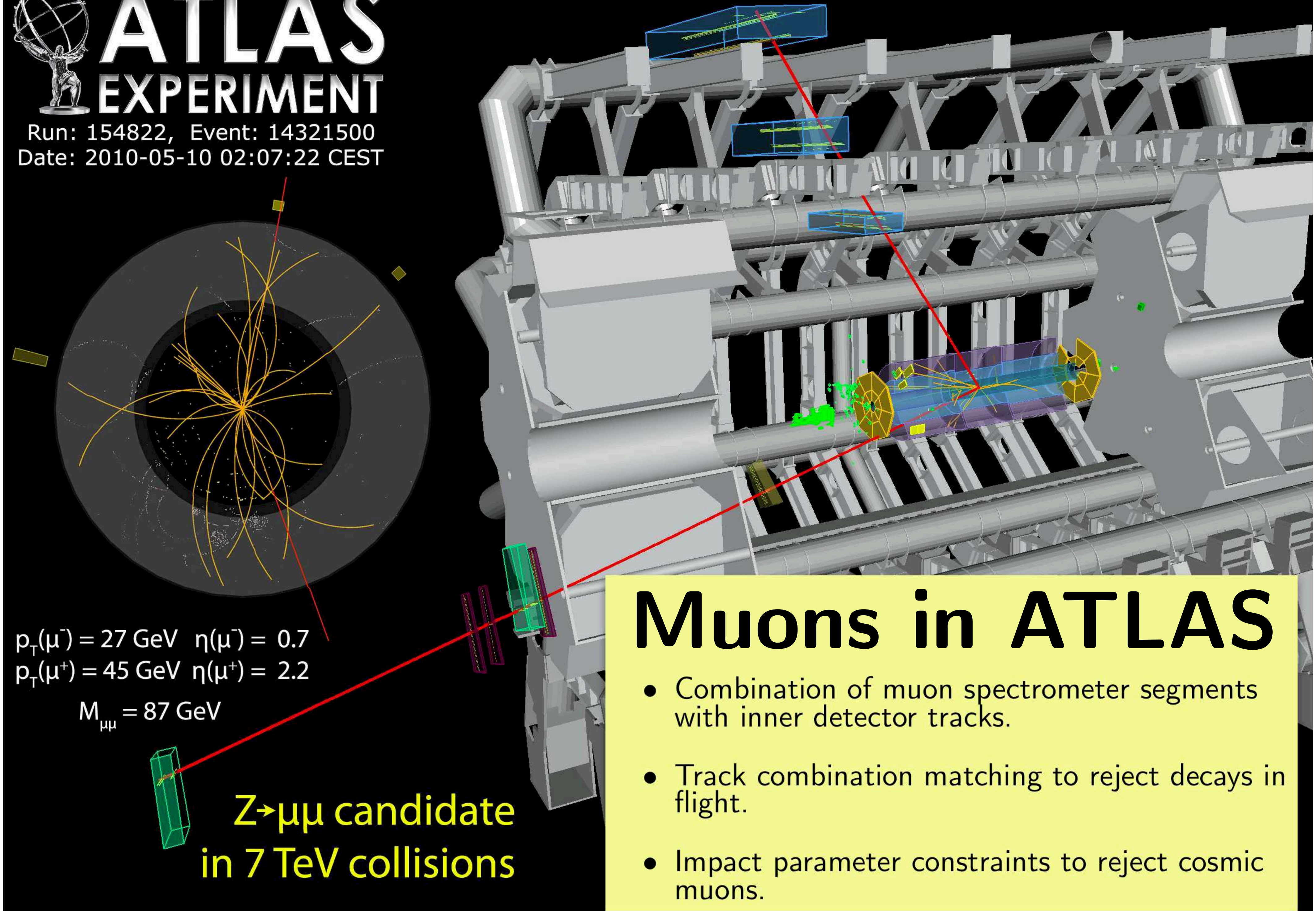
- 3000 scientists
- 170+ institutions
- tracking, calorimetry, muon spec.
- 100 M readout channels
- 1 MB/event written at 500 Hz
- $O(10^4)$ TB of data/year/exp.
- world-wide grid computing

ATLAS:

- 2T solenoid, 4T air-core toroid
- 3-layer Pb-LAr samp. EM-cal

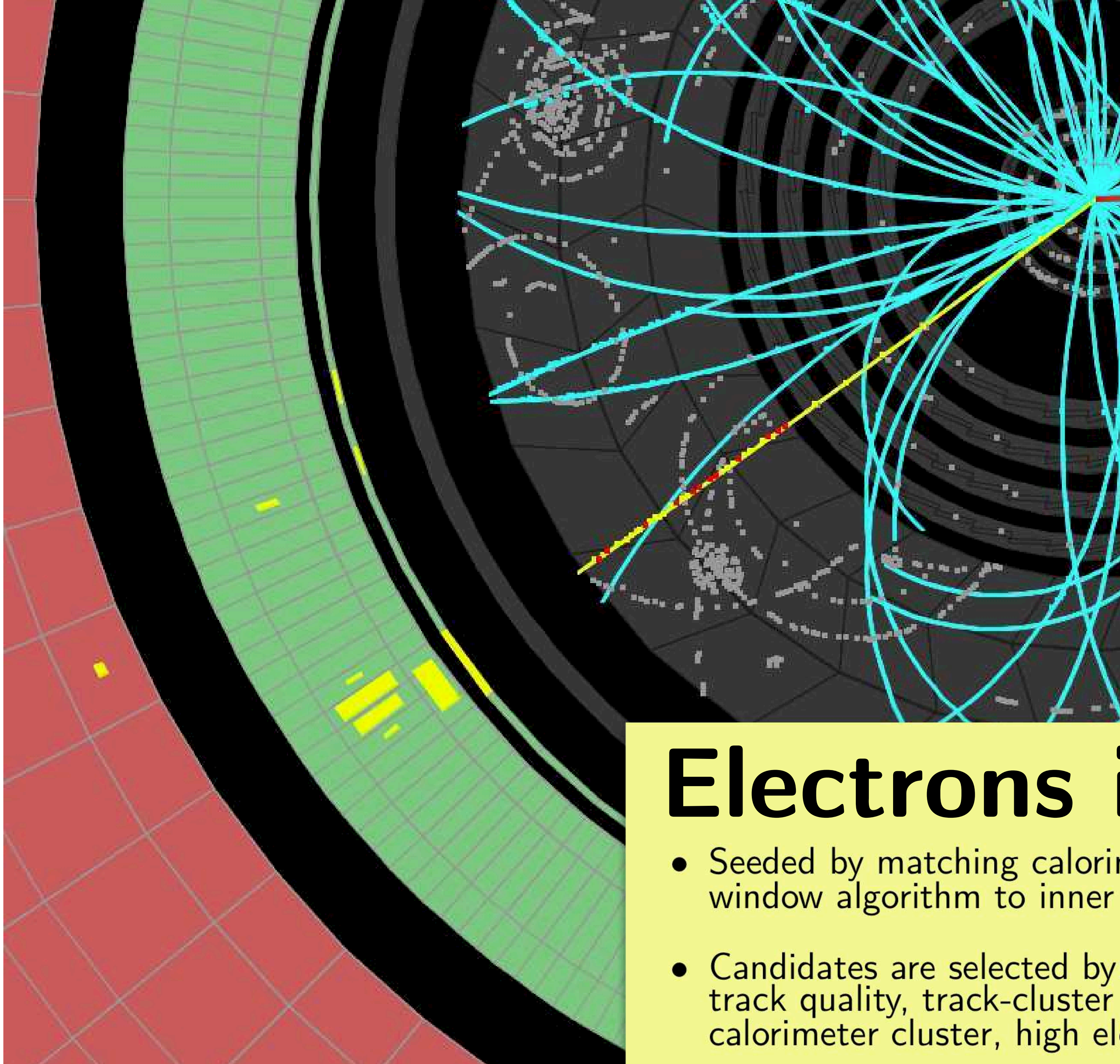
CMS:

- 3.8T solenoid
- $PbWO_4$ crystal EM-cal



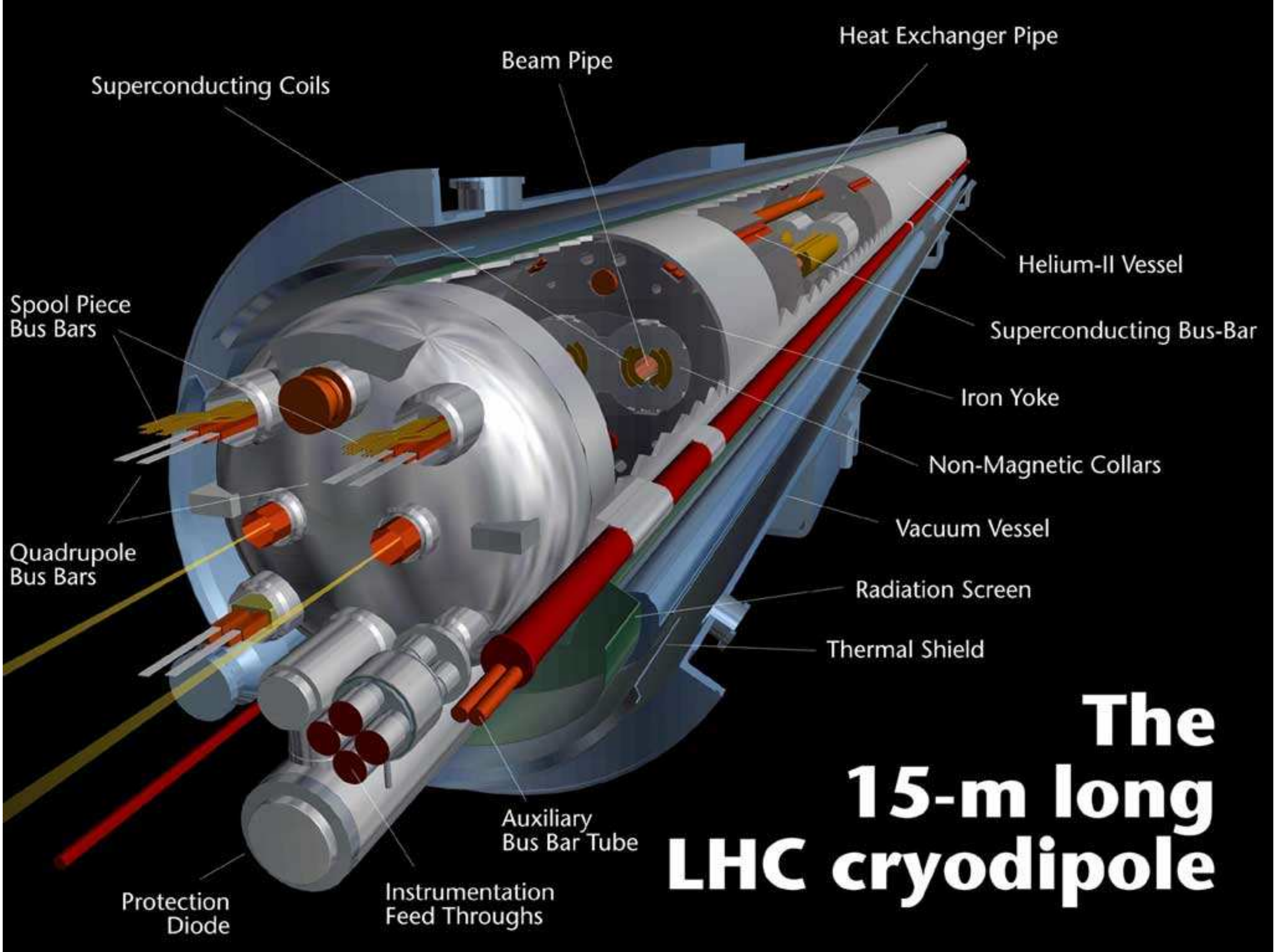
Muons in ATLAS

- Combination of muon spectrometer segments with inner detector tracks.
- Track combination matching to reject decays in flight.
- Impact parameter constraints to reject cosmic muons.

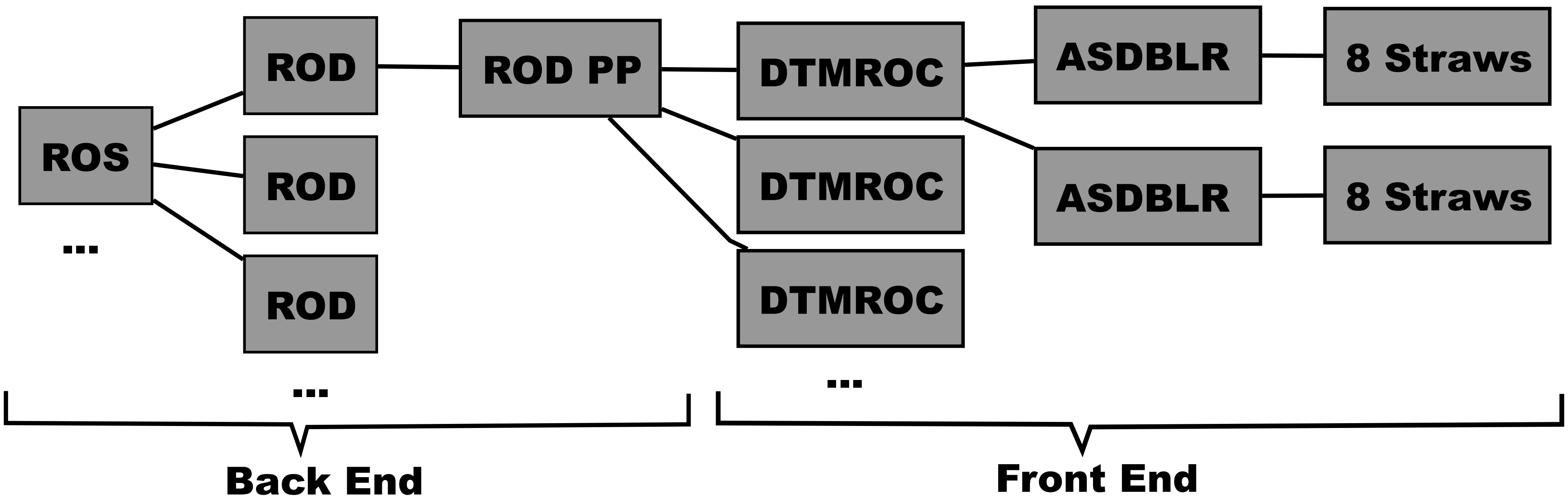


Electrons in ATLAS

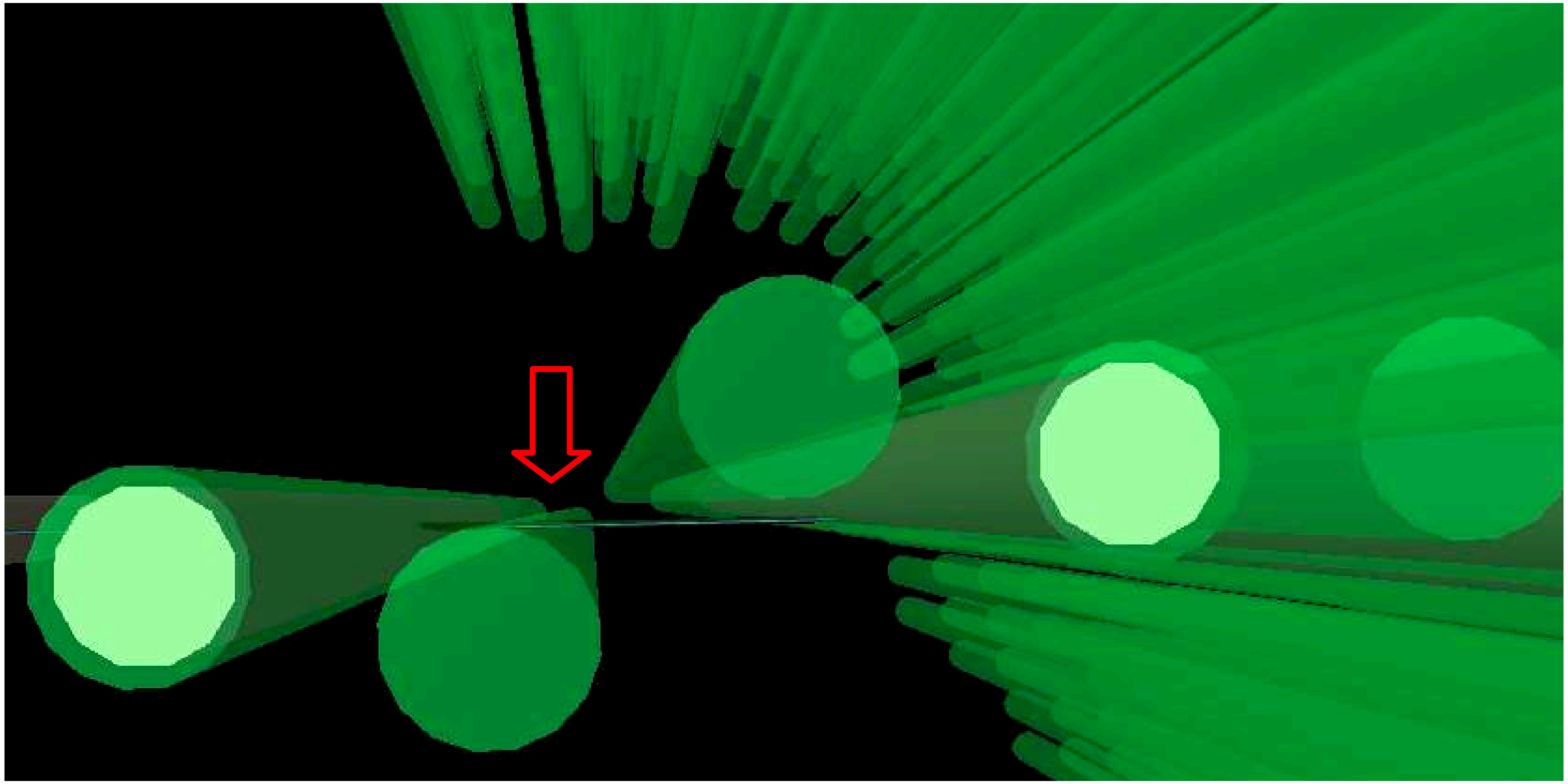
- Seeded by matching calorimeter clusters from a sliding-window algorithm to inner detector tracks.
- Candidates are selected by:
  track quality, track-cluster matching, narrow calorimeter cluster, high electromagnetic fraction
- Tight candidates have cuts on E/p and high thresholds hits from the transition radiation in the TRT.


Heat Exchanger Pipe
Beam Pipe
Superconducting Coils
Helium-II Vessel
Spool Piece
Bus Bars
Superconducting Bus-Bar
Iron Yoke
Non-Magnetic Collars
Vacuum Vessel
Quadrupole
Bus Bars
Radiation Screen
Thermal Shield
Auxiliary
Bus Bar Tube
Protection
Diode
Instrumentation
Feed Throughs
The
15-m long
LHC cryodipole

TRT Read-out


ROS
...
ROD
ROD
ROD
...
ROD PP
DTMROC
DTMROC
DTMROC
...
ASDBLR
ASDBLR
8 Straws
8 Straws
Back End
Front End

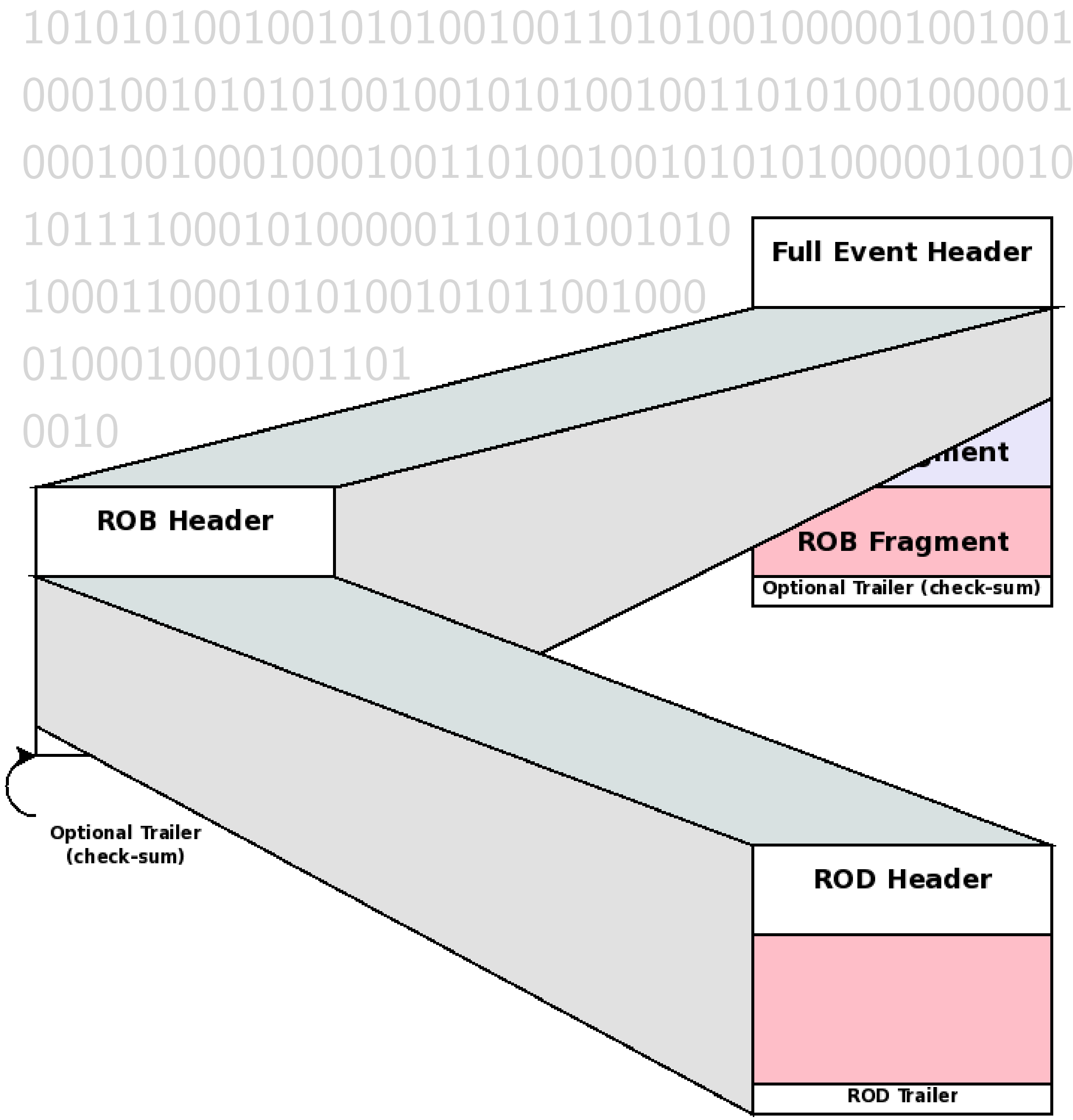
TRT hit efficiency



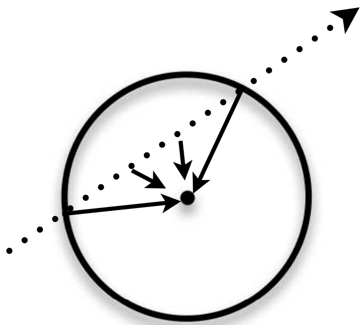
Parsing ATLAS raw data

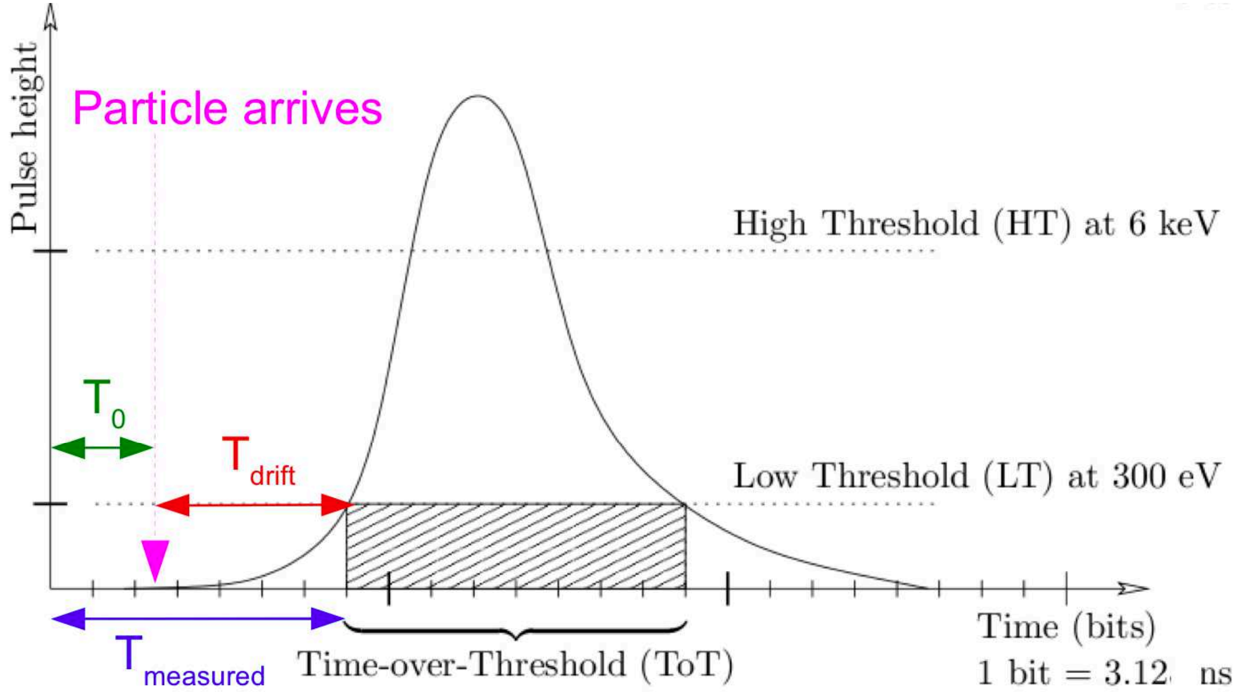
2007

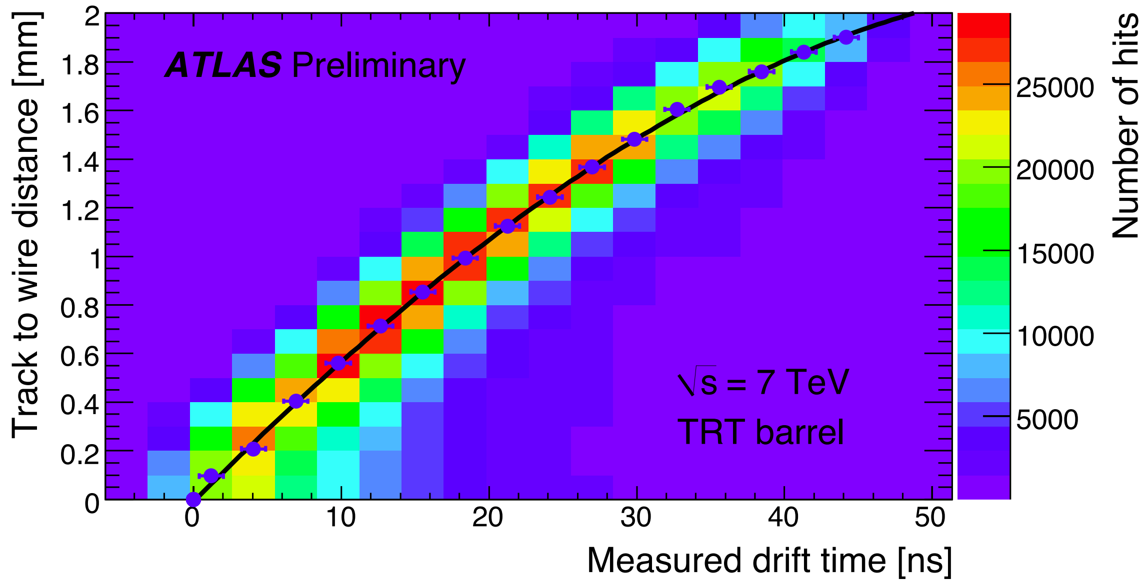
- Official libraries for parsing the ATLAS raw data were not finalized.
- Groups doing commissioning were still writing their own tools.
- My first software project was to write a library for parsing the ATLAS raw data for TRT commissioning purposes.
- Later, I used it to parse TRT threshold scans.



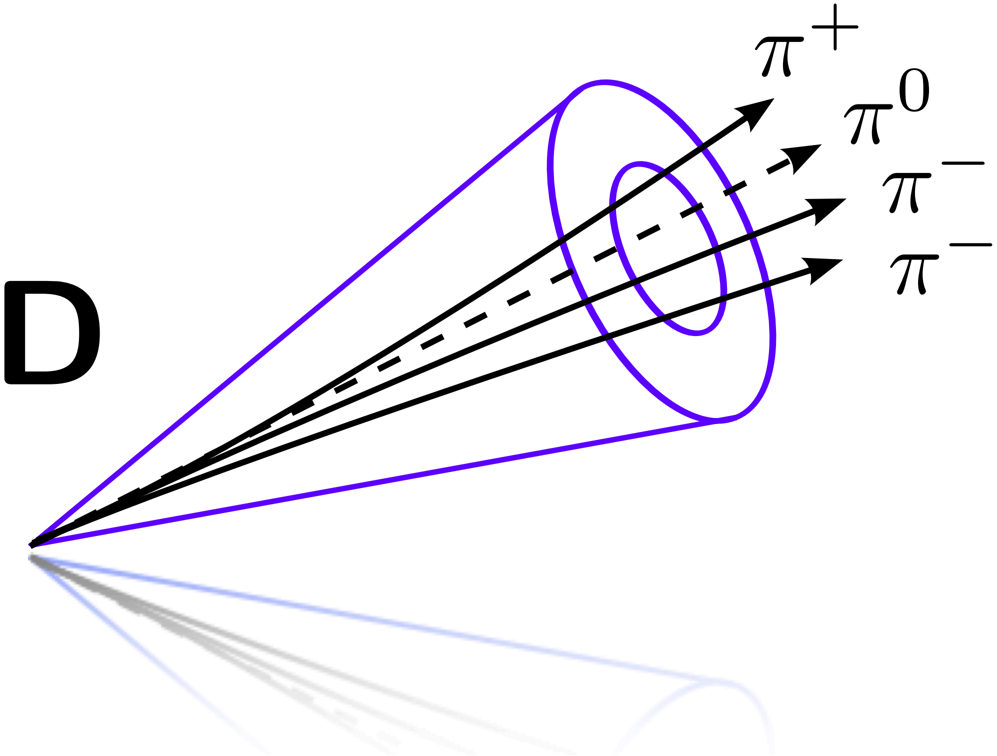
TRT Straw Hits





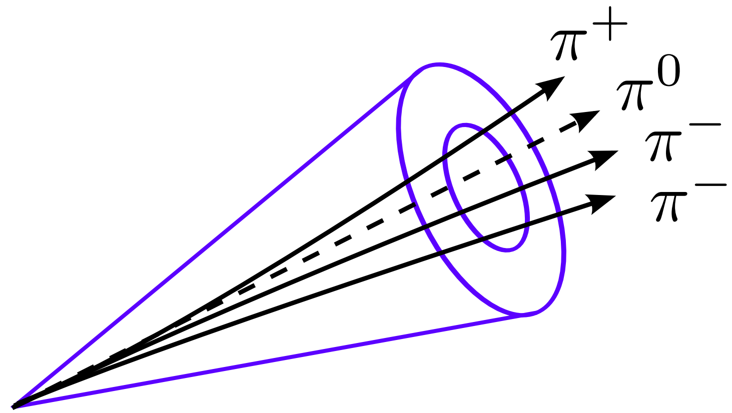


ATLAS Tau ID



Phenomenology of tau decays

$\tau^- \to$	$e^- \, \bar{\nu}_e \, \nu_\tau$	17.8%	leptonic 35.2%
	$\mu^- \, \bar{\nu}_\mu \, \nu_\tau$	17.4%	
	$\pi^- \, \pi^0 \, \nu_\tau$	25.5%	1 prong 49.5%
	$\pi^- \, \nu_\tau$	10.9%	
	$\pi^- \, 2\pi^0 \, \nu_\tau$	9.3%	
	$K^- \, (N\pi^0) \, (NK^0) \, \nu_\tau$	1.5%	
	$\pi^- \, 3\pi^0 \, \nu_\tau$	1.0%	
	$\pi^- \, \pi^- \, \pi^+ \, \nu_\tau$	9.0%	3 prong 15.2%
	$\pi^- \, \pi^- \, \pi^+ \, \pi^0 \, \nu_\tau$	4.6%	



Current tau identification variables

1. Core energy fraction*

$$f_{\rm core} = \frac{\sum_{\{\Delta R<0.1\}} E_{\rm T}^{\rm EM}({\rm cell})}{\sum_{\{\Delta R<0.2\}} E_{\rm T}^{\rm EM}({\rm cell})}$$

2. Leading track momentum fraction*
3. Track radius $R_{\rm track} = \frac{\sum_{\{\Delta R<0.4\}} p_{\rm T}({\rm track})\, \Delta R({\rm track,\ jet})}{\sum_{\{\Delta R<0.4\}} p_{\rm T}({\rm track})}$
4. Number of isolation tracks $N_{\rm trk}^{0.2<\Delta R<0.4}$
5. Leading track impact parameter significance $S_{\rm lead\ track} = \frac{d_0}{\sigma_{d_0}}$
6. Transverse flight path significance $S_{\rm T}^{\rm flight} = \frac{L_{\rm T}^{\rm flight}}{\sigma_{L_{\rm T}^{\rm flight}}}$
7. Mass of track system
8. Maximum ΔR between jet-axis and core tracks

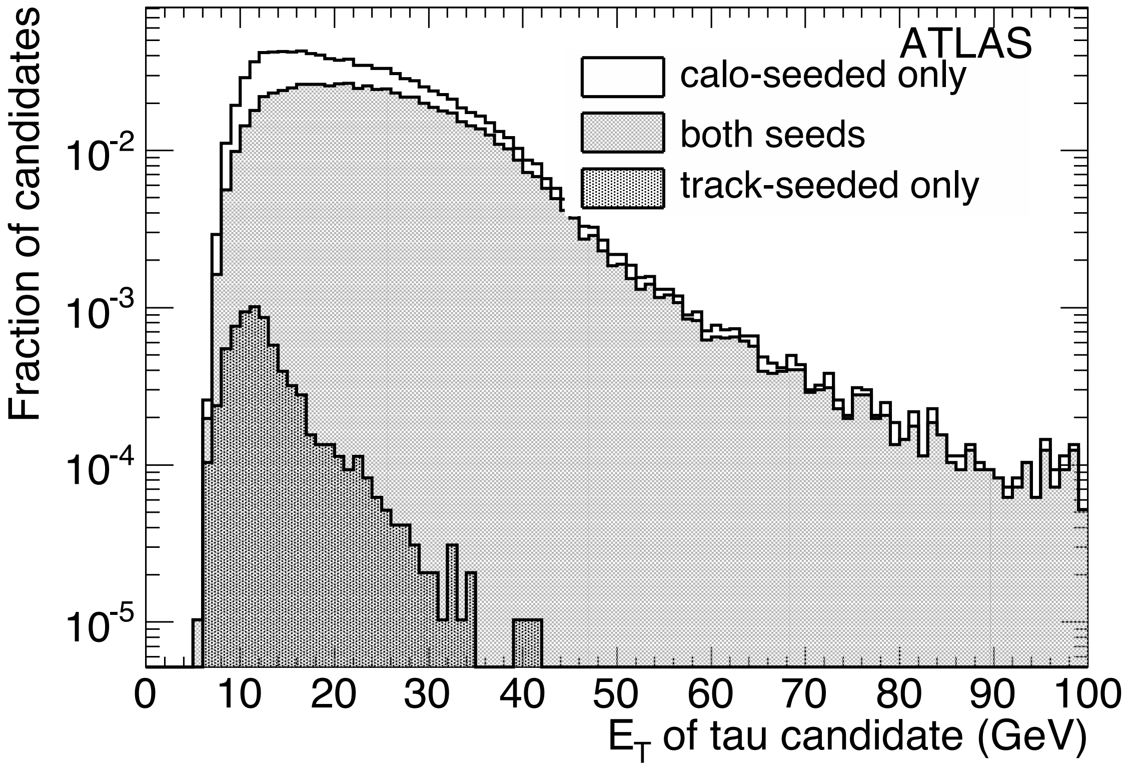
*has pile-up correction term linear in N(vertex)

Seeds of reconstruction

Once upon a time, there were two tau reconstruction algorithms. 2009

1. **tauRec** - seeded by $p_T > 10$ GeV anti-k_T 0.4 topo-jets. "calo-seeded"

2. **tau1p3p** - seeded by $p_T > 6$ GeV inner detector tracks. "track-seeded"



Since virtually all candidates have a calo-seed, we effectively merged the variable calculation of both algorithms, using only calo-seeds.

"Performance of the tau reconstruction and identification algorithm with 14.2.20 and mc08" [ATL-COM-PHYS-2009-229]

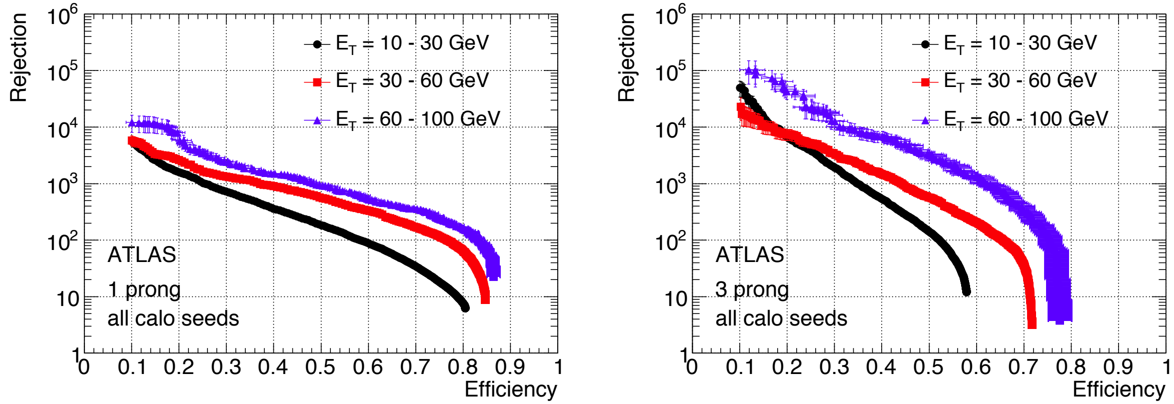
Early MV identification

2009

- **Jet-tau discrimination**

  Prefers narrow calorimeter jets, likelihood-based discriminant.

$$R_{\mathrm{EM}} = \frac{\sum_i^{\Delta R_i<0.4} E_{\mathrm{T},i}^{\mathrm{EM}} \Delta R_i}{\sum_i^{\Delta R_i<0.4} E_{\mathrm{T},i}^{\mathrm{EM}}},$$



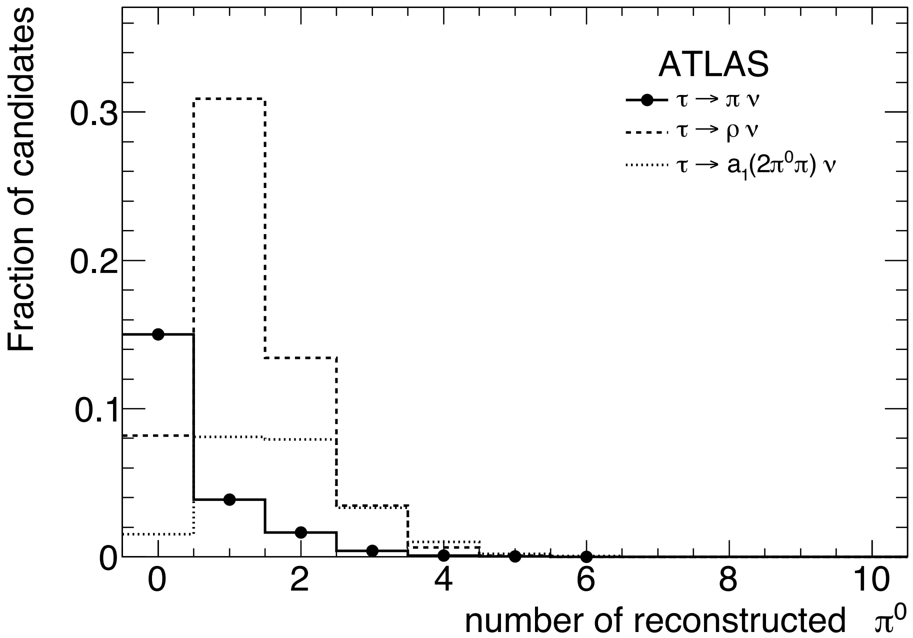
- **Electron-tau discrimination**

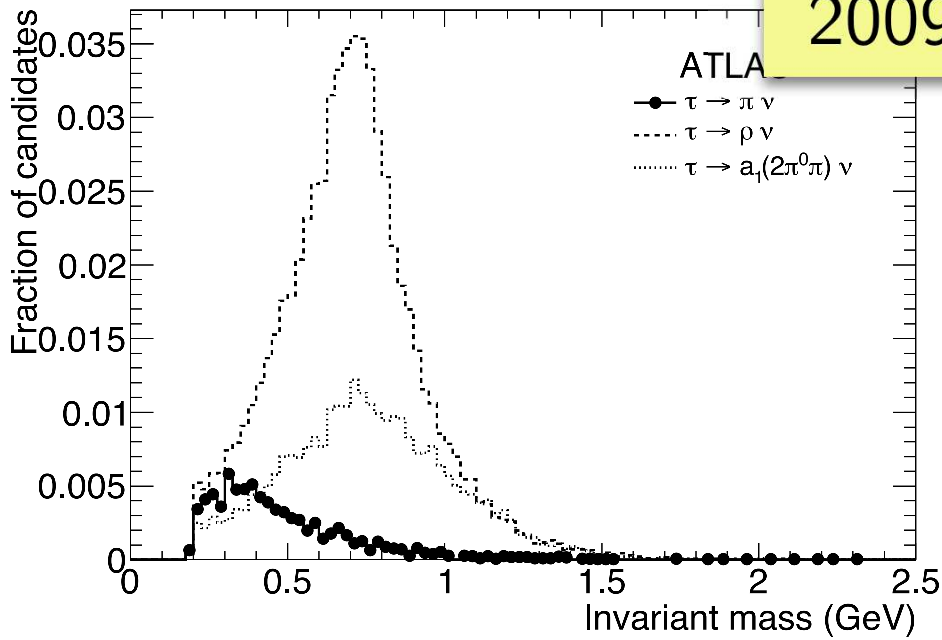
Candidate	IsEle(%) Overall	1P	3P	IsEle_eg(%) Overall	1P	3P
τ from W$\rightarrow\tau\nu$	93.2	92.7	95.3	99.8	99.8	99.8
τ form A$\rightarrow\tau\tau$	93.3	92.5	96.3	99.9	98.8	99.5
Electron form W$\rightarrow e\nu$	2.8	2.4	0.1	14.8	13.4	0.3
Electron form A$\rightarrow\tau\tau$	5.9	4.5	0.5	18.0	15.8	0.8

[ATL-COM-PHYS-2009-229]

Early sub-structure studies

2009

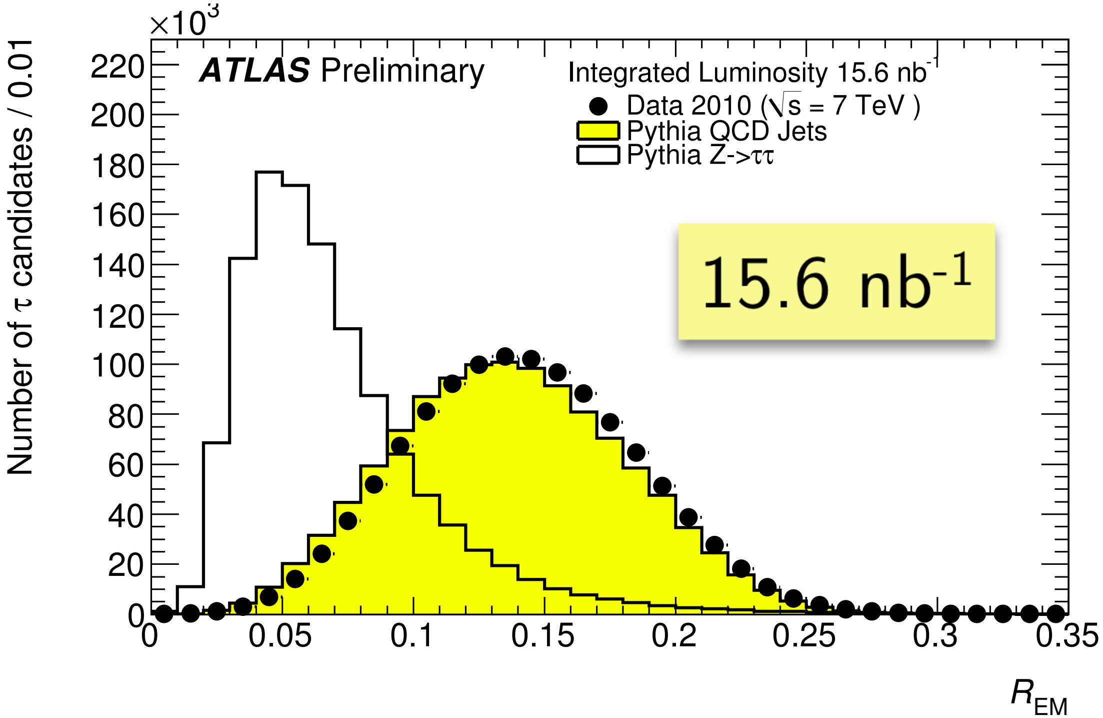


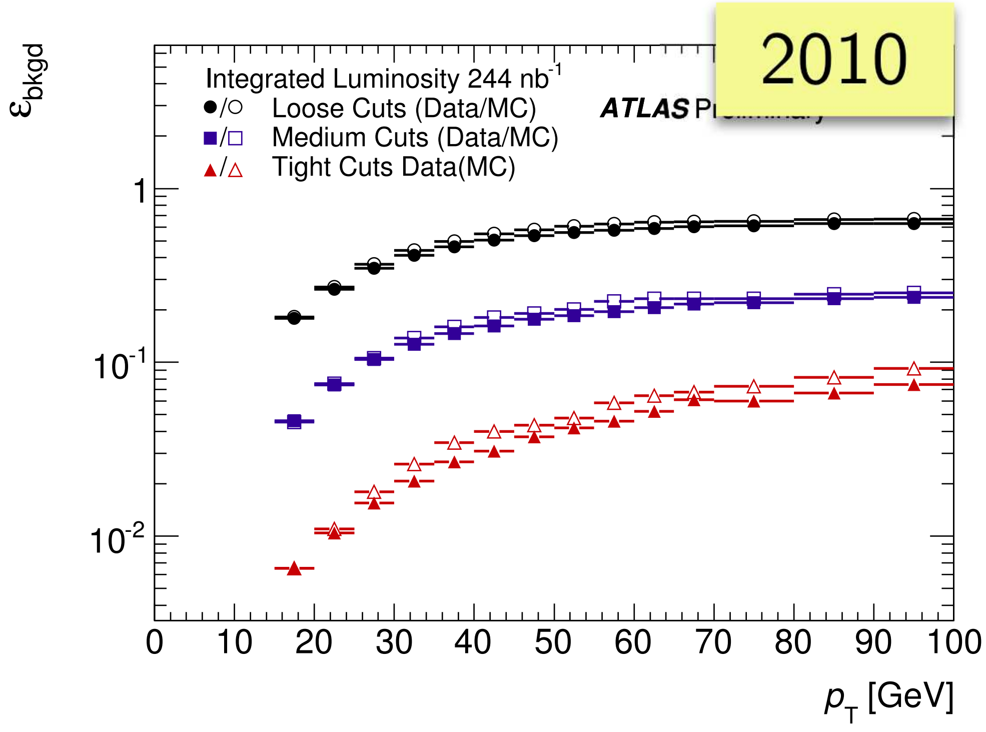


- Monte Carlo based substructure studies
- Cell-based shower-shape subtraction π^0 reconstruction.
- Still unvalidated with data.

[ATL-COM-PHYS-2009-229]

First data





- First comparisons of background distributions and the QCD fake-rate between data and Monte Carlo.
- Already see that MC over-estimates the jet fake-rate. $\Rightarrow k_W \approx 0.5$

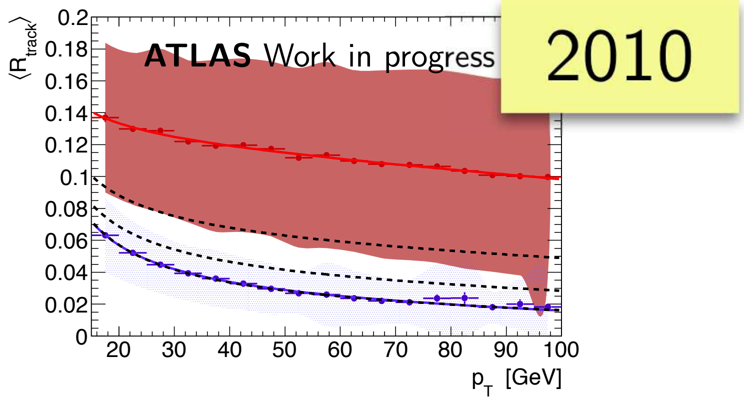
"Reconstruction of hadronic tau candidates in QCD events at ATLAS with 7 TeV pp collisions" [ATLAS-CONF-2010-059]

"Tau Reconstruction and Identification Performance in ATLAS" [ATLAS-CONF-2010-086]

Tau discriminants

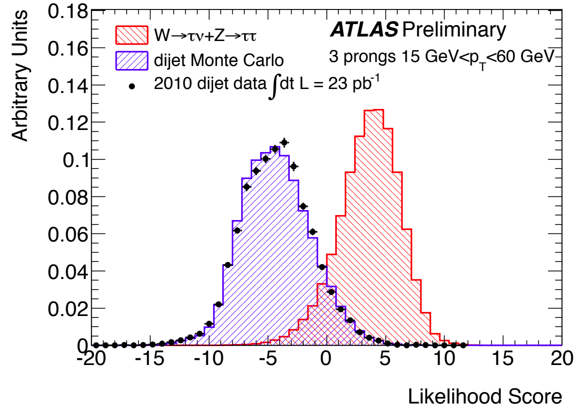
- **Cuts**

  p_{T}-parametrized cuts on R_{EM} and R_{track}, and a cut on f_{track}.

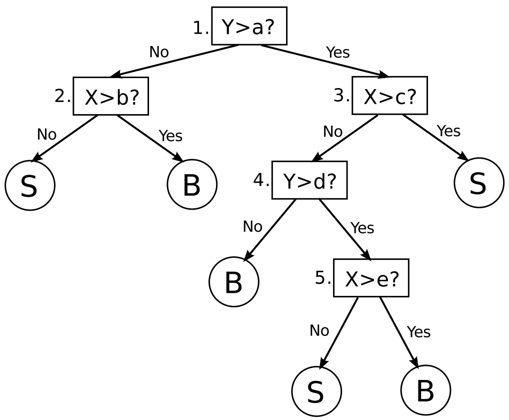


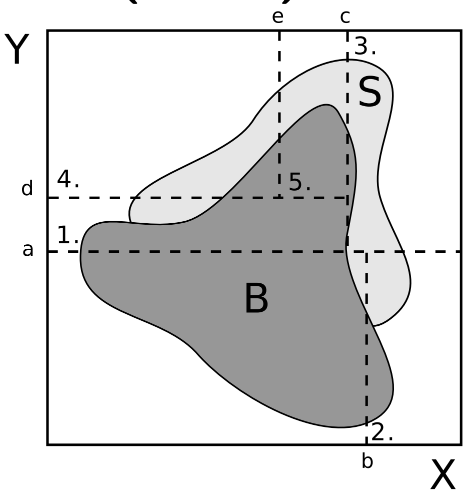
- **Projective likelihood**

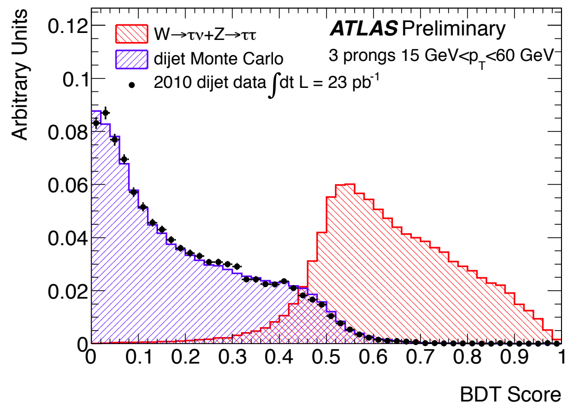
  $$d = \ln\left(\frac{L_S}{L_B}\right) = \sum_{i=1}^{N} \ln\left(\frac{p_i^S(x_i)}{p_i^B(x_i)}\right)$$



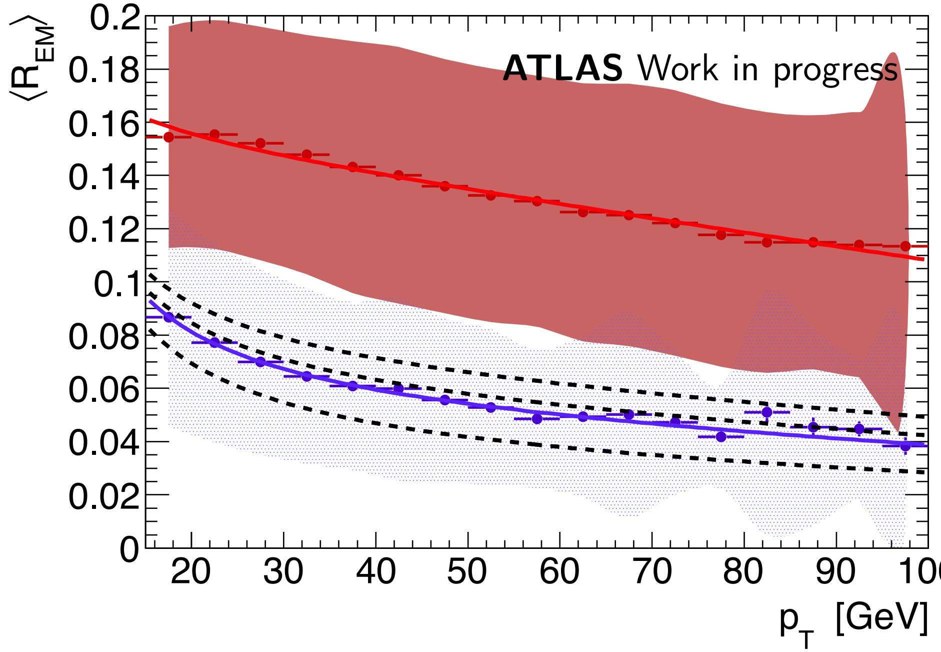
- **Boosted decision trees (BDT)**



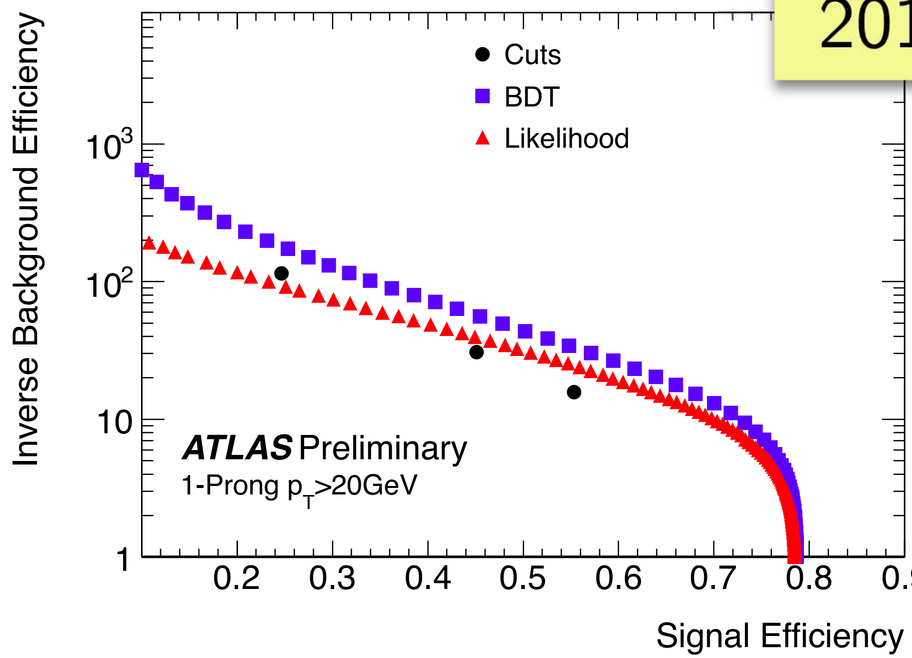




Maturing of discriminants



2010

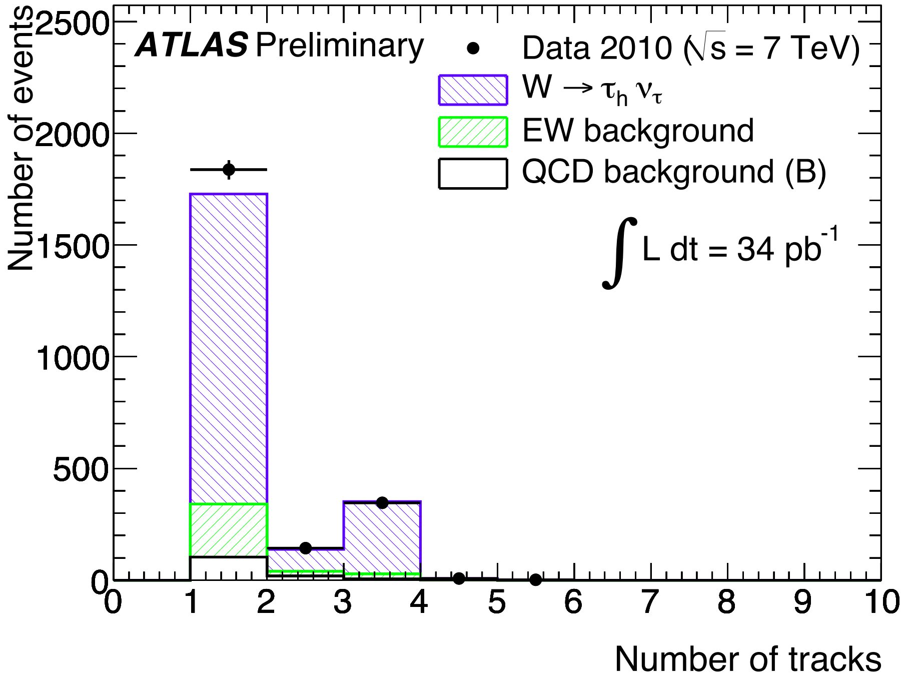


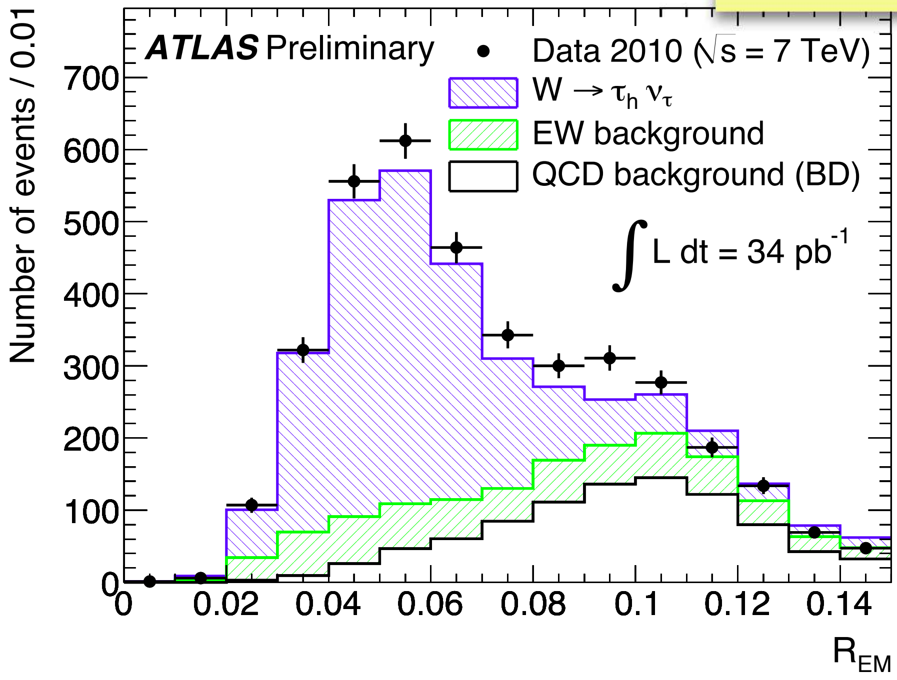
- Cuts are pt-parametrized to account for the Lorentz collimation of boosted taus.
- Experience grows with LLH and BDT discriminants, which become the preferred discriminants in 2011.

"Reconstruction, Energy Calibration, and Identification of Hadronically Decaying Tau Leptons in the ATLAS Experiment" [ATLAS-CONF-2011-077, ATL-PHYS-INT-2011-068]

Seeing first hadronic taus

2010





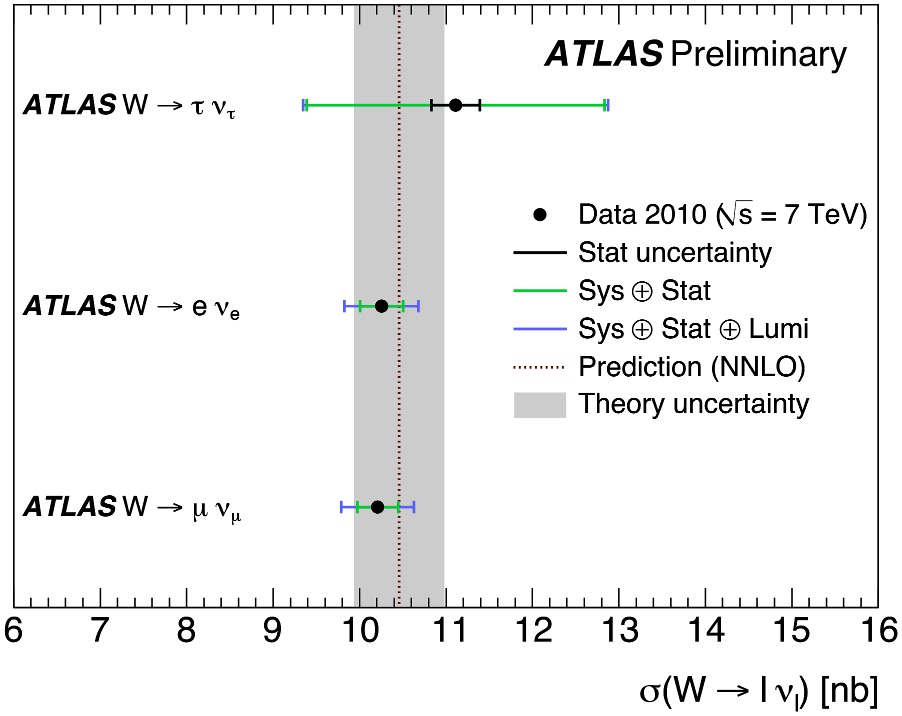
- Nov 2010: Observation of $W \to \tau_h \nu$ [ATLAS-CONF-2010-097]
- Feb 2011: Observation of $Z \to \tau_h \tau_l$ [ATLAS-CONF-2011-010]

$W \to \tau\nu$ cross section

2010

$$\sigma(W \to \tau\nu) = 11.1 \pm 0.3(\text{stat.}) \pm 1.7(\text{sys.}) \pm 0.4(\text{lumi.})\ \text{nb}$$

$$\sigma_{\text{theory}} = 10.46 \pm 0.52\ \text{nb at NNLO}$$



Dominant systematics

- τ_{h} efficiency 10.3%
- τ_{h} energy scale 8.0%
- τ_{h} + MET trigger efficiency 7.0%
- luminosity 3.4%
- acceptance 2.3%

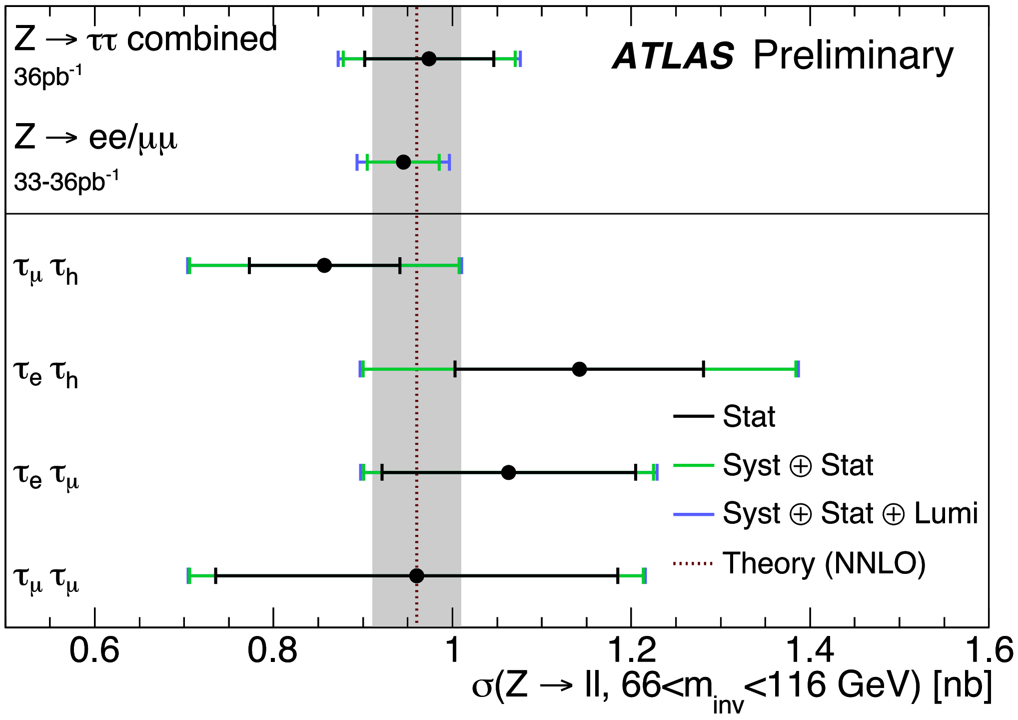
"Measurement of the $W \to \tau\nu$ cross section in pp collisions at sqrt(s)= 7 TeV with the ATLAS experiment" [arXiv:1108.4101]

$Z \to \tau\tau$ cross section

2011

$$\sigma_{\text{combined}} = 0.97 \pm 0.07(\text{stat.}) \pm 0.07(\text{sys.}) \pm 0.03(\text{lumi.})\ \text{nb}$$

$$\sigma_{\text{theory}} = 0.96 \pm 0.05\ \text{nb at NNLO}$$



Dominant systematics

- τ_h energy scale 11%
- τ_h efficiency 8.6%
- μ efficiency 8.6%
- e efficiency 3-10%
- acceptance 3%
- luminosity 3.4%

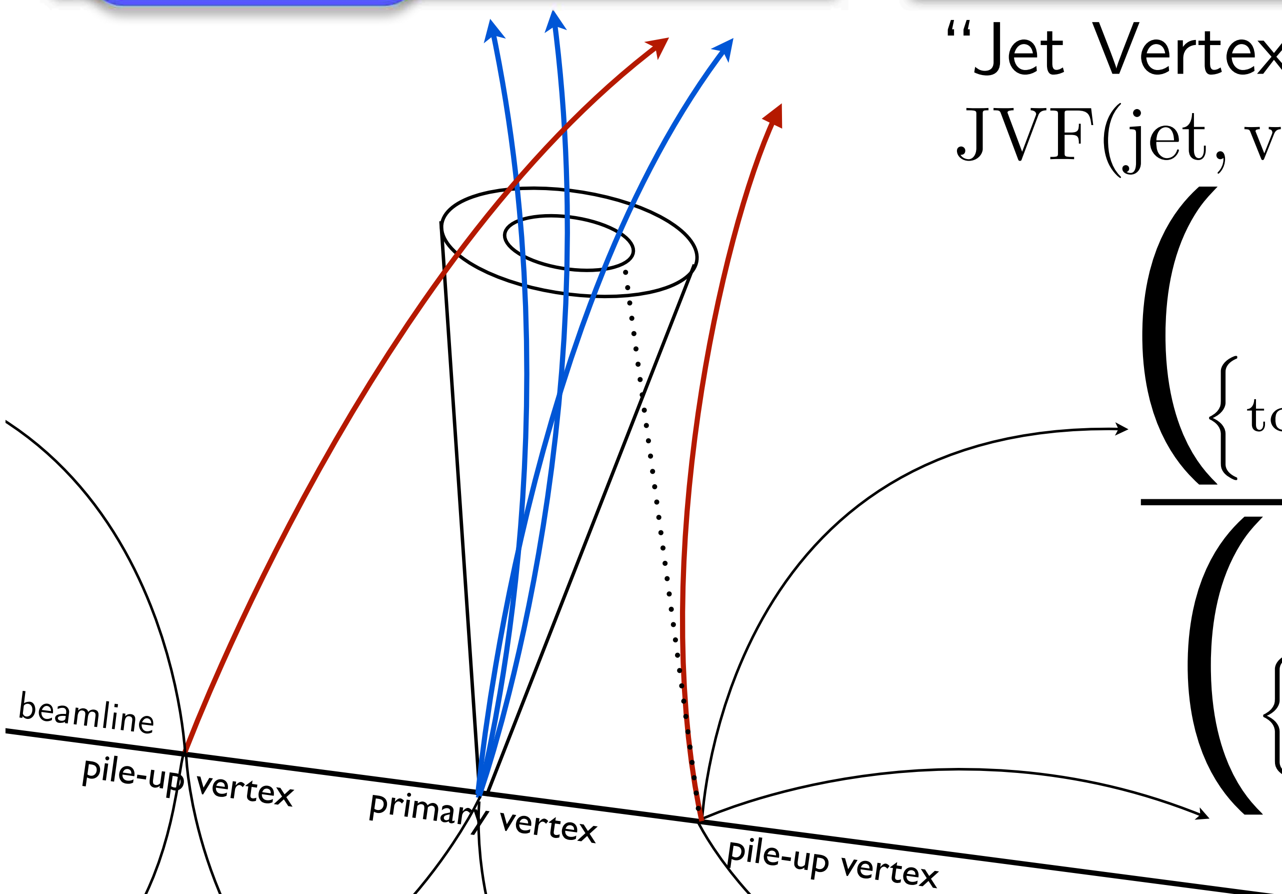
"Measurement of the $Z \to \tau\tau$ cross section in pp collisions at sqrt(s)= 7 TeV with the ATLAS detector" [arXiv:1108.2016]

Tau vertex association

Tau track selection

- $p_T > 1$ GeV,
- Number of pixel hits ≥ 2,
- Number of pixel hits + number of SCT hits ≥ 7,
- $|d_0| < 1.0$ mm,
- $|z_0 \sin\theta| < 1.5$ mm,

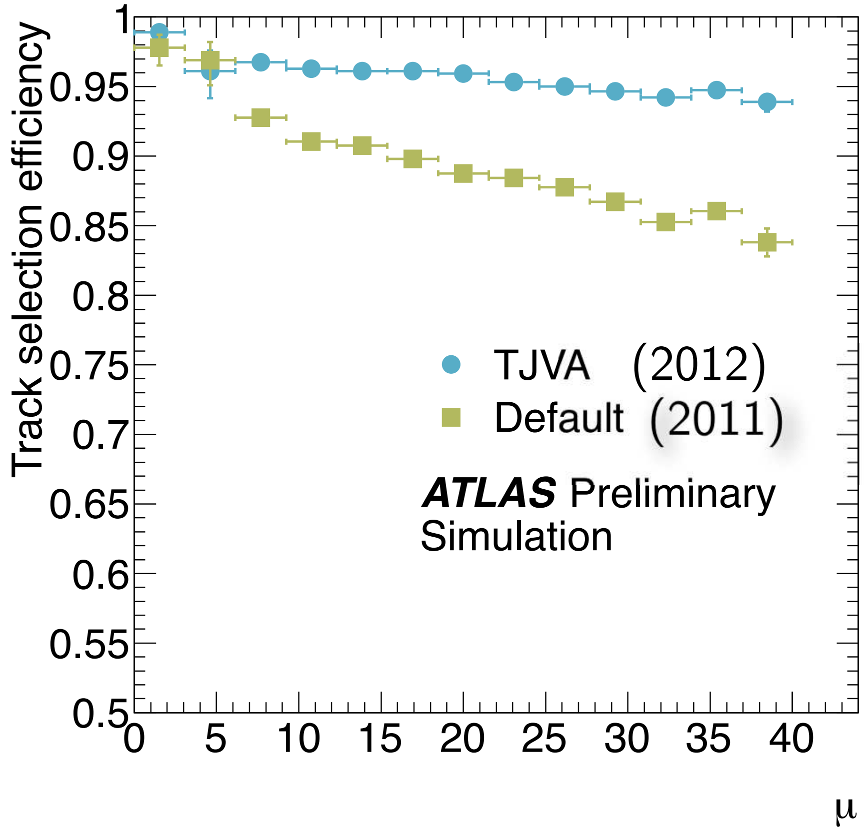
- The d_0 and z_0 requirements depend on the choice of vertex.
- Beginning in 2012, choose the vertex with the highest JVF for that tau candidate.

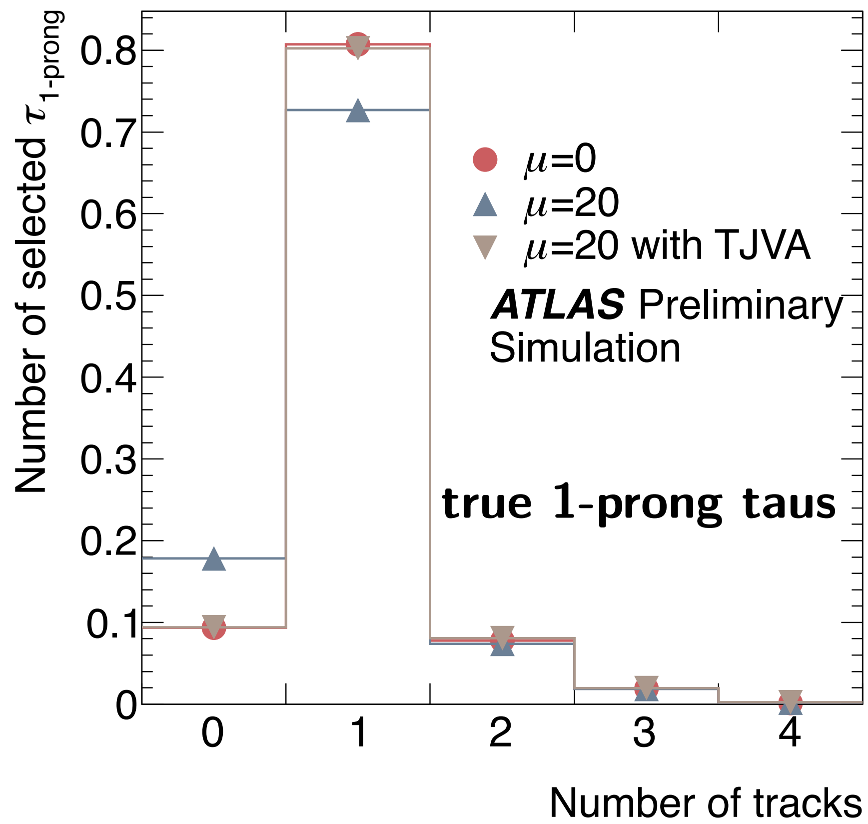


"Jet Vertex Fraction"

$$\mathrm{JVF(jet, vertex)} = \frac{\left(\sum_{\left\{\substack{\text{tracks matched}\\ \text{to jet and vertex}}\right\}} p_T(\text{track})\right)}{\left(\sum_{\left\{\substack{\text{tracks matched}\\ \text{to jet}}\right\}} p_T(\text{track})\right)}$$

Track selection efficiency

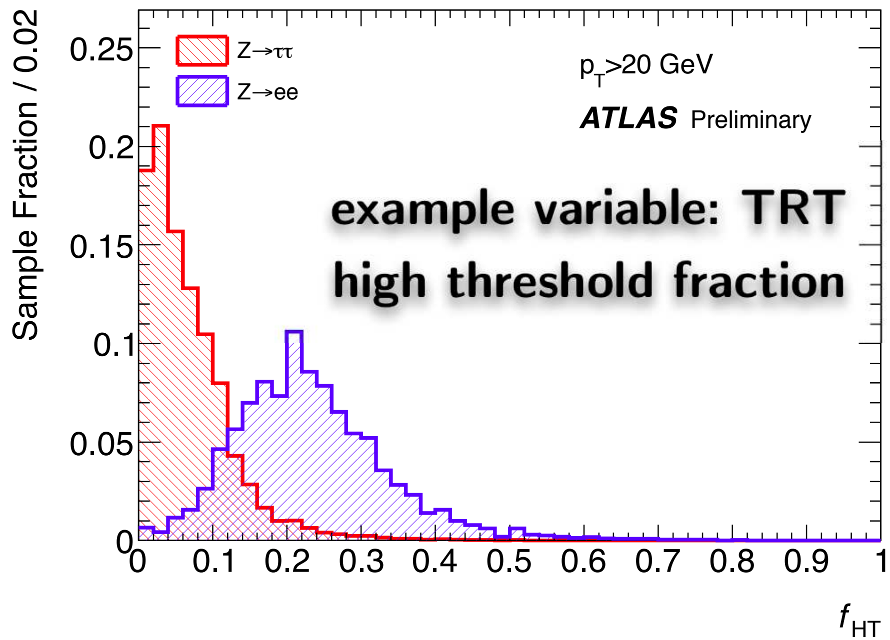


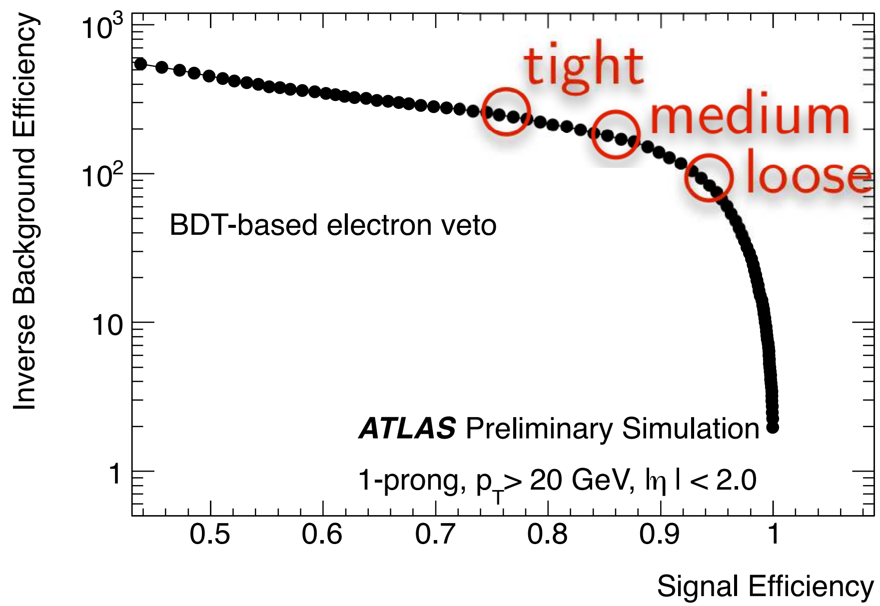


- In 2011, the track selection for tau candidates cut on the d_0 and z_0 with respect to the vertex with the highest $\sum p_T{}^2$.
- Selecting the vertex with the highest JVF recovers efficiency in high pile-up (Tau Jet Vertex Association).

Electron veto

- Electrons provide a track and calorimeter deposit that can fake hadronic tau decay identification.
- ATLAS provides a BDT to discriminate electrons from tau candidates, even after removing overlaps with selected electrons.
- Tight/Medium/Loose working points are defined (≈75%, 85%, 95% efficient).
- In 2012, the BDT is being re-optimized to have better efficiency at high-p_T.



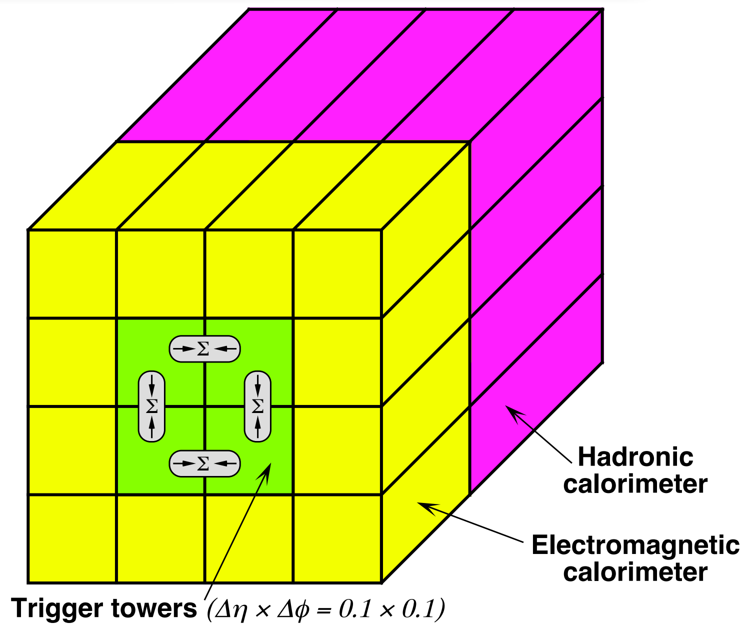


Tau triggering

1. **Level 1:** (latency 2.5 μs)

   Coarse EM+Had calorimeter trigger towers $\Delta\eta\times\Delta\phi = 0.1\times0.1$. Candidate passing thresholds on the sum of energies:

   1. highest 2×1 towers
   2. surrounding 4×4 isolation ring

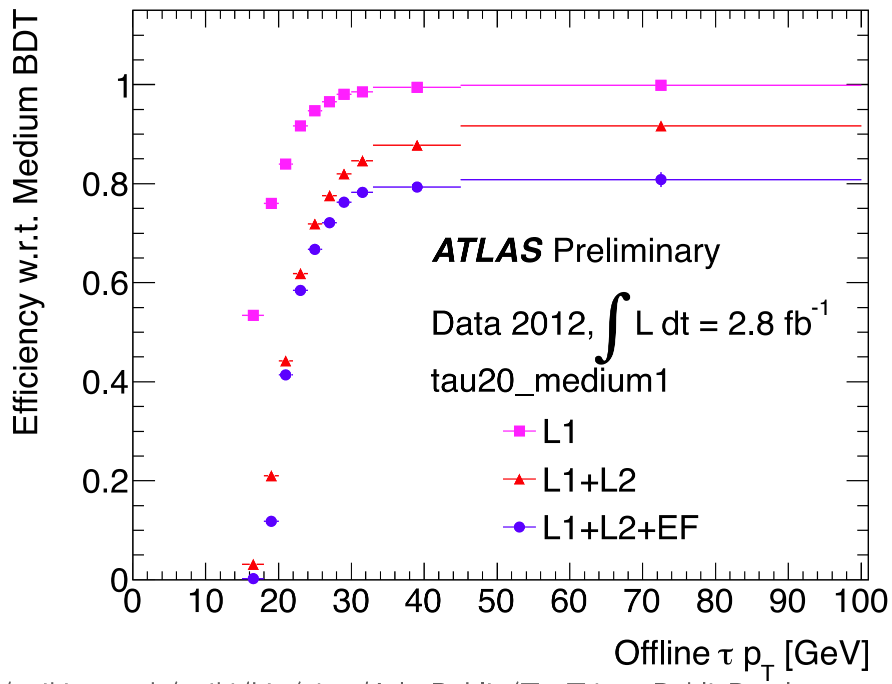


2. **Level 2:** (latency 40 ms)

   Fast tracking. Region-of-interest (RoI) calculation of track- and calorimeter-based ID variables. Similar selection to offline cut-based ID.

3. **Event Filter:** (latency 4 s)

   Beginning in 2012, started using the offline BDT algorithm at the EF trigger.



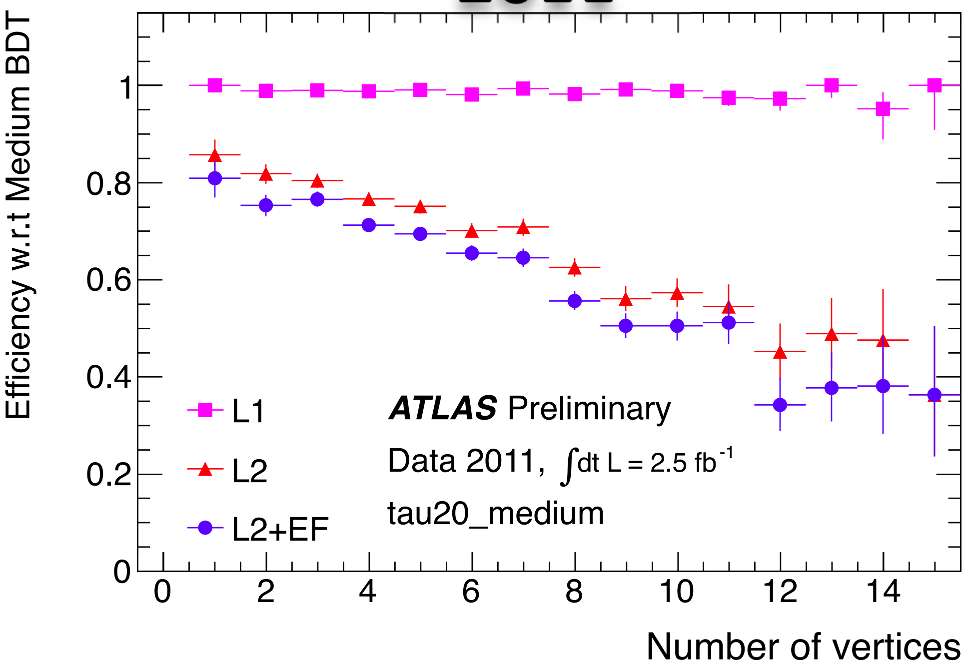
https://twiki.cern.ch/twiki/bin/view/AtlasPublic/TauTriggerPublicResults

L2 pile-up robustness

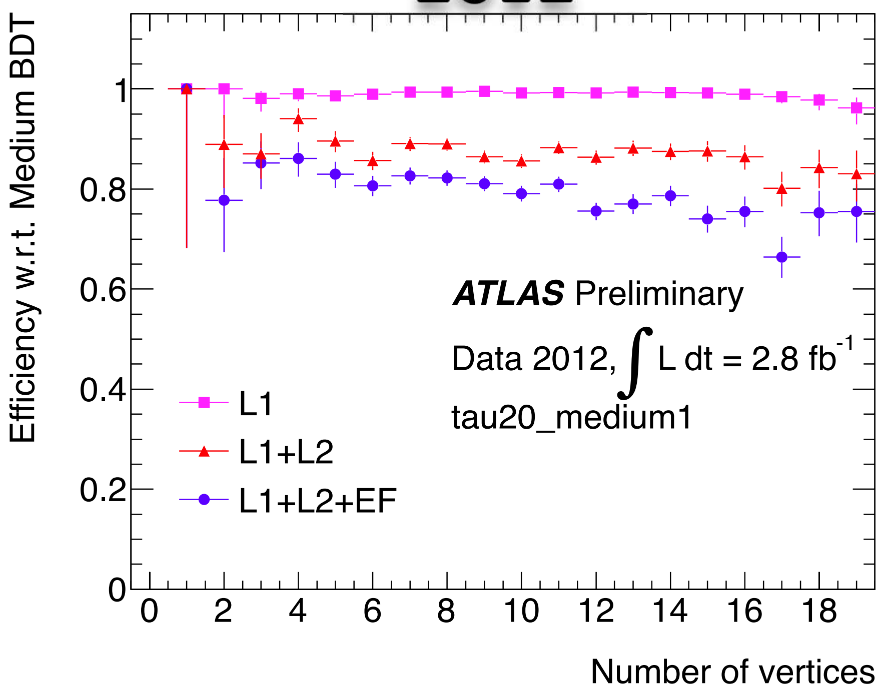
Example improvements to variable definitions to lessen sensitivity to pile-up:

- Smaller ΔR cone for calculating EM radius 0.4→0.2
- Select tracks within $\Delta z < 2$ mm of the highest-p_{T} track within the RoI (cannot vertex at L2).

2011

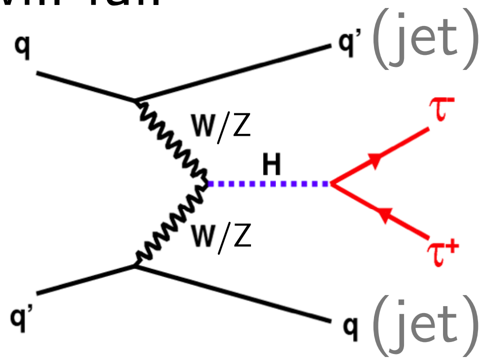


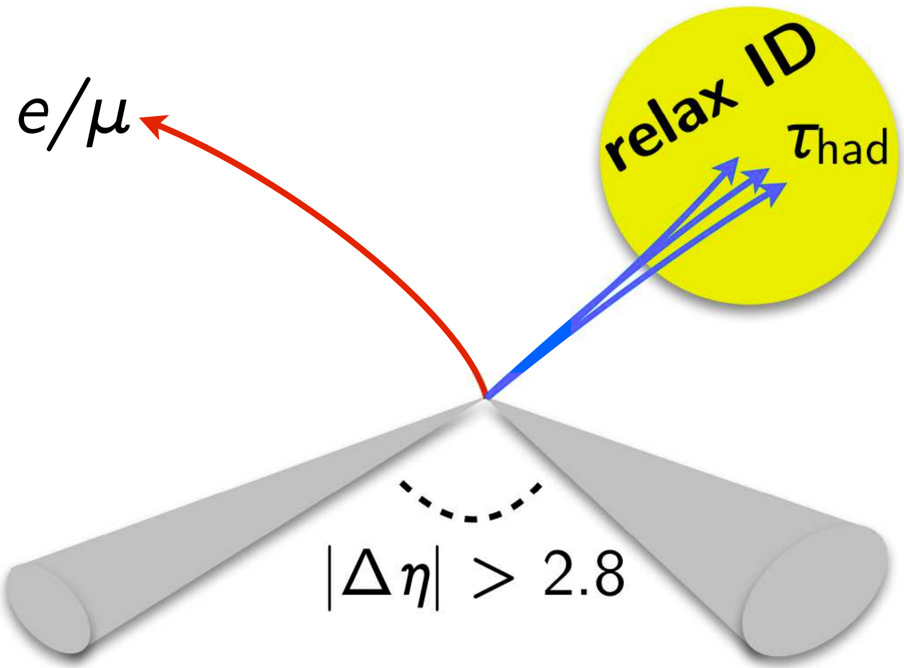
2012



VBF triggers

- New VBF triggers *relax tau identification* required at L2 and the EF by adding requirements for forward jets.
- This increases the control sample of tau candidates that will fail identification, used to ***estimate the fake contribution***.
- Being evaluated for the $H \to \tau\tau \to$ lep+τ_{had} search.





New as of periods G1, H2

VBF reqs:

- 2 L2 jets $p_{\mathrm{T}} > 15$ GeV, $|\Delta\eta| > 2.5$
- 2 EF jets $p_{\mathrm{T}} > 25$ GeV, $|\Delta\eta| > 2.8$, $M_{jj} > 400$ GeV

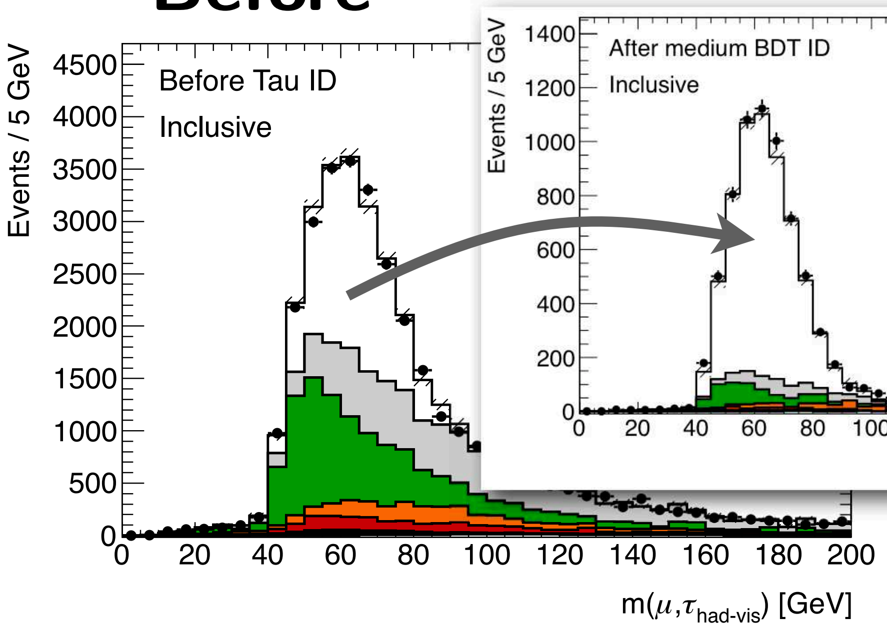
Trigger menu

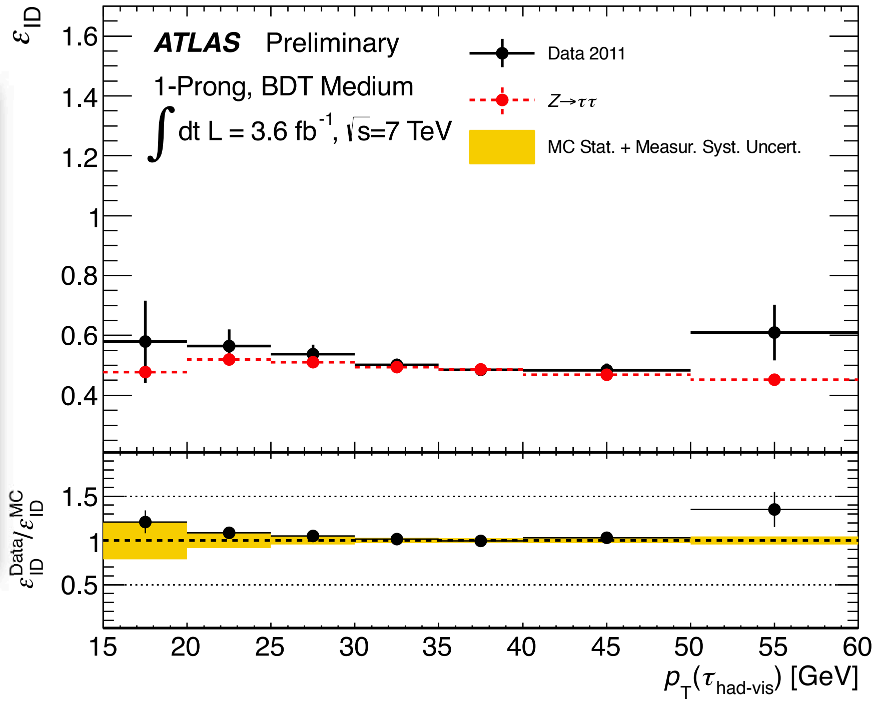
$\tau_{\mathrm{had}}+\mu$	$\tau_{\mathrm{had}}+e$
tau20_medium1_mu15	tau20Ti_medium1_e18vh_medium1
mu15_vbf_L1TAU8_MU10	e18vh_medium1_vbf_2L1TAU11I_EM14VH

Identification efficiency

- ***Tag-and-probe*:** selecting a sample of a known composition without some ID, so one can probe its efficiency.
- For the case of tau ID, select $Z \to \tau\tau \to \mu\tau_h 3\nu$ by triggering on the muon and selecting events with muon + tau candidate.

Before **After Tau ID**

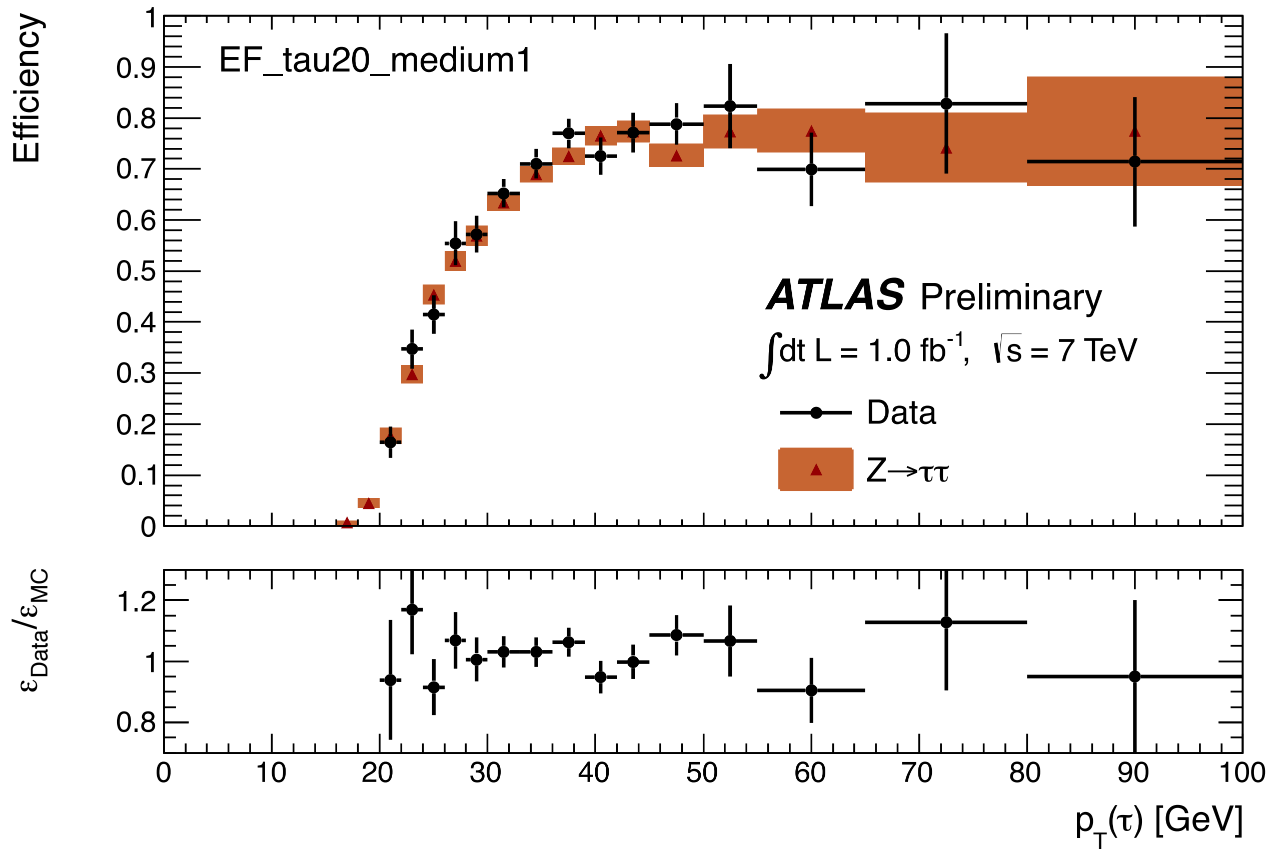




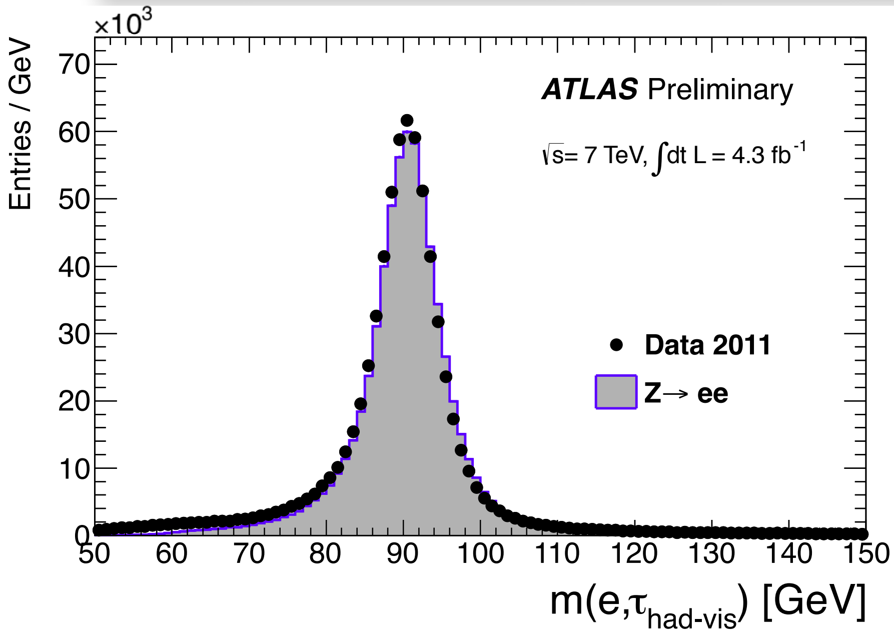
- Scale factor ≈ 1, known to a few percent, 2-3% (1-prong), 5-6% multi-prong.

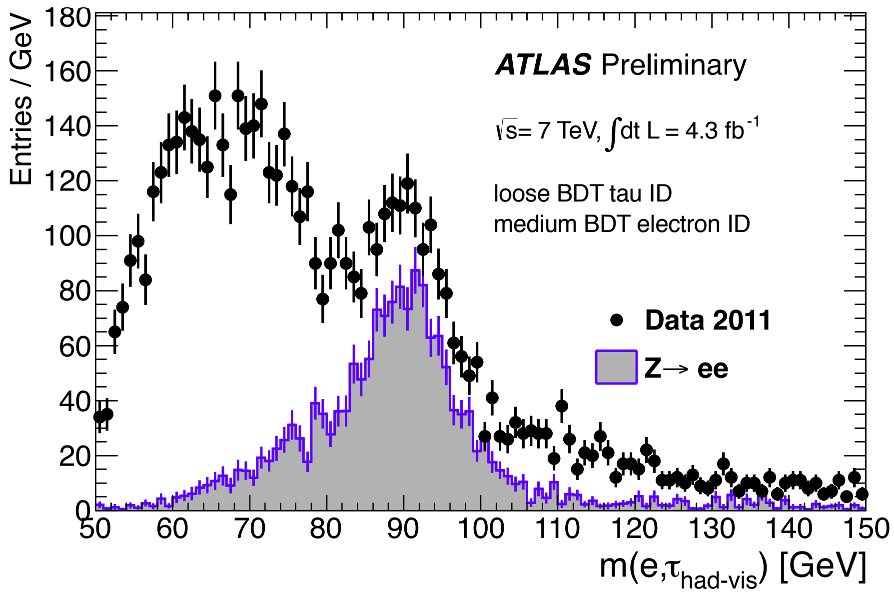
Trigger efficiency

- The same $Z \to \tau\tau \to \mu\tau_h 3\nu$ tag-and-probe sample is used to measure the efficiency of the tau triggers.
- Known to $O(5\%)$ in the turn-on.
- Improving with statistics in 2012.



Electron veto fake rate





- Tag e + tau candidates
- Probe the e-veto efficiency after removing overlap with selected electrons.

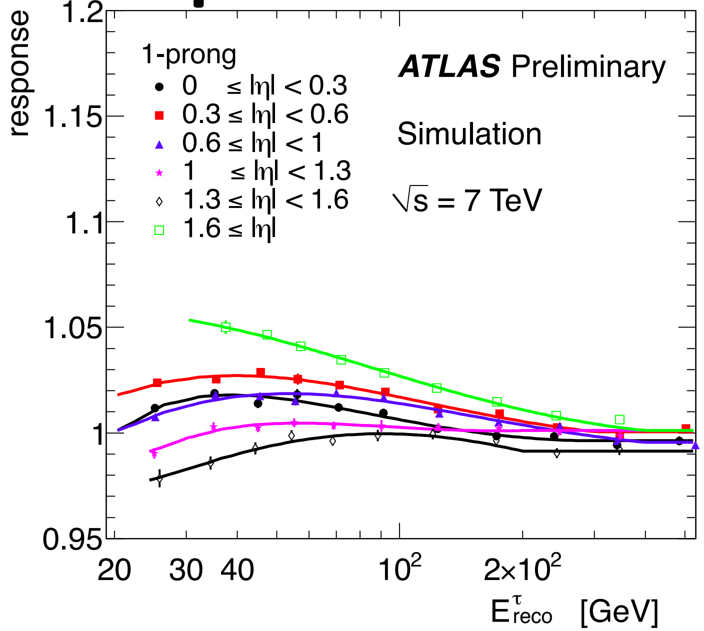
data/MC scale factor and uncertainty
from Z→ee tag-and-probe with 2.6/fb from 2011

electron BDT veto	$\lvert\eta_{trk}\rvert < 1.37$	$1.37 < \lvert\eta_{trk}\rvert < 1.52$	$1.52 < \lvert\eta_{trk}\rvert < 2.00$	$\lvert\eta_{trk}\rvert > 2.00$
loose	0.96±0.22	0.8±0.3	0.47±0.14	1.7 ±0.4
medium	1.3 ±0.5	-	0.5 ±0.4	2.8 ±1.3

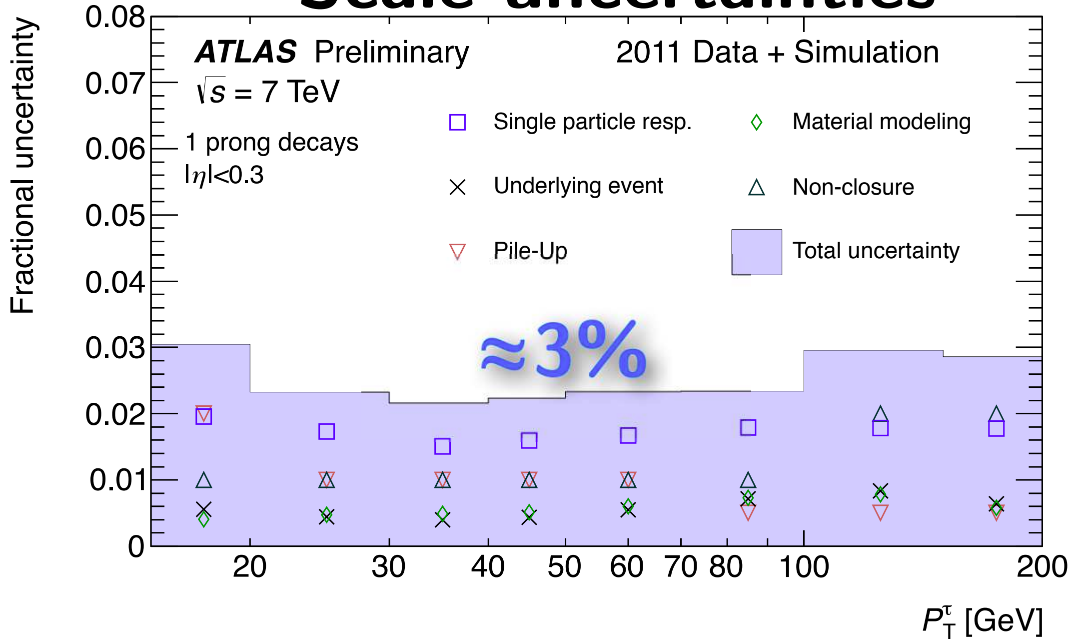
- Statistically limited by the sample that pass the veto, giving uncertainties ≈ 50-100%.
- Improving with the data added in 2012.

Energy scale

Response functions



Scale uncertainties



- Tau candidates are first brought from the EM to the Jet Energy Scale with LC calibration of the clusters within $\Delta R < 0.2$ (from 0.4 to be pile-up robust).
- Then response functions are calibrated with tau Monte Carlo to make final corrections of a few percent.
- Uncertainties are determined by smearing the Monte Carlo truth according the tau decays true composition, using uncertainties constrained by single particle response measurements (CTB, E/p, $Z \rightarrow ee/\pi^0$-resp.)

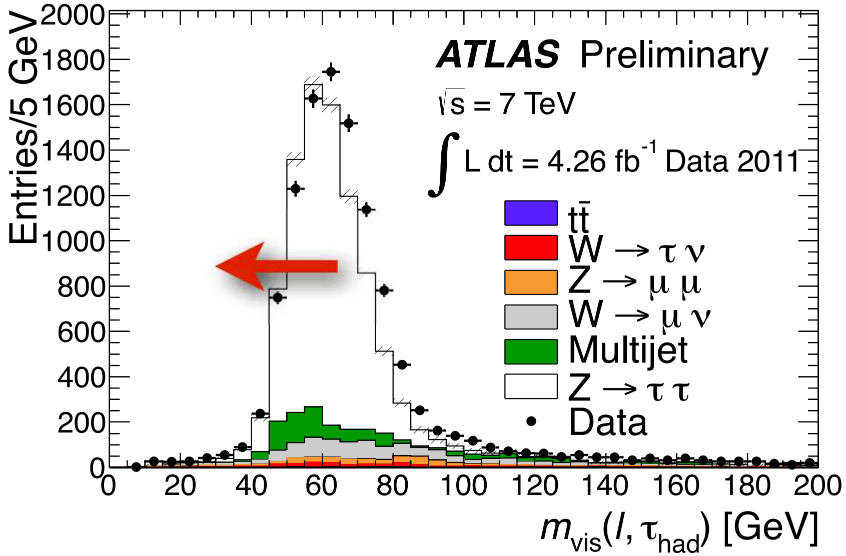
Energy scale cross check

- Tau energy scale is manually shifted in the modeling.
- Median of the visible mass peak is used to decide which scale matches the data.
- Toy experiments are used to estimate the uncertainty.

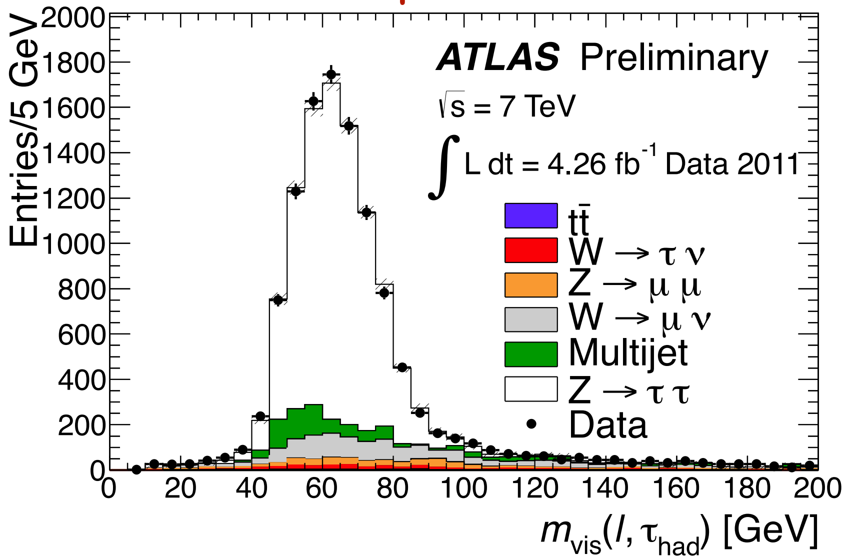
$\|\eta\|$	best scale	uncert.
0.0-0.8	-1.5%	3.3%
0.8-2.5	+1.5%	2.8%

- Scale consistent with 1 within single-particle-response uncertainties ≈ 3%.
- May become primary method with more data.

scale shifted -10%

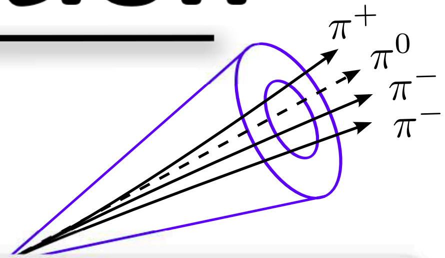


best fit +1.5%



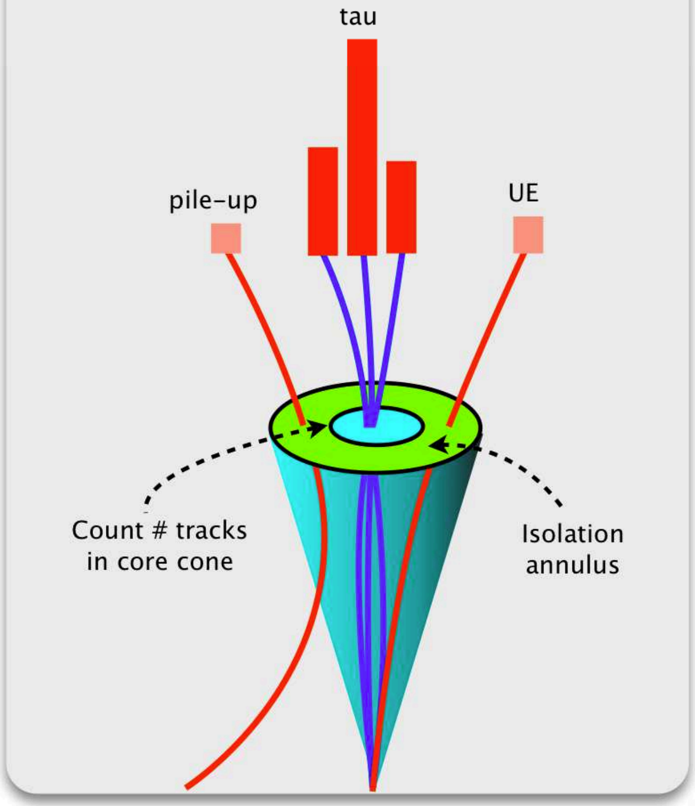
CMS Tau ID

High-p_T τ_h reconstruction



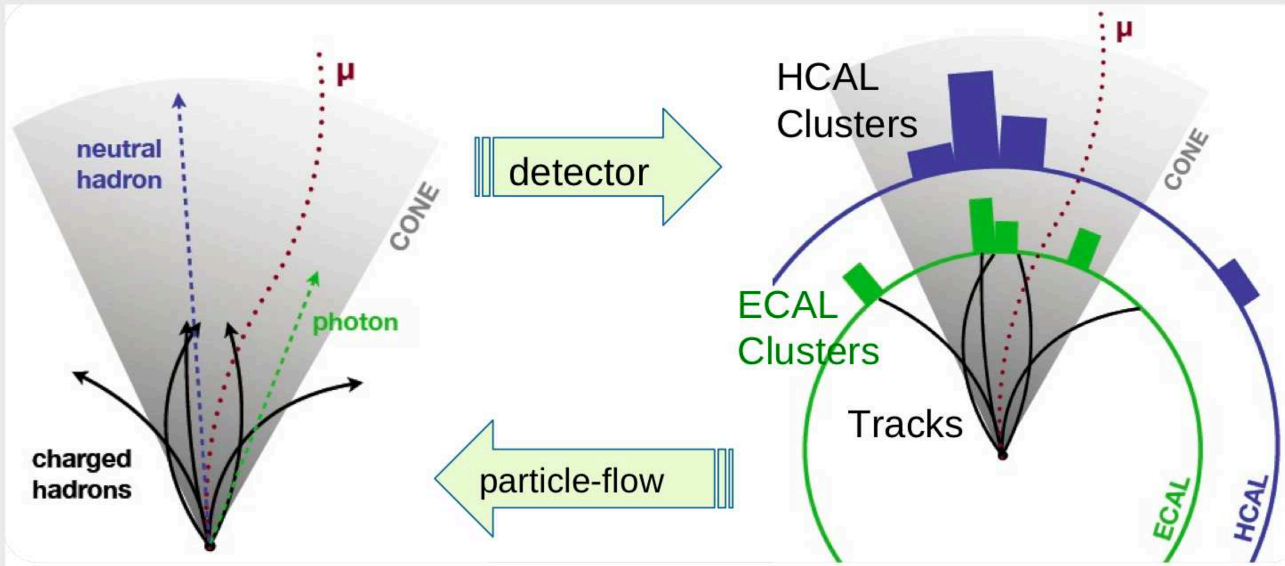
- Hadronic decays dominantly to 1 or 3 $\pi^\pm$ and possibly a few additional π^0s
- Decay in beam-pipe: $c\tau \approx 87$ μm

ATLAS



- τ_h reco seeded by calorimeter jets
- associate tracks in $\Delta R < 0.2$, select 1 or 3
- combine calorimeter and tracking information in a BDT or likelihood discriminant, preferring narrow clustering, hadronic activity

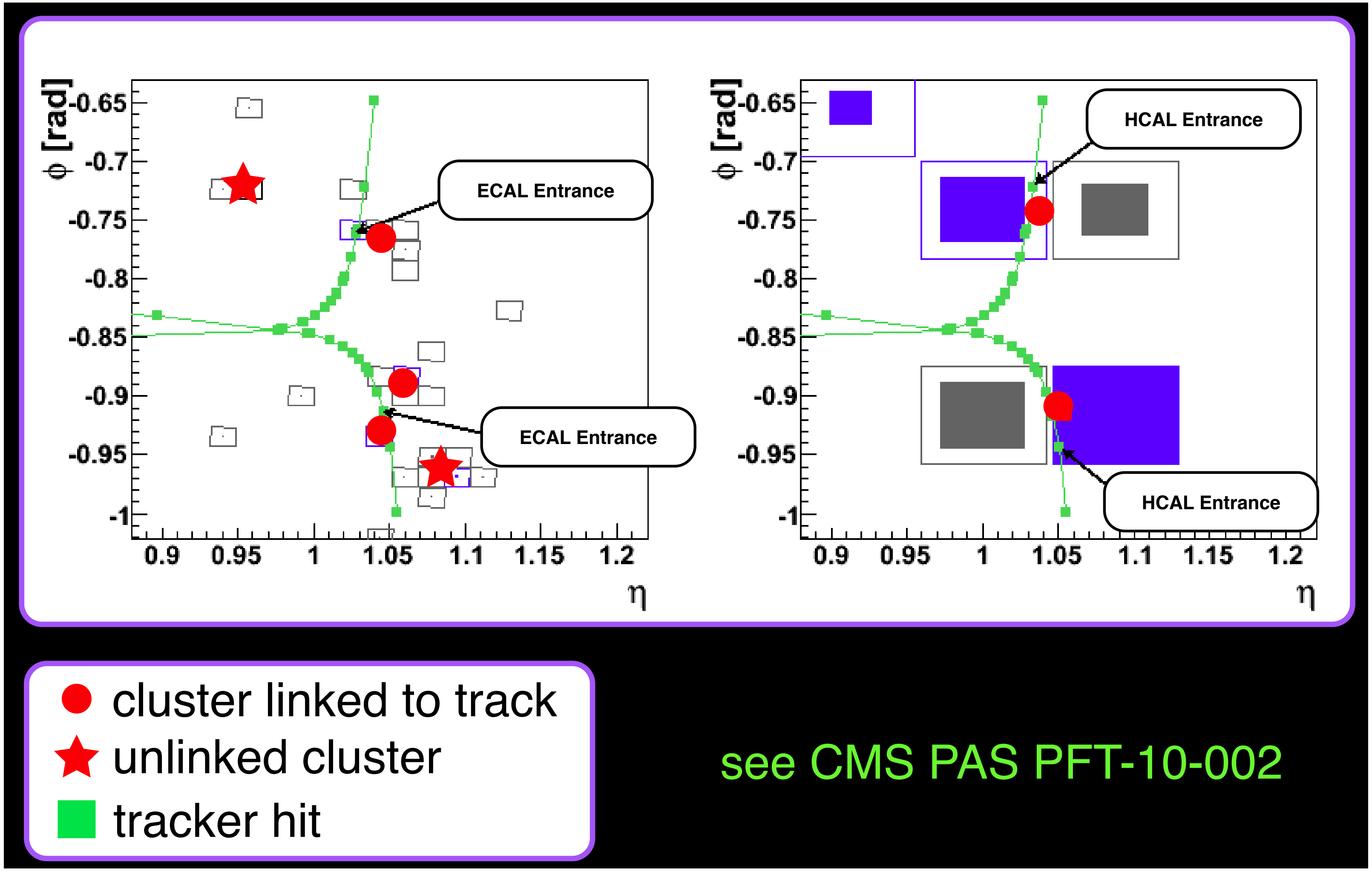
CMS



- particle-flow reconstructs constituent 4-vectors
- τ_h reco seeded by particle-flow hadrons
- Hadron Plus Strip (HPS) algorithm for counting π^0s
- isolation cone for rejecting QCD jets

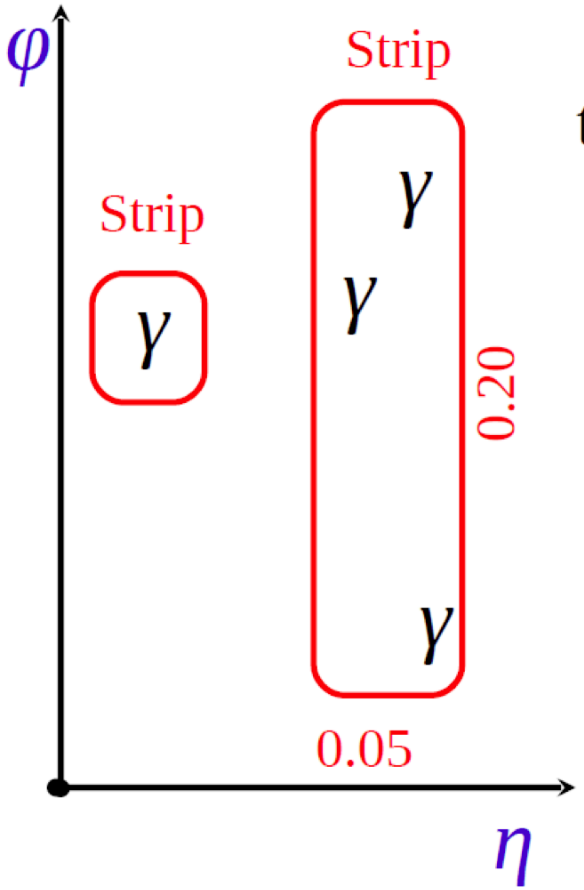
[ATLAS-CONF-2011-152, CMS PAS TAU-11-001]

CMS Particle Flow



- Matches track to clusters to form charged and neutral PF objects.
- PF objects are used as input for all CMS tau reconstruction.

CMS: Hadron Plus Strip (HPS)



Build all possible taus
that have a 'tau-like'multiplicity
from the seed jet

π^+

$\pi^+ \pi^0$

$\pi^+ \pi^+ \pi^-$

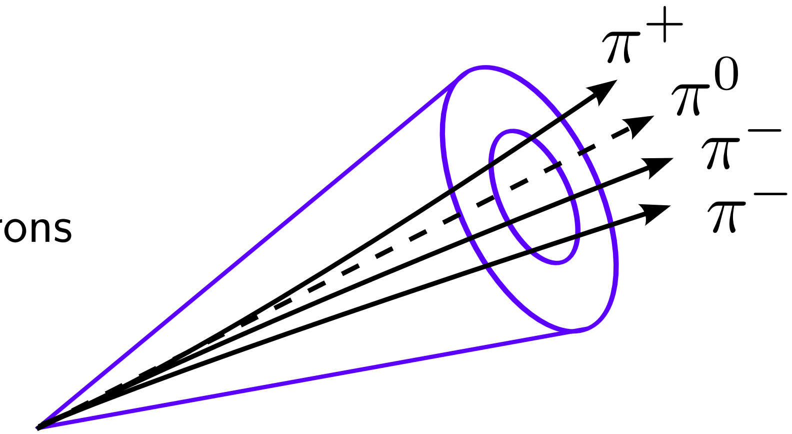
tau that is 'most isolated'
with compatible m_{vis}
is the final tau candidate
associated to the seed jet

Discrimination with calorimeter based isolation $\Delta R < 0.5$.

[CMS PAS TAU-11-001]

CMS: Tau Neural Classifier (TaNC)

- Uses a ***shrinking core-cone***:
  - $\Delta R(\text{photons}) < 0.15$ for photons
  - $\Delta R(\text{charged}) < (5\ \text{GeV})/E_T$ for charged hadrons
  - $\Delta R(\text{charged}) < \Delta R(\text{isolation}) < 0.5$



- Immediately discarded if the candidate doesn't match an expected tau decay mode.

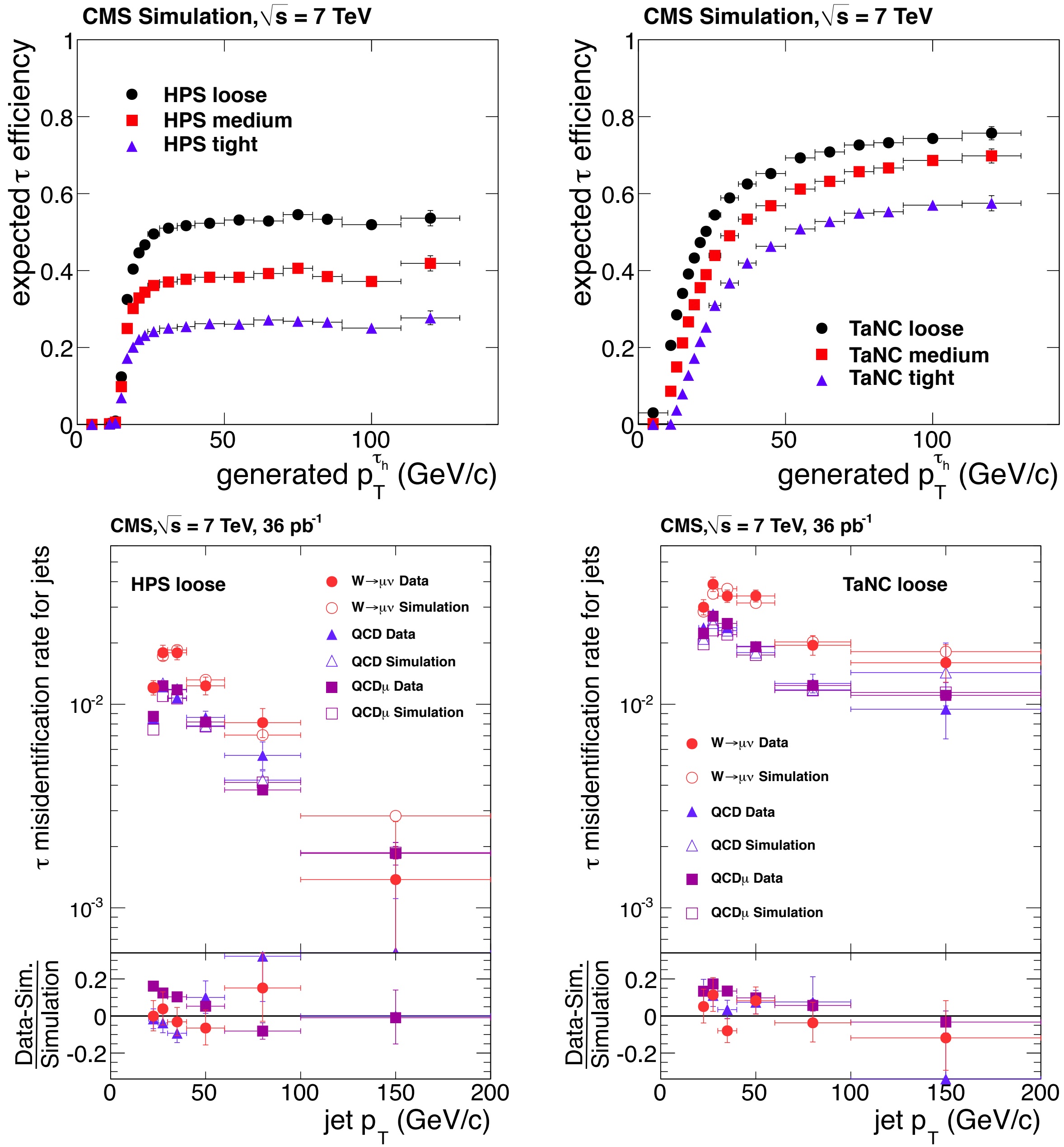
Decay mode	Resonance	Mass (MeV/c^2)	Branching fraction (%)
$\tau^- \to h^- \nu_\tau$			11.6%
$\tau^- \to h^- \pi^0 \nu_\tau$	ρ^-	770	26.0%
$\tau^- \to h^- \pi^0 \pi^0 \nu_\tau$	a_1^-	1200	9.5%
$\tau^- \to h^- h^+ h^- \nu_\tau$	a_1^-	1200	9.8%
$\tau^- \to h^- h^+ h^- \pi^0 \nu_\tau$			4.8%

- Dedicated Neural-net classifier for each decay mode

[CMS PAS TAU-11-001]

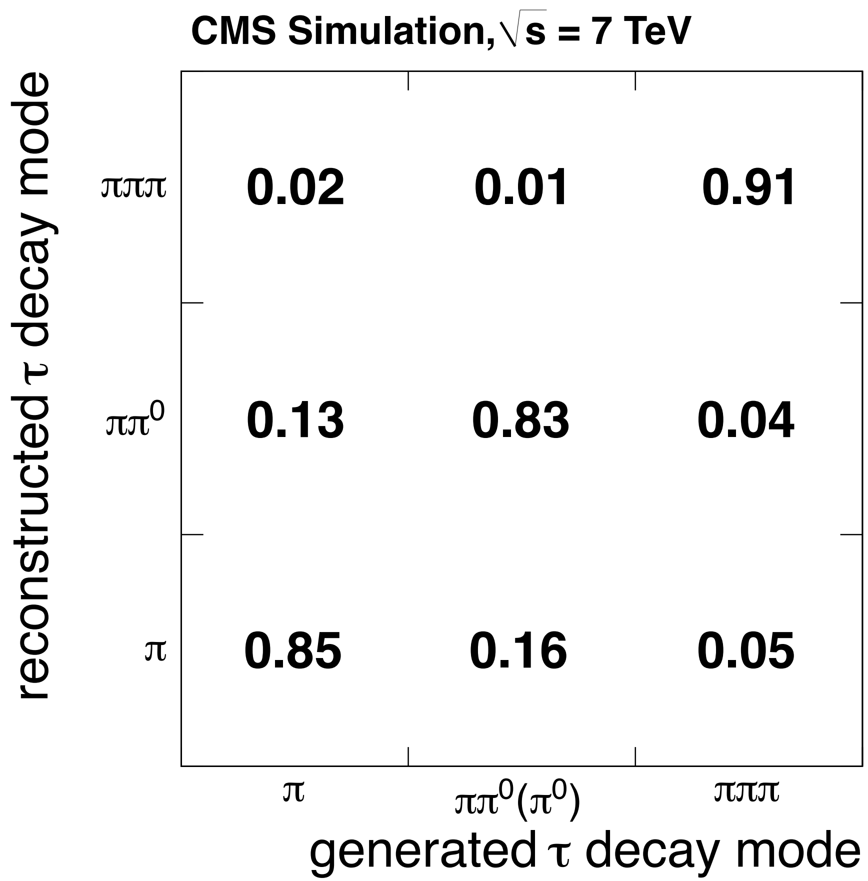
CMS Performance

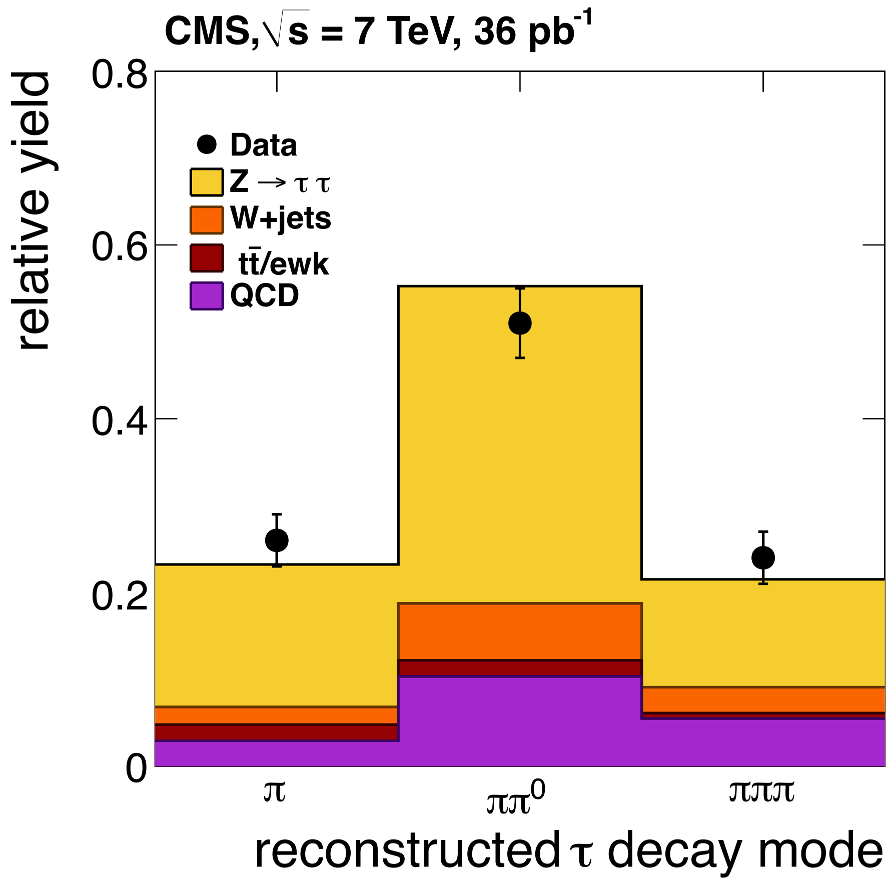
- Not trivial to compare ATLAS and CMS tau performance because we bin fake-rates in N(track) instead of categorizing the decay mode.



[CMS PAS TAU-11-001]

CMS decay mode ID





[CMS PAS TAU-11-001]

Calorimeter granularity

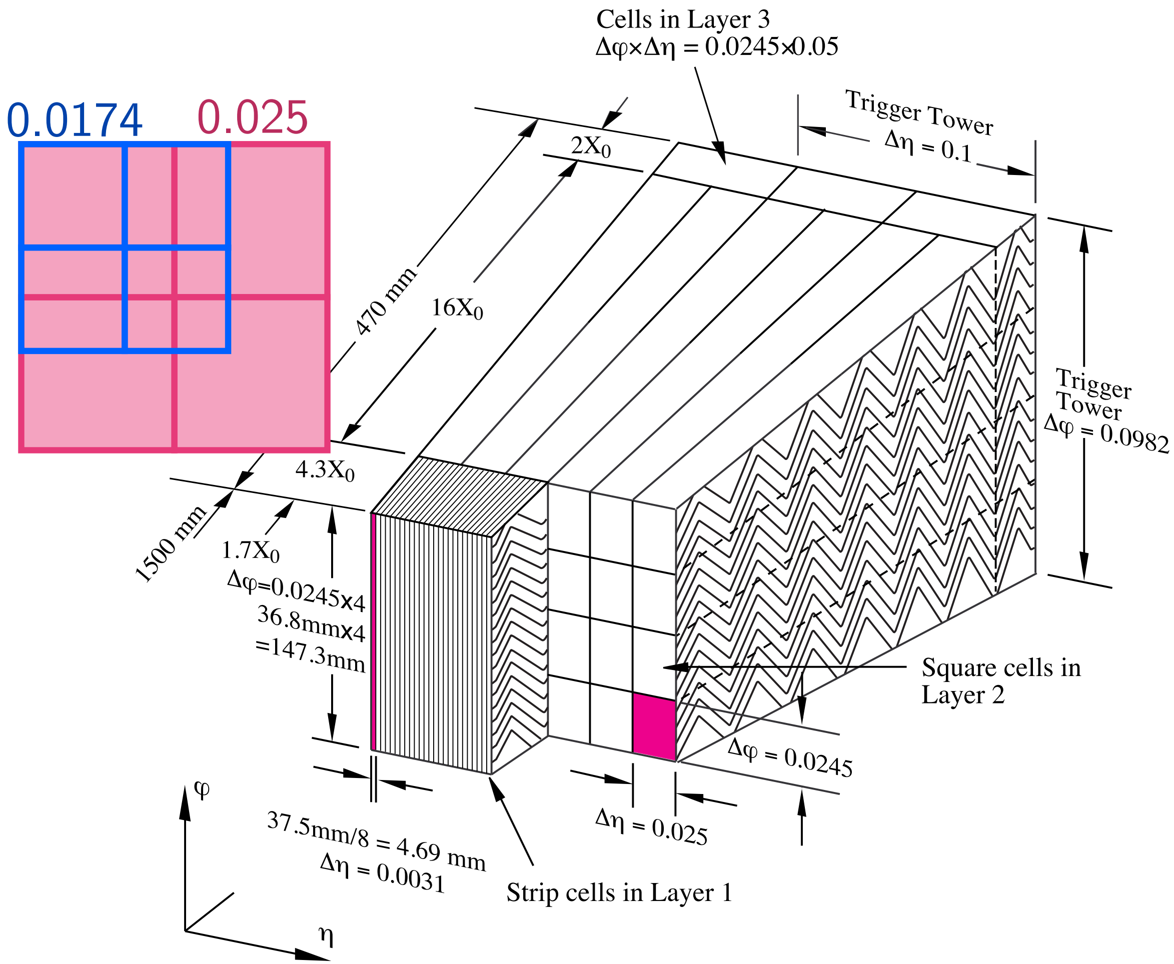
ATLAS

- $B = 2.0$ T
- $\Delta\eta \times \Delta\phi =$ 0.025×0.0245
- R = 0.4 anti-k_T topo-jets

CMS

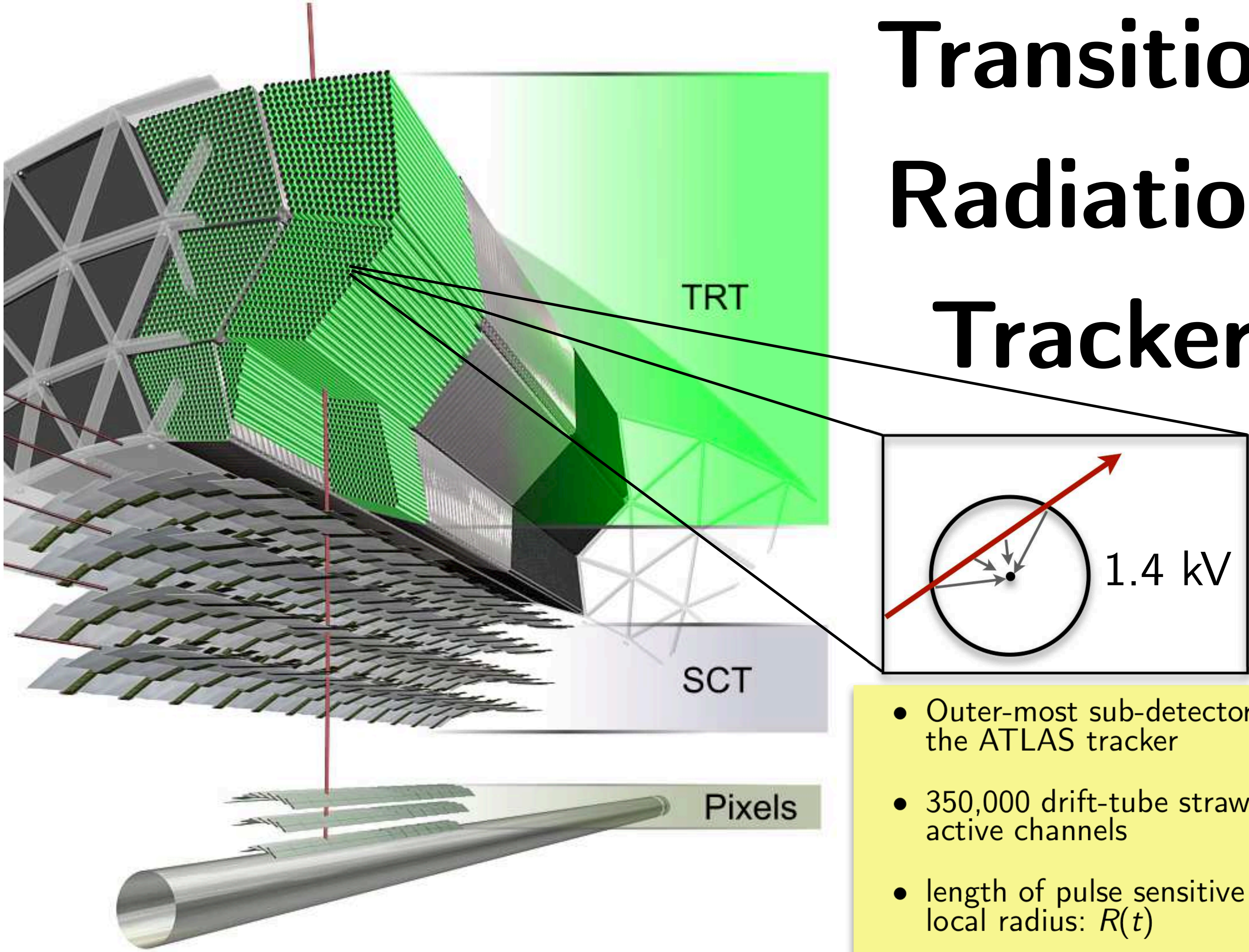
- $B = 3.8$ T
- $\Delta\eta \times \Delta\phi =$ 0.0174×0.0174
- R = 0.5 anti-k_T PF-jets

ATLAS Barrel EM Calorimeter

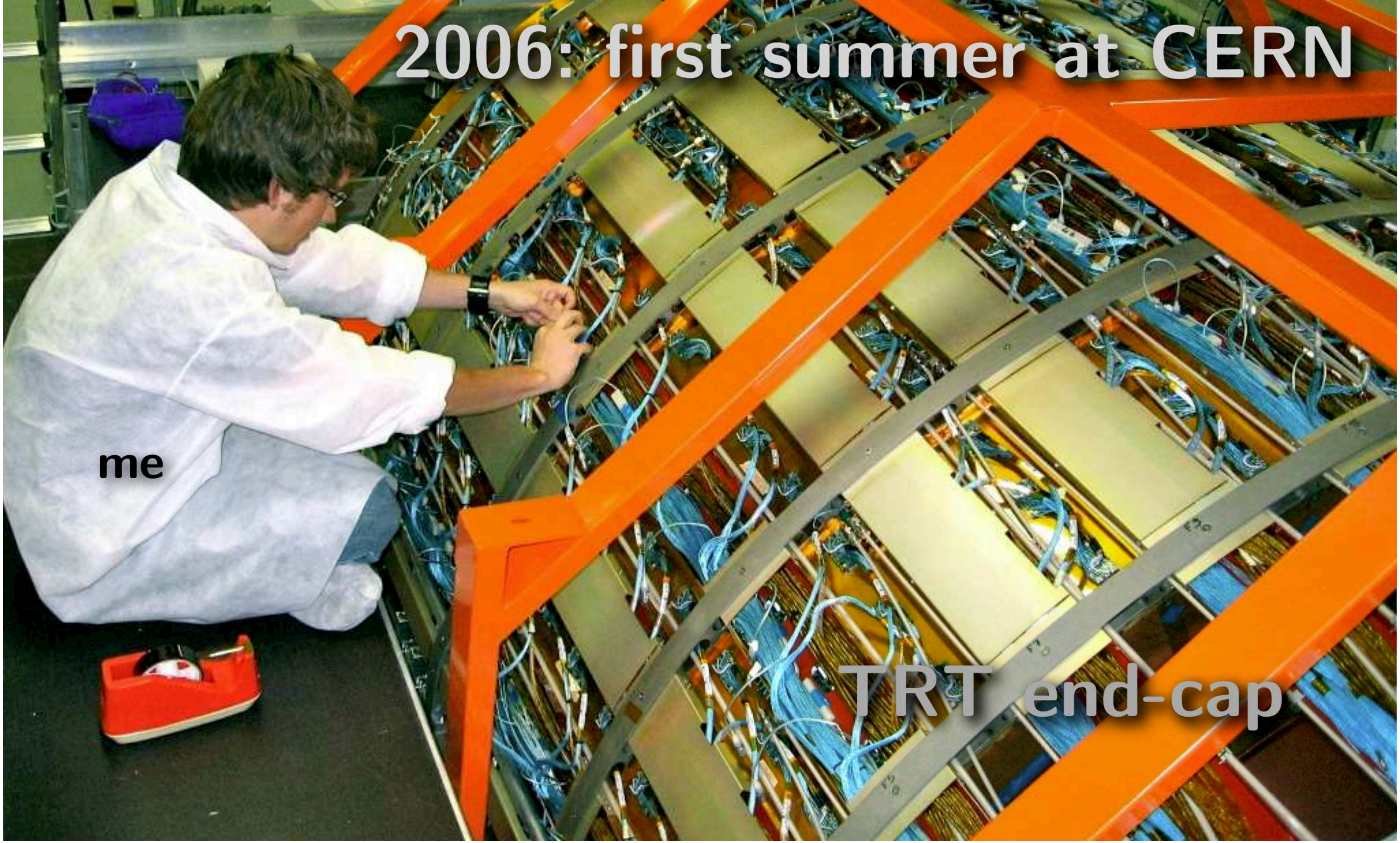


Granularity could fundamentally limit our capacity to reconstruct sub-structure / π^0s.

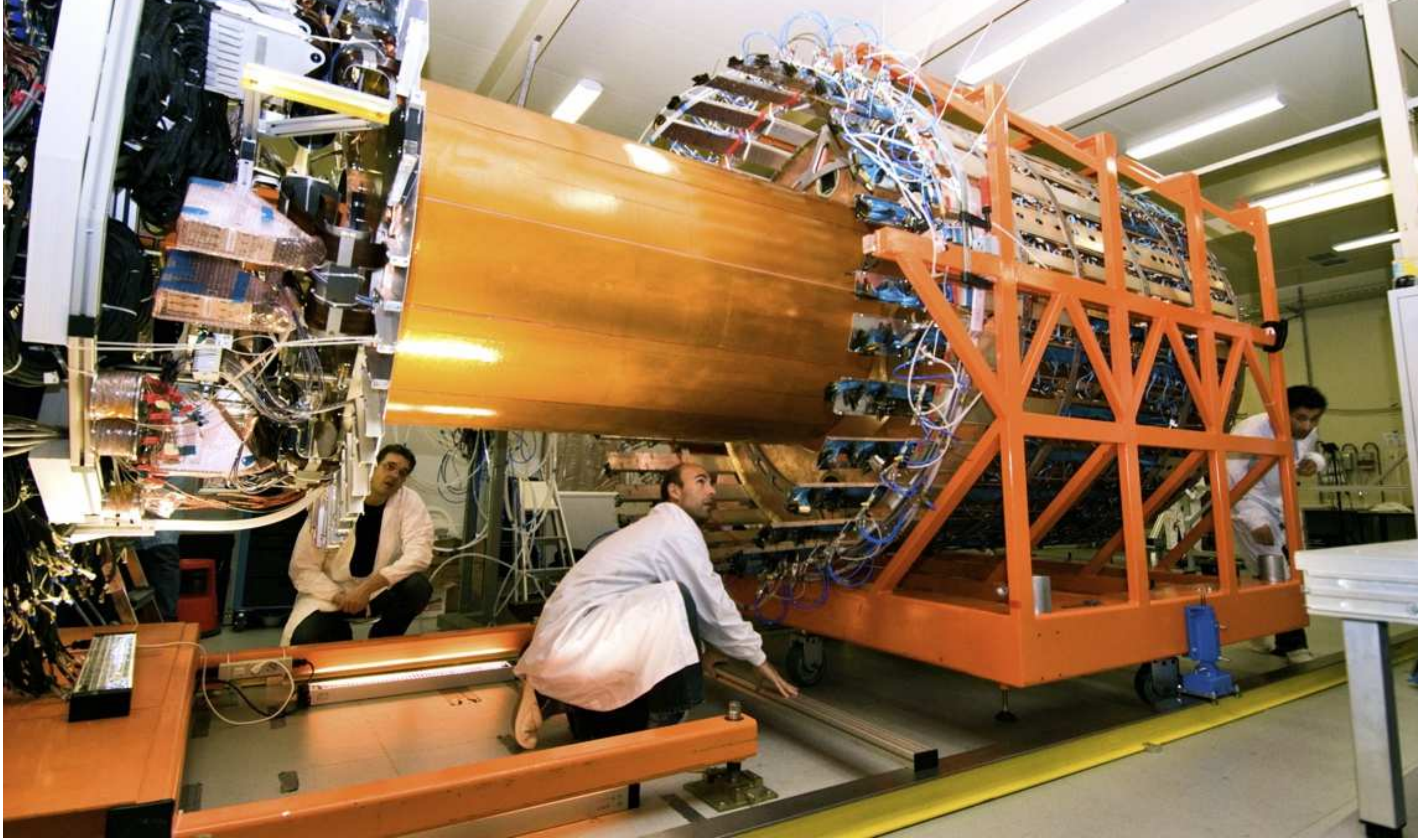
Transition Radiation Tracker



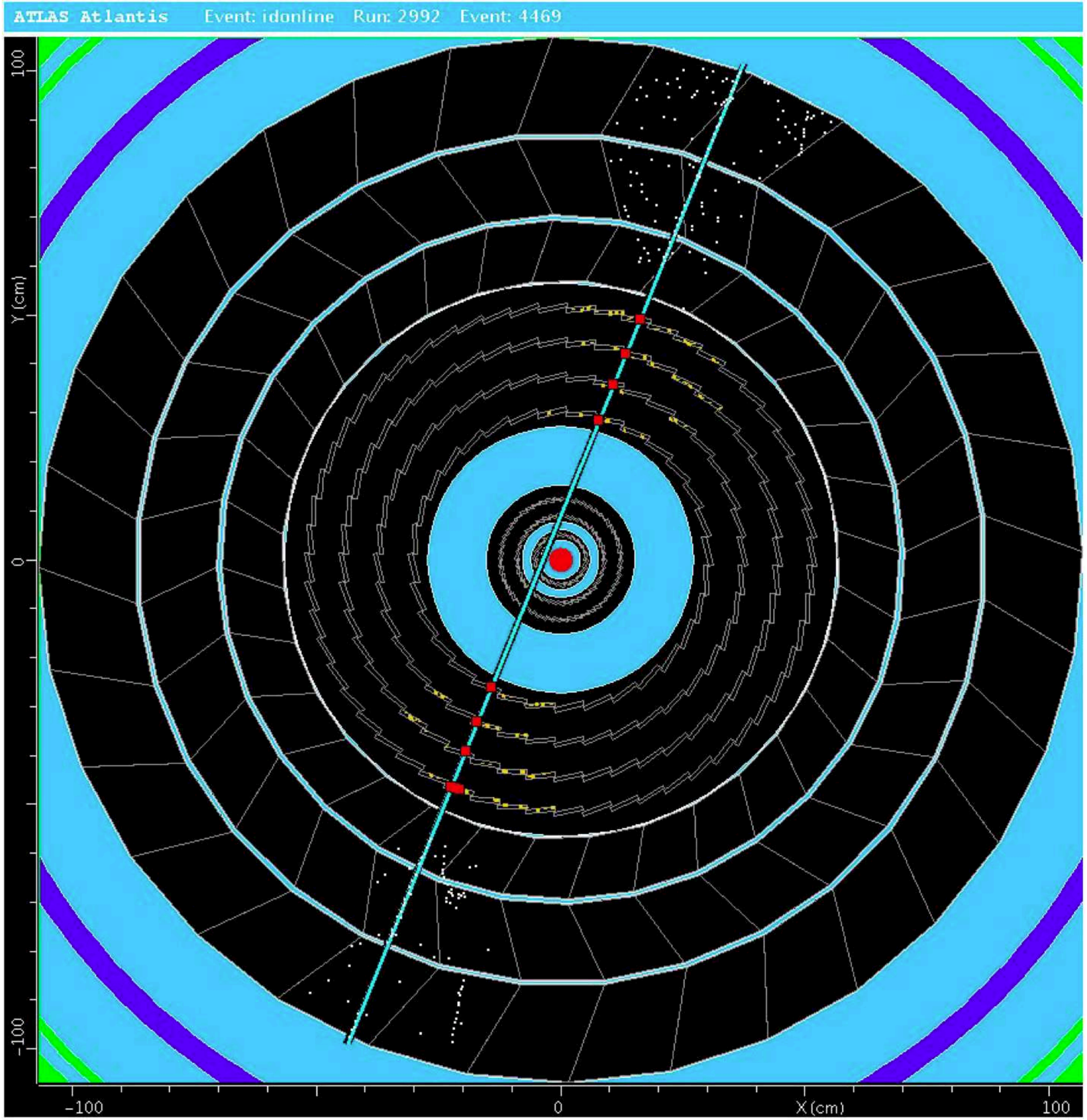
- Outer-most sub-detector of the ATLAS tracker
- 350,000 drift-tube straws = active channels
- length of pulse sensitive to local radius: $R(t)$
- 120 μm resolution



- Participated TRT integration in the above-ground cleanroom (SR-1).
- Helped with cabling and connectivity tests.



- Witnessed the marrying of the SCT and TRT barrel and end-caps.
- Amazed by all the efforts I found at CERN.

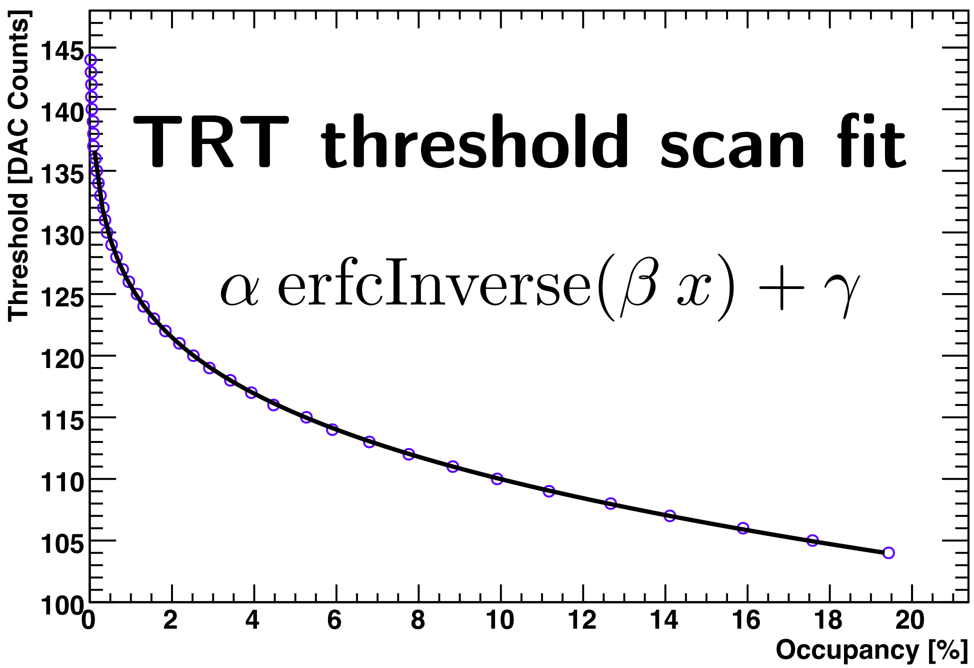


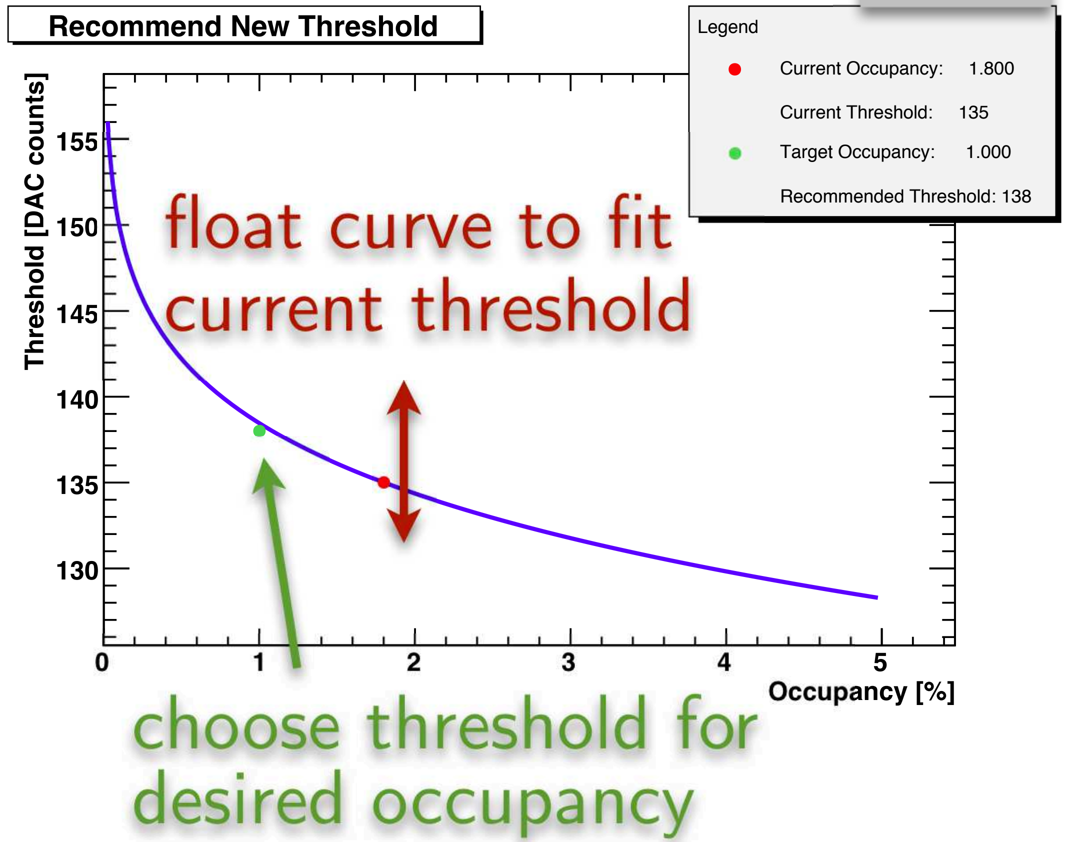
2006

- Did my first “shift” as we saw the first cosmics in the few ϕ-slices that were fully instrumented.

TRT threshold calibration

2008



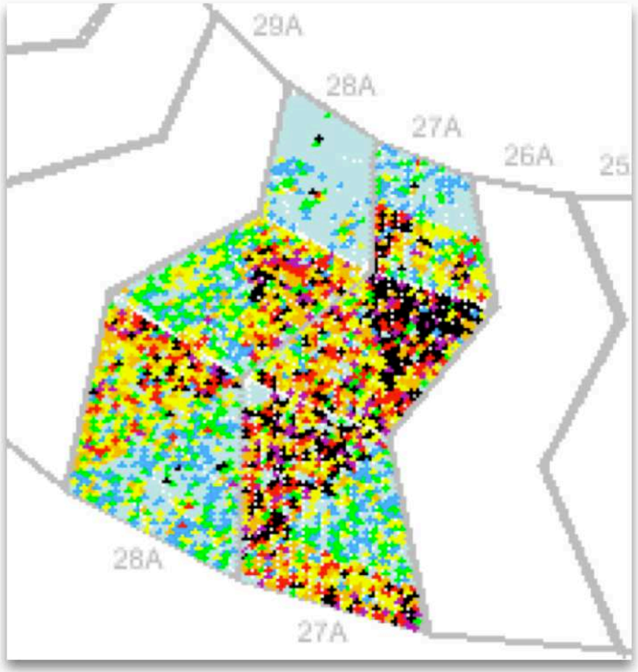


- Fit a parametrization to the data from a scan varying the analog-to-digital thresholds in the discriminator chips on the TRT front-end.
- Wrote an algorithm for calibrating the thresholds *channel-by-channel* to reach a uniform noise occupancy by floating the functional form.

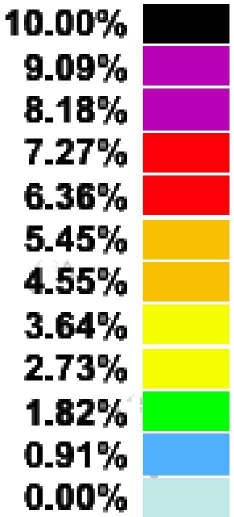
TRT threshold calibration

2008

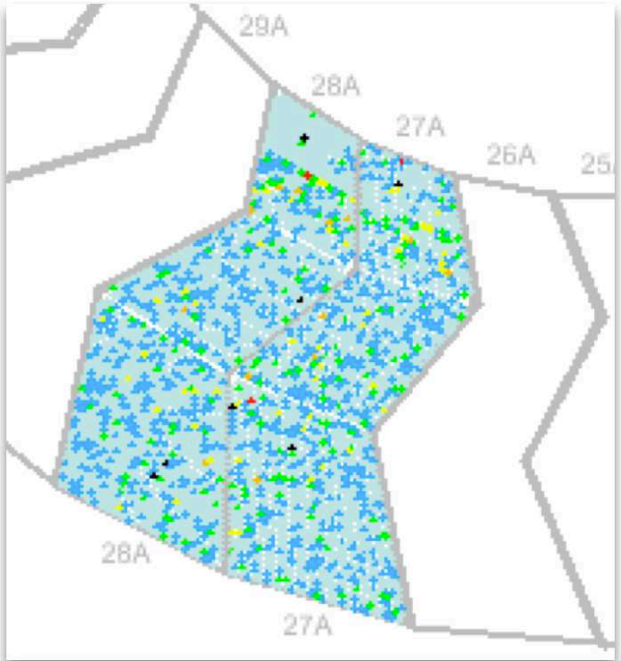
Before



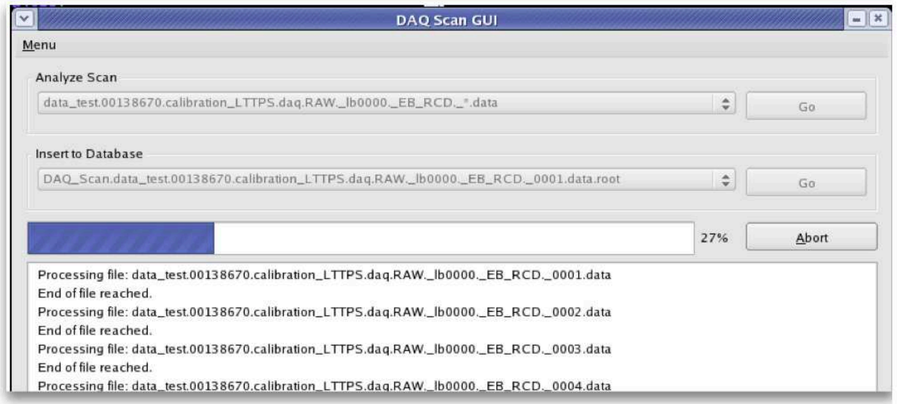
noise occupancy



After

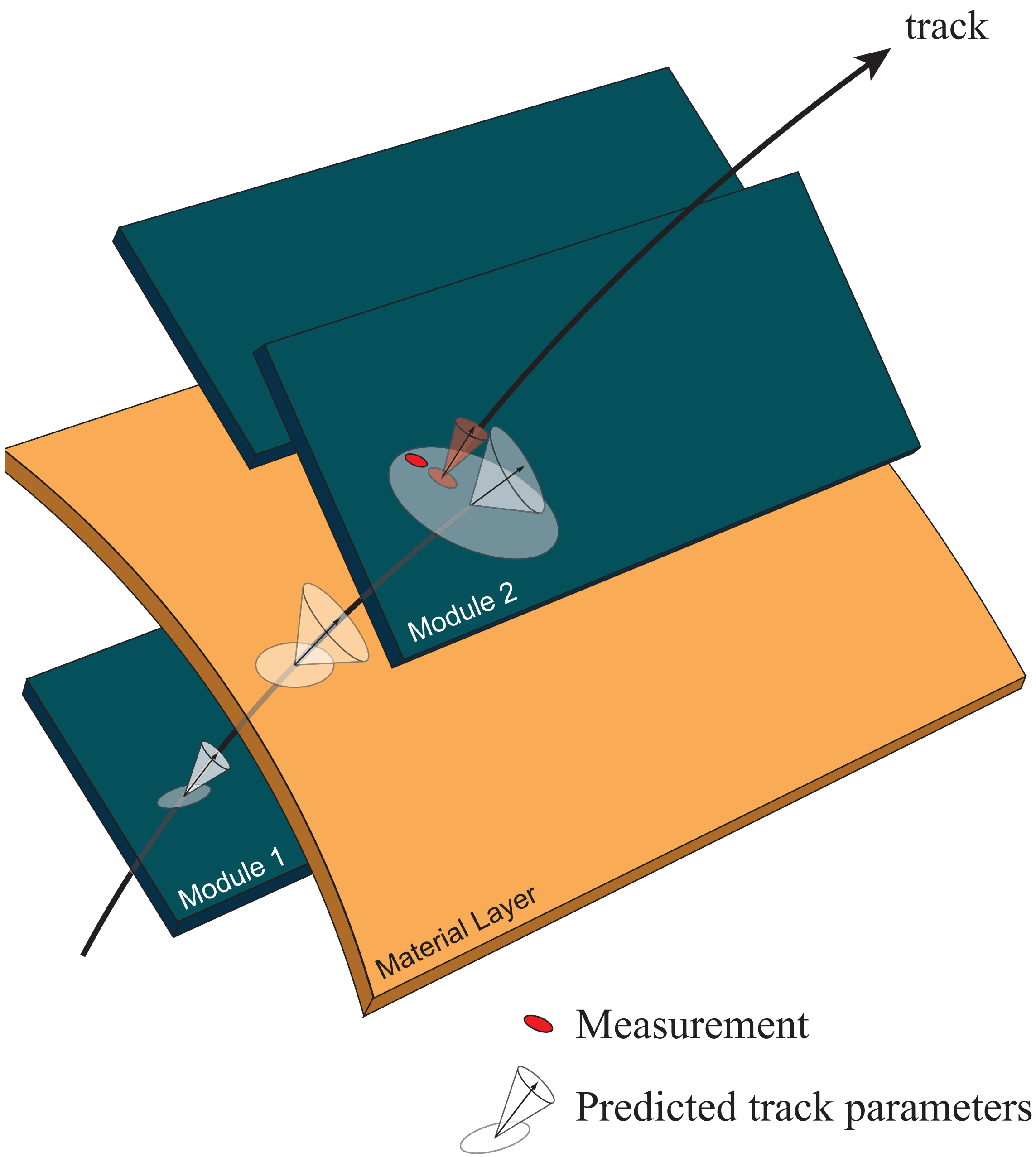


- Developed a GUI making it easy for shifters to archive scans to a database for monitoring long-term detector health.
- Still used in the regularly scheduled calibration periods between beam fills.
- Supported TRT as part of DAQ on-call team.



TRT hit efficiency

'09-'10

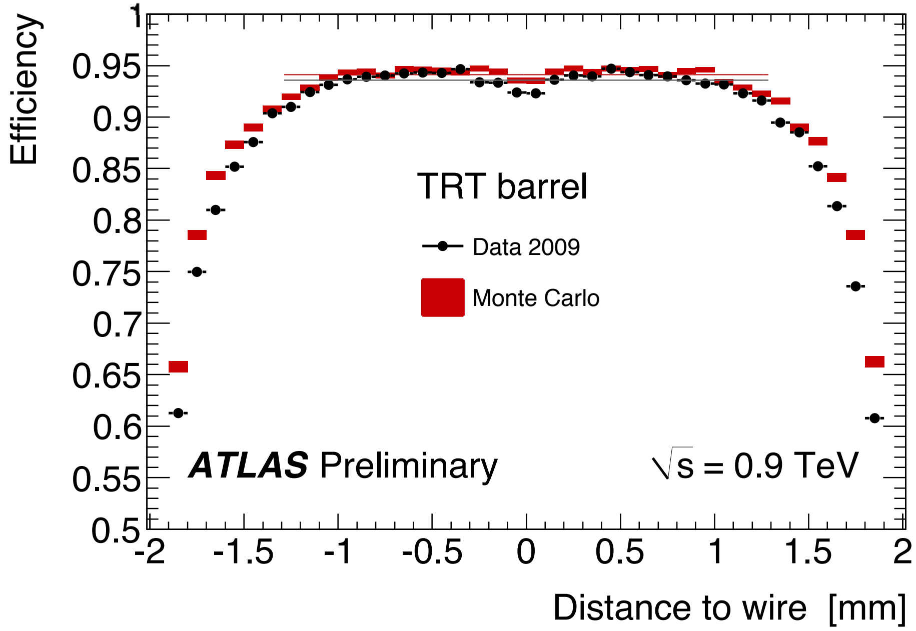


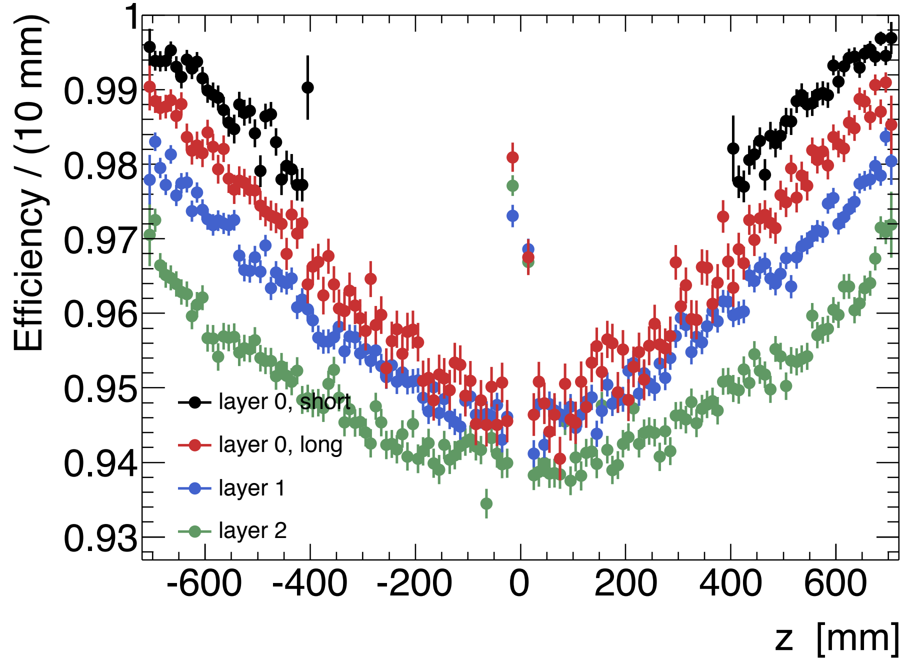
- $$\varepsilon(x) = \frac{n_{\text{hits}}(x)}{n_{\text{hits}}(x) + n_{\text{holes}}(x)}$$
- ***Hits*** are easy to count directly from the data.
- ***Holes*** are counted by extrapolating along a track to determine which straws were crossed.
- ***Extrapolation tools*** use a detailed description of detector material to stochastically model bremsstrahlung and multiple scattering.

[ATL-SOFT-PUB-2007-005, ATL-COM-INDET-2011-088]

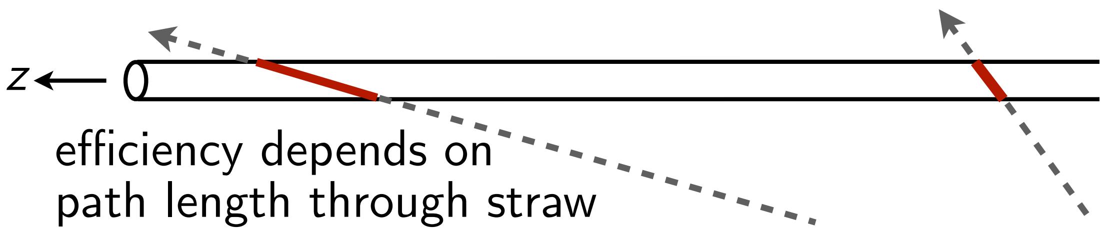
TRT hit efficiency

'09-'10





- I wrote an algorithm that uses extrapolator tools to calculate the TRT straw-hit efficiency.
- Gives an important data/MC comparison to test the TRT digitization, the step in the MC production where the response detector and electronics are simulated.
- ***Published in the first JHEP paper*** documenting the ATLAS detector performance with the first 900 GeV collision from 2009.



efficiency depends on path length through straw

[J. High Energy Phys. 09 (2010) 056 arXiv:1005.5254]